### A High-Resolution View of Fast Radio Burst Host Environments

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## ABSTRACT

We present Hubble Space Telescope (HST/WFC3) ultraviolet and infrared observations of eight fast radio burst (FRB) host galaxies with sub-arcsecond localizations, including the hosts of three known repeating FRBs. We quantify their spatial distributions and locations with respect to their host galaxy light distributions, finding that they occur at moderate host normalized-offsets of  $1.4r_e$  ([0.6, 2.1] $r_e$ ; 68% interval), occur on fainter regions of their hosts in terms of IR light, but overall trace the radial distribution of IR light in their galaxies. The FRBs in our tested distribution do not clearly trace the distributions of any other transient population with known progenitors, and are statistically distinct from the locations of LGRBs, H-poor SLSNe, SGRBs, and Ca-rich transients. We further find that most FRBs are not in regions of elevated local star formation rate and stellar mass surface densities in comparison to the mean global values of their hosts. We also place upper limits to the IR flux at the FRB positions of  $m_{\rm IR} \gtrsim 24.8 - 27.6$  AB mag, constraining both satellite and background galaxies to luminosities well below the host luminosity of FRB 121102. We find that 5/8 FRB hosts exhibit clear spiral arm features in IR light, and that the positions of all well-localized FRBs located in such hosts are consistent with their spiral arms, although not on their brightest regions. Our results do not strongly support the primary progenitor channel of FRBs being connected either with the most massive (stripped-envelope) stars, or with events which require kicks and long delay times (neutron star mergers).

Keywords: radio transient sources – Hubble Space Telescope

#### 1. INTRODUCTION

Since the discovery of fast radio bursts (FRBs) in 2007 (Lorimer et al. 2007) – dispersed, millisecondtimescale transients, traced primarily to cosmological distances (Thornton et al. 2013; Cordes & Chatter-

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jee 2019) – their definitive origins have remained elusive. The emerging association of FRBs with magnetized neutron stars ("magnetars") was catalyzed by the discovery and sub-arcsecond localization of the repeating FRB 121102 (Spitler et al. 2016; Chatterjee et al. 2017; Tendulkar et al. 2017), which was found to be spatially consistent with a compact, persistent radio source (Chatterjee et al. 2017) postulated to be a synchrotron nebula powered by an embedded magnetar (Marcote et al. 2017; Margalit & Metzger 2018) or an offset active galactic nucleus (AGN; Marcote et al. 2017; Eftekhari et al. 2020). The association of at least some FRBs with magnetars was further solidified by the detection of multiple FRB events from the Galactic magnetar SGR 1935+2154 (CHIME/FRB Collaboration et al. 2020; Bochenek et al. 2020). However, a myriad of progenitor models still exist (Platts et al. 2019; Petroff et al. 2019), and the precise connection to magnetars for the extragalactic population has yet to be made definitive. Moreover, given the diversity of their observed FRB properties, it is not clear if there is one or multiple progenitor channels for FRB production.

The advent of dedicated FRB experiments are fueling an unprecedented rise in their detection rates (CHIME/FRB Collaboration et al. 2018; Macquart et al. 2010). One of the most promising ways to make progress on their origins is through robust associations to host galaxies, which generally requires precise localizations of  $\leq 1''$  (Eftekhari & Berger 2017). Indeed, different progenitor channels will yield distinct host galaxy demographics and host stellar population properties (e.g., Margalit et al. 2019). It is also expected that different production pathways will be imprinted in how FRBs are distributed with respect to their host galaxies.

Locations have historically played an important role in delineating the progenitors of a wide range of transients. For instance, the spatial distributions and strong correlation with the UV light distributions of their hosts for super-luminous supernovae (SLSNe) and long-duration gamma-ray bursts (LGRBs) is commensurate with their massive star origins (Woosley 1993; Fruchter et al. 2006; Lunnan et al. 2015; Perley et al. 2016), while the significant offsets of short-duration gamma-ray bursts (SGRBs) and weaker correlation to UV light is indicative of an older stellar progenitor (e.g. Fong & Berger 2013). Moreover, quantifying the relationship between the locations of core-collapse SNe (CCSNe, Types Ib/c, II) and Type Ia SNe, to detailed morphological features such as spiral arms, can serve as an indirect indicator of the age of their stellar and/or white dwarf progenitors (Audcent-Ross et al. 2020).

For FRBs, two primary pathways that have been considered for magnetar formation are "prompt" magnetars, formed from massive star progenitors, and "delayed"-channel magnetars, formed from existing compact objects and their interactions, e.g., neutron star mergers or accretion-induced collapse (AIC) of a white dwarf to a NS (Nicholl et al. 2017; Margalit et al. 2019). Detailed studies using ground-based observations based on the first ~dozen well-localized FRBs and their host galaxies have shown that their spatial distributions are inconsistent with engine-driven massive star explosions (LGRBs, SLSNe; Li et al. 2019; Bhandari et al. 2020a; Heintz et al. 2020).

One is therefore motivated to characterize, as precisely as possible, the local environments of FRBs within their host galaxies. Furthermore, the competing progenitor models offer distinct predictions for the ages and masses of the responsible compact object(s). For young progenitors, there must be a direct link to ongoing or recent star-formation activity. Older progenitor channels, meanwhile, may track the underlying stellar mass of the galaxy. Thus motivated, we have designed an experiment to examine the local environments of FRBs, in both active star-formation (via near-UV light) and stellar mass (near-IR) and at the highest spatial resolution afforded by space-borne instrumentation.

Against this backdrop, we present the first comprehensive sample of HST observations for eight FRB host galaxies, six of which are newly presented in this work. In Section 2, we describe our sample, observations and data reduction. In Section 3, we present the spatial distribution of FRBs (angular, physical and hostnormalized offsets) and the locations of FRBs with respect to their host light distributions. In Section 4, we present the results of surface brightness profile fits, including the revelation of spiral structure in several hosts; constraints on the star formation rate and stellar mass densities at the FRB positions; and deep limits on possible satellite or background galaxies. In Section 5, we discuss the implications of our results in terms of FRB progenitors. We highlight our main conclusions in Section 6. Throughout the paper, we employ a Planck cosmology with  $H_0 = 67.8 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_M = 0.308$ , and  $\Omega_{\Lambda} = 0.692$  (Planck Collaboration et al. 2016). All of the data and analysis code are made available via GitHub<sup>1</sup>.

### 2. DATA

### 2.1. Sample of FRB Host Galaxies

Here we present observations of eight FRB host galaxies obtained with the Wide-Field Camera 3 using the infrared and ultra-violet-visual channels (WFC3/IR and WFC3/UVIS). The data for six of the FRB host galaxies were collected between October 2019 and April 2020 as part of programs 15878 (PI: Prochaska) and 16080 (PI: Mannings), which targeted galaxies for which FRB events have been detected and localized by the Commensal Real-time ASKAP Fast Transients (CRAFT) survey

<sup>1</sup> https://github.com/FRBs/FRB

### HST FRB HOST GALAXIES

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 Table 1. Log of Hubble Space Telescope FRB Host Galaxy observations

FRB	$\mathrm{RA}_{\mathrm{Host}}$	$\mathrm{Dec}_{\mathrm{Host}}$	$\sigma_{ m Host}$	z	$M_*$	Instrument	Filter	Date	Exp. Time
	(J2000)	(J2000)	(mas)		$(10^9 M_{\odot})$			(UT)	(sec)
121102	$05^{\rm h}31^{\rm m}58.69^{\rm s}$	$+33^{\circ}08'52.43''$	6.3	0.1927	$0.14\pm0.07$	WFC3/IR	F160W	2017 Feb 23	1197
						WFC3/UVIS	F763M	$2017 \ {\rm Feb} \ 23$	1940
180916	$01^{ m h}58^{ m m}00.29^{ m s}$	$+65^{\circ}42'53.09''$	1.8	0.0337	$2.15\pm0.33$	WFC3/IR	F110W	2020 Jul 17	5623
						WFC3/UVIS	F673N	2020 Jul 16	2877
180924	$21^{\rm h}44^{\rm m}25.256^{\rm s}$	$-40^{\circ}54'00.80''$	0.4	0.3212	$13.23\pm5.06$	WFC3/IR	F160W	$2019~{\rm Nov}~27$	2470
			3.1			WFC3/UVIS	F300X	$2019~{\rm Nov}~26$	2492
190102	$21^{\rm h}29^{\rm m}39.577^{\rm s}$	$-79^{\circ}28'32.52''$	14.2	0.2912	$3.39 \pm 1.02$	WFC3/IR	F160W	2020Jan $14$	2470
			†			WFC3/UVIS	F300X	$2019 \ \mathrm{Oct} \ 07$	2776
190608	$22^{\rm h}16^{\rm m}04.903^{\rm s}$	$-07^{\circ}53'55.91''$	0.8	0.1177	$11.57\pm0.84$	WFC3/IR	F160W	$2019 \ \mathrm{Dec} \ 01$	2295
			0.8			WFC3/UVIS	F300X	2019  Oct  11	2400
190711	$21^{h}57^{m}40.613^{s}$	$-80^{\circ}21'29.05''$	5.2	0.522	$0.81\pm0.29$	WFC3/IR	F160W	$2020 {\rm \ May\ } 11$	2470
			7.7			WFC3/UVIS	F300X	$2020 {\rm \ May\ } 09$	2780
190714	$12^{\rm h}15^{\rm m}55.090^{\rm s}$	$-13^{\circ}01'15.96''$	1.1	0.2365	$14.92 \pm 7.06$	WFC3/IR	F160W	$2020~{\rm Apr}~30$	2295
			†			WFC3/UVIS	F300X	$2020 {\rm \ May\ } 19$	2396
191001	$21^{h}33^{m}24.440^{s}$	$-54^{\circ}44'54.53''$	0.5	0.2340	$46.45 \pm 18.80$	WFC3/IR	F160W	$2020~{\rm Apr}~28$	2296
			†			WFC3/UVIS	F300X	$2020~{\rm Apr}~25$	2580

NOTE—Data are from programs 15878 (FRB 180924, 190102, 190608, PI: Prochaska), 16080 (FRBs 190711, 190714, and 191001, PI: Mannings), 14890 (FRB 121102, PI: Tendulkar), and 16072 (FRB 180916, PI: Tendulkar).

 $\sigma_{\text{Host}}$  is the 1 $\sigma$  positional uncertainty of the host (RA and Dec components added in quadrature).  $M_*$  is the host stellar mass.

<sup>†</sup> S/N of host galaxy is not sufficient to obtain an uncertainty on the position.

Redshift References- Tendulkar et al. (2017); Bannister et al. (2019); Marcote et al. (2020); Bhandari et al. (2020a); Heintz et al. (2020)

on the Australian Square Kilometer Array Pathfinder (ASKAP; Bannister et al. 2019; Day et al. 2020; Bhandari et al. 2020a; Chittidi et al. 2020; Macquart et al. 2020). These bursts were localized to sub-arcsecond precision, with  $\sigma_{\rm FRB} \approx 0.1 - 0.7''$ .

We also include two additional FRB hosts with HSTobservations,  $FRB121102^2$  (Bassa et al. 2017) under program 14890 (PI: Tendulkar) taken in February 2017 and FRB 180916 (Tendulkar et al. 2020) under program 16072 (PI: Tendulkar) taken in July 2020. FRB 121102 was discovered by the Arecibo telescope (Spitler et al. 2016), and subsequently localized via repeating bursts with the Very Large Array (VLA; Chatterjee et al. 2017), and with Very Long Baseline Interferometry (EVN VLBI; Marcote et al. 2017) with  $\sigma_{\text{FRB}} = 0.0045''$ . FRB 180916 is the closest and most precisely localized FRB with  $\sigma_{\text{FRB}} = 0.0023''$  (Marcote et al. 2020). Our sample thus comprises all eight FRB host galaxies for which there exist available HST observations. Table 1 summarizes all of these data and Table 2 summarizes coordinates and the localization errors of the FRBs.

All of the host galaxies in our HST sample have spectroscopically-confirmed redshifts. These are considered secure associations<sup>3</sup> with probabilities of chance coincidence of  $P_{\text{chance}} \leq 0.05$  (Heintz et al. 2020) with their most likely host galaxy. The HST data for FRB 121102 and FRB 180916 were previously published in Bassa et al. (2017) and Tendulkar et al. (2020), respectively, while the WFC3/UVIS image for FRB 190608 and its local environment was previously published and analyzed in Chittidi et al. (2020). All of the remaining HST observations are newly presented here. Three of the bursts are known "repeating" FRBs (FRBs 121102, 190711 and 180916; Spitler et al. 2016; Kumar et al. 2020) while the remaining bursts are apparent "non-repeaters".

We supplement this sample with ground-based data from two other FRB hosts presented in Heintz et al. (2020) with secure host associations (FRB 190611 and FRB 200430) when computing cumulative distributions of offsets in Section 3.2. Both of these FRBs in the ground-based sample are apparent non-repeaters. Combined, our ground-based and HST sample comprises all FRBs with sub-arcsecond localizations discovered over

 $<sup>^2</sup>$  The Transient Name Server (TNS) name for this burst is FRB 20121102a.

 $<sup>^3</sup>$  See also Aggarwal et al. (2021) for a Bayesian analysis that reaches similar conclusions.

Table 2. FRB Sample and Localizations

FRB	$RA_{FRB}$	$\mathrm{Dec}_{\mathrm{FRB}}$	$a_{\rm stat}$	$a_{\rm sys}$	$b_{\rm stat}$	$b_{\rm sys}$	PA
	(J2000)	(J2000)	('')	('')	('')	('')	(deg)
121102	82.994589	33.1479316	0.004		0.002		90.0
180916	29.503126	65.7167542	0.002		0.002		0.0
180924	326.105229	-40.9000278	0.07	0.09	0.06	0.07	0.0
190102	322.415667	-79.4756944	0.21	0.52	0.17	0.44	0.0
190608	334.019875	-7.8982500	0.19	0.19	0.18	0.18	90.0
190611	320.745458	-79.3975833	0.34	0.60	0.32	0.60	0.0
190711	329.419500	-80.3580000	0.12	0.38	0.07	0.32	90.0
190714	183.979667	-13.0210278	0.17	0.32	0.10	0.23	90.0
191001	323.351554	-54.7477389	0.13	0.11	0.08	0.10	90.0
200430	229.706417	12.3768889	0.01	0.02	0.24	1.00	0.0

NOTE— $a_{\text{stat}}, a_{\text{sys}}$  are the angular size of the semi-major axis describing the  $1\sigma$  statistical and systematic uncertainties respectively. *b* refers to the semi-minor axis. PA is the sky position angle of the error ellipse. Sources without reported systematic error have been incorporated in the statistical. Data from Day et al. (2020); Tendulkar et al. (2017); Marcote et al. (2020); Heintz et al. (2020).

2012-2020. The exception is FRB 190614D which does not have a clear host galaxy association (Law et al. 2020), and is not included in the sample.

#### 2.2. Observations

For the WFC3/UVIS observations under programs 15878 and 16080, we use the ultra-wide F300X filter to sample the rest-frame near- and far-ultraviolet (NUV/FUV) wavelengths with the aim of capturing the distribution of star formation in the host galaxies. This filter provides increased throughput in the NUV compared to the standard wide filters (although it has a minor red tail out to  $\sim 4000$  Å), and is chosen to maximize the signal-to-noise (S/N) in a single orbit of HSTimaging. To minimize the effects of charge transfer efficiency (CTE) degradation, we set up the observations to position the targets near the readout on amplifier C located on chip 2. We used a 4-point dither pattern to sub-sample the point spread function (PSF) and remove detector artifacts. We increase the line and point spacing by a factor of 5 over the standard box pattern to remove residual background patterns as described in Rafelski et al. (2015). The data from program 15878 include a 9 e<sup>-</sup> post-flash per exposure to reach 12 e<sup>-</sup> per pixel background. Recently there was a new recommendation to reach a background of 20e<sup>-</sup>, and therefore the

data obtained in program 16080 included a  $17 e^-$  post-flash to reach this level (Anderson et al. 2012).

For the WFC3/IR observations under programs 15878 and 16080, we use the F160W filter, the reddest wide filter available with HST, to cover the rest-frame optical band to assess the distribution of the stellar mass as traced by older stellar populations. We use SPARS25 and NSAMP15 to ensure that the observations remain in the linear count regime, and obtain clean images by dithering over the known IR "blobs" with a seven-point dither pattern with a factor 3 increase in spacing of the 7 point wide dither pattern provided in Anderson (2016). Finally, the data for FRB 121102 under program 14890 employ a two-point and four-point dither pattern for the WFC3/UVIS and IR observations, while the data for 180916 employ a three point and four-point dither pattern, respectively (Bassa et al. 2017; Tendulkar et al. 2020). The details of these observations are listed in Table 1.

#### 2.3. Image Processing

The data were retrieved from the Barbara A. Mikulski Archive for Space Telescopes (MAST), and the WFC3/UVIS data are custom calibrated. These data have degraded CTE, and therefore require pixel-based CTE corrections (Anderson et al. 2012). In addition, the degradation requires improved dark, hot pixel, and bias level corrections. First, we use a new correction for the CTE. Second, we apply concurrent superdarks to the data, reducing the blotchy pattern otherwise present (Rafelski et al. 2015). Third, we identify hot pixels in the darks such that the number of hot pixels is consistent as a function of the distance to the readout amplifiers based on the number of hot pixels identified close to the readout. This is accomplished by modifying the threshold for hot pixel detection as a function of distance to the readout (Prichard et al. in prep). Lastly, we normalize the amplifiers since the applied superbias is based on bias files with insufficient background levels for a pixel based CTE correction. We measure the background level in each amplifier after masking sources, and match the background level between the amplifiers.

To combine the images for each FRB and in each filter, we used the AstroDrizzle routine as part of the DrizzlePac software package (Avila et al. 2015) employing a pixfrac= 0.8, pixscale= 0.033" for UVIS images, and 0.064" for IR images. As part of AstroDrizzle, we also perform cosmic ray removal and sky subtraction. The final drizzled images are shown in Figures 1-3.

### 3. FRB LOCATIONS & OFFSETS

In this section, we present the locations of the FRBs in our sample with respect to their host galaxy centers. We introduce their angular and physical offsets ( $\theta$  and  $\delta R$ , respectively), "host-normalized" offsets ( $\delta R/r_e$ ) which are normalized by the half-light radii  $r_e$  of their host galaxies, the location with respect to their host galaxy light distribution ("fractional flux";  $F_{\rm F}$ ), and the fraction of light enclosed  $F_{\rm E}$  within the radius of the FRB.

## 3.1. Astrometry & Uncertainties

We first perform absolute astrometry using sources in common between available deep, optical groundbased imaging and the Gaia DR2 catalog. The ground-based imaging is sourced from the Gemini-North Observatory (FRB 121102), Gemini-South Observatory (FRB 190711). Very Large Telescope (VLT: FRBs 190102, 190714, 191001), the Dark Energy Survey (DES; FRB 180924) and the Sloan Digital Sky Survey (SDSS; FRB 190608). We then perform relative astrometry to tie the ground-based images to the WFC3/IR images. This series of astrometric ties ensures that there are enough sources in common with the WFC3 imaging, which has a significantly narrower field-of-view than the ground-based imaging, to properly calculate the total astrometric uncertainty. For astrometry, we employ a six-order polynomial accounting for linear shifts, rotation, and skew, using IRAF/ccmap and ccsetwcs. We calculate a range of tie uncertainties in each coordinate, of  $\sigma_{\rm tie,RA} \approx 0.014 - 0.073''$  and  $\sigma_{\rm tie,Dec} \approx 0.015 - 0.097''$ .

We also use the SourceExtractor tool (Bertin & Arnouts 1996) to determine the FRB galaxy centroid positions and associated uncertainties ( $\sigma_{\text{host}}$ ). These positions and values for  $\sigma_{\text{host}}$  are listed in Table 1. The final source of uncertainty is the positional uncertainty of the FRB ( $\sigma_{\text{FRB}}$ ), which is derived from the statistical and systematic uncertainties from the FRB detections (Table 2).

## 3.2. Physical & Host-normalized Offsets

Given that FRB localizations are typically noncircular (elliptical) in shape, and that their values span a range ( $\sigma_{\rm FRB} \approx 0.0023 - 0.7''$ ), it is necessary to take their shape, size, and orientation into account when calculating the angular, physical, and host-normalized offsets. While the synthesized beam and hence localization ellipse of the FRB can in principle have any position angle, most FRB detections have been made with beams that are close to circular and report the positional uncertainties projected onto the right ascension and declination axes, and we construct our localization ellipses using these projected values. To determine the total uncertainty on offset measurements, we sum each of the RA and Dec components of the three sources of uncertainty,  $\sigma_{\rm tie}, \sigma_{\rm host}, \text{ and } \sigma_{\rm FRB}$  in quadrature. We use the total uncertainties in RA and Dec to define an ellipsoidal region that represents the FRB location on the HST image.

The estimated angular offset  $\theta$  is then the convolution of the offset from the galaxy centroid  $\alpha_g$  with the FRB localization:

$$\theta = \int d\omega \, |\alpha_g - \omega| L(\alpha_{\rm FRB} - \omega) \tag{1}$$

with L a 2-D Gaussian set by the ellipsoidal region described above. To evaluate this convolution, we divide each 5- $\sigma$  region around the FRB into four million grid points by imposing a 2000 × 2000-point sub-grid. We measure the angular offset between each grid-point *i* and the host galaxy center to obtain a distribution of angular offsets  $\theta_i$  for each FRB. Finally, we apply a 2-D Gaussian probability distribution within the FRB localization ellipse, centered on the central RA and Dec of the FRB, and weight the angular offset distribution by the corresponding values. We estimated the variance in  $\theta$  in a similar manner and report the RMS in Table 3.

For each FRB, we determine the median offset and standard deviation. We find a range of projected angular offsets of  $\theta \approx 0.23 - 7.87''$  with a population median and 68% interval in the IR of 1.1'' and [0.6'', 2.9'']. The values for each FRB are listed in Table 3. We note that we only obtained offsets for observations in which the host galaxy center could be well determined, so this in-



Figure 1. *HST* imaging of three of the host galaxies in our sample, for FRB 180924, FRB 190102, and FRB 190106. Images with a blue color-bar were taken with the UV channel F300X filter, while images in the red color-bar were taken with the IR channel F160W filter. The black ellipse in each image denotes the FRB position ( $2\sigma$  uncertainty in each coordinate). All images are oriented with North up and East to the left.



Figure 2. *HST* imaging of three of the host galaxies in our sample, for FRB 190711, FRB 190714, and FRB 191001. Color scheme and ellipses are as in Figure 2.



Figure 3. *HST* IR imaging of the host galaxies of FRB 121102 (F160W) and FRB 180916 (F110W). The black dashed ellipse represents 30 times the  $1\sigma$  uncertainty from the VLBI localization for FRB 121102 (Marcote et al. 2017) and 200 times the  $1\sigma$  uncertainty for FRB 180916 (Marcote et al. 2020).



Figure 4. Left: Cumulative distribution of projected physical offsets,  $\delta R$ , for the 10 FRBs in the HST and ground-based samples (black line). The gray shaded region is a bootstrap estimate of the RMS of the distribution, which accounts for both uncertainties on individual measurements, as well as statistical uncertainties due to the sample size. Comparison samples are included for SGRBs (Fong et al. 2010; Fong & Berger 2013), LGRBs (Blanchard et al. 2016), Ca-rich transients (Lunnan et al. 2017; De et al. 2020), Type Ia SNe (Uddin et al. 2020), CCSNe (Schulze et al. 2020), and SLSNe (Lunnan et al. 2015; Schulze et al. 2020) for events at z < 1. The computed *P*-values from a two-sided KS test are listed for each population relative to the FRB sample. Right: The same as the right-hand panel but for the host-normalized offsets ( $\delta R/r_e$ ). This plot also shows the profile of an exponential disk.

cludes all hosts for which there are IR images, as well as the host galaxy of FRB 190608 in both the IR and UV.

We convert the angular offsets to projected physical offsets using the redshift of each FRB host galaxy and a Planck cosmology with  $H_0 = 67.8 \text{ km s}^{-1} \text{ Mpc}^{-1}$ ,  $\Omega_M = 0.308$ ,  $\Omega_{\Lambda} = 0.692$  (Planck Collaboration et al. 2016). For the *HST* offsets, we find a range of  $\delta R \approx 0.75 - 10.5$  kpc with the lower and upper bounds set by

FRB 121102 and FRB 191001, respectively<sup>4</sup>. The median and 68% interval on the projected physical offset is 3.2 kpc and [2.0, 8.7] kpc. Finally, we use the host galaxy half-light radii ( $r_e$ ), as measured from *HST* imaging (see Section 4.1) to determine the host-normalized

 $<sup>^4</sup>$  We note that the ground-based determination for FRB 190611 is the largest physical offset, with  $\approx 11.4$  kpc.

### HST FRB HOST GALAXIES

Table 3.	Offsets	and	Light	Locations	of	FRBs
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FRB	Filter	θ	$\delta R$	$\delta R/r_e$	F	$F_{\rm E}$
		('')	(kpc)			
121102	F160W	$0.23\pm0.02$	$0.75{\pm}~0.05$	$1.14\pm0.08$	$0.70\pm0.07$	$0.24\pm0.12$
	F763M				$0.67\pm0.10$	
180916	F110W	$7.869 \pm 0.002$	$5.461{\pm}~0.002$	$1.531 \pm 0.000$	$0.32\pm0.07$	$0.90\pm0.09$
	F673N				$0.32\pm0.25$	
180924	F160W	$0.71\pm0.12$	$3.37{\pm}~0.56$	$1.23 \pm 0.20$	$0.24\pm0.11$	$0.66\pm0.07$
	F300X					
190102	F160W	$0.80\pm0.39$	$2.26{\pm}~2.22$	$0.51\pm0.50$	$0.39\pm0.25$	$0.25\pm0.16$
	F300X				$0.36\pm0.29$	
190608	F160W	$2.98\pm0.27$	$6.52{\pm}~0.60$	$2.30\pm0.21$	$0.19\pm0.06$	$0.82\pm0.08$
	F300X				$0.39 \pm 0.28$	
190611	$GMOS-S_r$	$2.23\pm0.65$	$11.36\pm3.53$	$5.29 \pm 1.64$		
190711	F160W	$0.53\pm0.27$	$1.94{\pm}~2.30$	$0.66\pm0.78$	$0.55\pm0.27$	$0.61\pm0.22$
	F300X					
190714	F160W	$0.61\pm0.29$	$1.97{\pm}~1.18$	$0.50\pm0.30$	$0.34\pm0.23$	$0.23\pm0.10$
	F300X				$0.38\pm0.31$	
191001	F160W	$2.74\pm0.15$	$10.49 \pm\ 0.59$	$1.89\pm0.11$	$0.09\pm0.03$	$0.91\pm0.09$
	F300X				$0.29\pm0.25$	
200430	$\operatorname{Pan-STARRS}_r$	$1.30\pm0.79$	$2.97{\pm}~2.36$	$1.81\pm1.44$		
Median (IR)		1.1	3.2	1.4	0.33	0.6
$16{,}84\%$ Interval		[0.6, 2.9]	[2.0, 8.7]	[0.6, 2.1]	[0.20, 0.53]	[0.2, 0.9]
Median~(UV)					0.37	
$16{,}84\%$ Interval					[0.32, 0.45]	

NOTE—FRBs 190611 and 200430 are derived from ground-based imaging as reported in Heintz et al. (2020)

offsets,  $\delta R/r_e$ . The values for the projected angular, physical, and host-normalized offsets for the eight FRBs in our sample are listed in Table 3.

We supplement the FRB distributions with two FRB host galaxies in Heintz et al. (2020), FRBs 190611 and 200430, both of which have offsets determined from ground-based imaging with seeing of  $\sim 0.8''$ . To determine the uncertainty on the cumulative distribution, we follow the method by Palmerio et al. (2019) and create 10,000 realizations of asymmetric Gaussian PDFs using the errors on the offset measurements for each FRB, derived from the previously described weighted grid analysis. We then use a bootstrap method to sample from the PDF in each realization, allowing us to compute a CDF of the bootstrapped sample. Finally, we compute the median of all the resulting CDFs, as well as the upper and lower bounds for each bin. We perform this same analysis for the projected physical and host-normalized offset distributions. The resulting median cumulative distributions, and bootstrap estimate of the uncertainty (shown as the shaded gray region) are shown in Figure 4.

To compare the FRB distribution to the offset distributions of other transients, we draw relevant comparison samples from the literature. Included are longduration gamma-ray bursts (LGRBs; Blanchard et al. 2016, short-duration gamma-ray bursts (SGRBs; Fong et al. 2010; Fong & Berger 2013), Ca-rich transients (Lunnan et al. 2017; De et al. 2020), Type Ia supernovae (Type Ia SNe; Uddin et al. 2020), core-collapse SNe (CCSNe; Schulze et al. 2020) and super-luminous SNe (SLSNe; Lunnan et al. 2015; Schulze et al. 2020). To align with the redshift distributions of the FRBs, we only include values for events with z < 1. We perform a two-sided KS-test between the median FRB distribution and each of the transient populations to test the null hypothesis that the (median) distribution of FRBs and each transient population is drawn from the same underlying distribution. Using this analysis for projected offsets, we find  $P_{\rm KS} < 0.05$  for both LGRBs and SLSNe, rejecting the null hypothesis that they are drawn from the same continuous distributions. We caution, however, that we only report KS-test results on the median FRB distribution.

However, for host-normalized offsets, we also find  $P_{\rm KS} < 0.05$  for Ca-rich transients and for LGRBs. The remaining *p*-values are  $P_{\rm KS} > 0.05$ , and thus we cannot reject the null hypothesis for any other population tested. Finally, we compare the distribution to an exponential disk light profile (light purple curve in Figure 4). While the distribution overall appears to be at larger offsets, the  $P_{\rm KS} = 0.066$  value is not conclusive. In this analysis, we note that we are treating the FRB population as a single distribution with a dominant progenitor population. However, the sample size considered here prevents meaningful constraints on the presence of multiple, equally dominant progenitor populations, each of which have distinct offset distributions.

#### 3.3. Fractional Flux

We now explore the location of the FRBs with respect to their host galaxy light distributions ("fractional flux";  $F_{\rm F}$ ). The brightness of the burst site in relation to how its rest-frame UV and optical host light is distributed is a crucial tool for determining how star formation activity and stellar mass are tracked (Fruchter et al. 2006). Compared to offsets, which can depend on host size and morphology, the fractional flux method is independent of these physical characteristics. Specifically, the measurement determines the fraction of host light fainter than the flux at the burst position, where a value of unity corresponds to the brightest light level of the host ( $F_{\rm F} = 1$ ).

Foreground stars in the field of FRB 180916 preclude the direct application of this analysis to the field. Since PSF photometry and subtraction of many of these stars were difficult as they were saturated; in this case we decided to use an alternate, "brute-force" approach. We first performed an isophote fit to the galaxy using photutils (Bradley et al. 2020), clipping pixels which were over the  $3\sigma$  level compared to the local mean. This left all the foreground stars and subtracted most of the galaxy light. Then we created a segmentation map from the residual image with a threshold level of  $4\sigma$  over the sky background and minimum source area of 5 pixels (the default value). From the objects extracted, we selected those with a peak pixel value of  $1 e^-$  per sec or greater and created a masked image. This masked out all star light above the segmentation map threshold value. We then replaced the masked pixels with the isophote fit from earlier and used this stitched image for the fractional flux analysis.

We center a 2D cutout on each host galaxy, making sure that we include a sufficient amount of background. We then determine which pixels lie within the  $3-\sigma$  FRB localization ellipse. We note that unlike how we determined the offsets (Section 3.2), we do not apply a sub-grid to the localization, as the main limitation is the pixel scale, and we cannot resolve the fractional flux below this scale. The fractional flux for each *i*th pixel on N pixels within the localization is then calculated as

$$F_{\rm F} = \frac{\sum_{i} (F_i < \text{limit})}{\sum_{i} F_i}.$$
 (2)

We then use a 2-D Gaussian distribution to develop a weighting scheme with each FRB localization ellipse in the same manner as that described in Section 3.2. From the distribution of  $F_{\rm F}$  values for each FRB, we determine the median  $F_{\rm F}$  and its standard deviation. The values for the fractional flux for each FRB can be found in Table 3.

Figure 5 shows the  $F_{\rm F}$  CDF for the sample of eight HST hosts from the IR images (corresponding to restframe optical; right panel), and four hosts from the UV images (left panel). The gray region was generated in the same manner described in Section 3.2. We exclude the UV imaging of FRBs 180924 and 190711 from this analysis because they are effectively non-detections with very large error, while we include the  $H\alpha$  imaging for FRB 180916 as a proxy for star formation (Tendulkar et al. 2020). Given that the UV and IR imaging can be used as proxies for star formation and stellar mass, respectively, we keep the two wavelength regimes separated. In this method, the 1:1 dashed line represents a population of events which traces the light of its host galaxy in that band. Adherence to the 1:1 line would indicate that FRBs may trace the distribution of starforming regions and stellar mass of their galaxies, respectively.

For the IR (rest-frame optical) distribution which can be used as a proxy for stellar mass, the  $F_{\rm F}$  values span a wide range,  $\approx 0.09 - 0.70$ , where the lower and upper bounds are set by FRB 190714 and FRB 121102, respectively (Table 3). The median of the distribution is 0.33 with a 68% interval of [0.20, 0.53].



Figure 5. Left: Cumulative distribution of UV fractional flux  $(F_{\rm F})$  of the 6 FRB hosts in the HST sample with WFC3/UVIS imaging. The gray shaded region is a bootstrap estimate of the RMS of the distribution, which accounts for both uncertainties on individual measurements, as well as statistical uncertainties due to the sample size. For comparison, the corresponding distributions for SGRBs (Fong et al. 2010; Fong & Berger 2013), LGRBs (Blanchard et al. 2016), Type Ia-SNe (Wang et al. 2013), CCSNe (Svensson et al. 2010), and SLSNe (Lunnan et al. 2015) are shown. Also shown is the dashed, 1:1 line, representing the distribution of host galaxy light. The *p*-values from two-sided KS tests between each population and the median FRB distribution are listed. Right: The same as the left-hand panel but for all eight FRB hosts with HST/IR imaging.

Figure 5 highlights that the median distribution of FRBs overall traces the fainter rest-frame optical regions of their host galaxies, with a location to the left/above of the 1:1 line. However, the relatively large positional uncertainties which extend to  $\sigma_{\rm FRB} \approx 0.7''$ , coupled with the small sample size of eight events leads to a non-trivial uncertainty in the distribution, which is consistent with the 1:1 line. Thus, while FRBs appear to trace the fainter regions of their hosts in terms of stellar mass, at present it is not possible to make a strong statistical statement.

For comparison, we draw  $F_{\rm F}$  measurements from the literature for SGRBs (Fong et al. 2010; Fong & Berger 2013), LGRBs (Blanchard et al. 2016), SLSNe (Lunnan et al. 2015), CCSNe (Svensson et al. 2010), and Type Ia SNe (Wang et al. 2013) and divide them into rest-frame UV and optical measurements for direct comparisons. Performing two-sided KS tests with respect to the median FRB distribution, we can rule out the null hypothesis that the LGRBs, SGRBs and FRBs are from the same underlying population as they all yield *p*-values of  $P_{KS} < 0.05$ . For CCSNe, Type Ia SNe, and the 1:1 distribution, we find  $P_{KS} > 0.05$  and cannot rule out the null hypothesis.

For the UV distribution, which can be used as a proxy for the distribution of current star formation, there are six data points and their median values fall in a fairly narrow range of  $\approx 0.29$  (for FRB 191001) to  $\approx 0.67$  (for FRB 121102). Despite the small sample size, the KS test does reject the null hypothesis that both LGRBs and SLSNe and the median FRB distribution come from the same underlying population with  $P_{KS} < 0.05$ . However, we caution that the very small sample size coupled with large localization uncertainties (which translate to large measurement uncertainties in  $F_{\rm F}$ ) effectively means that almost the entire parameter space of  $F_{\rm F}$  is included in the uncertainty region. We also note that the IR  $F_{\rm F}$ contains all three repeating FRBs, whereas the UV  $F_{\rm F}$ contains two.

#### 3.4. Fraction of enclosed light

While the  $F_{\rm F}$  metric is designed to assess the local environment of the FRB in the galaxy *independent* of its morphology, we introduce an alternative metric to better assess its global position: the fractional flux enclosed

$$F_{\rm E} = \frac{\sum\limits_{i < r_{\rm FRB}} F_i}{\sum\limits_{i} F_i}$$
(3)

with  $r_i$  the radius of pixel i and  $r_{\rm FRB}$  the distance of the FRB from the galaxy centroid. In practice, we estimate  $F_{\rm E}$  using the isophotal analysis of Section 4.1. Specifically, we find the isophote closest to the FRB localization and measure the flux enclosed within it (and all interior isophotes). Furthermore, we allow for the FRB localization error by performing a weighted average of individual  $F_{\rm E}$  evaluations across the FRB localization.

Figure 6. The median FRB cumulative distribution of enclosed flux (black line), which is the fraction of host light enclosed within a radius set by the position of the FRB. The values are derived for eight host galaxies with IR images. The gray shaded region represents the  $1\sigma$  uncertainty in the CDF, while the blue dashed line indicates a 1:1 mapping. A KS test between the median distribution and 1:1 line does not reject the null hypothesis that the distributions are drawn from the same underlying population.

Figure 6 shows the  $F_{\rm E}$  results derived from the WFC3/IR images and a 1:1 line corresponding to the null hypothesis that FRBs are biased in tracking the host's light. The close correspondence between the two lends credence to the null hypothesis and the resultant  $P_{\rm KS} = 0.73$  value offers statistical support. Similar to Safarzadeh et al. (2020), we infer that FRBs track the general distribution of light and, by inference, stellar mass in their host galaxies.

# 4. MORPHOLOGICAL & FRB SITE PROPERTIES 4.1. Galaxy Light Profile Fitting

We fit the light profile for the eight FRB hosts with HST data to determine the half-light radii  $(r_e)$ , which are used to compute the host-normalized offsets. The half-light radii from HST imaging are valuable in comparison to previous ground-based imaging, where the measurements are based on seeing-limited images, allowing for more accurate estimates. The increased sensitivity of HST also presents the opportunity to search for alternate, fainter host galaxy candidates. Thus, in what follows, we also use light profile fitting to develop a galaxy model, and subtract it from the images to determine constraints on possible alternative, host galaxy candidates (e.g., low-luminosity galaxies or background galaxies) at the position of each FRB.

First, we use the elliptical isophotal model from photutils.isophote to map the light of the eight FRB hosts. We begin with an initial guess, providing values for the central position, ellipticity, semi-major axis, and position angle. The function then fits a series of isophotes which we then use to create a model and residual image. We determine the value of  $r_e$  from our isophotal fits, taking this to be the semi-major axis in which half of the total light is enclosed. These values are listed in Tables 4-5, and are used in our calculation of the host-normalized offsets (see Section 3.2).

In addition to the isophotal fits, we also compute residuals from Sérsic profile fits using GALFIT v.3 (Peng et al. 2010) for the IR images. Our model has two Sérsic components for all galaxies except the hosts of FRBs180924 and 190711. These two components roughly correspond to a central bulge and an outer disk. In the case of FRB 180924, we use a single Sérsic component because there is no obvious improvement in the residuals relative to performing a multi-component fit. In the case of the host galaxy of FRB 180924, the two components converged to the same effective radii, similar Sérsic indices and similar magnitudes, implying that GALFIT could not distinctly identify separate core and disk components. For FRB 190711, the relatively low signal-to-noise of the host galaxy image precludes the identification of two distinct components.

The residual images from the fits are shown in Figure 7. The GALFIT half-light radii  $(r_e)$  reported in Table 5 correspond to the larger component (i.e. the disk component) in the case of the two-component fits. While the residual images of FRBs 121102, 190711 and 190102 do not show clear, symmetric structures, there are such morphological structures in the residual images of the remaining hosts. In particular, spiral arm structure becomes very apparent for FRB 180924 and 190714, as well as the previously-known structure for FRBs 180916 (Tendulkar et al. 2020), 190608 (Chittidi et al. 2020) and 191001 (Bhandari et al. 2020b). Therefore we can see in Figure 7 that *all* of the FRBs which are localized to hosts with spiral structure land on or very near to a spiral arm; this point is discussed further in Section 5.

### 4.2. Star Formation Rate and Stellar Mass Constraints

In addition to relative photometric measures, these data enable precise photometric measurements at the burst positions. Furthermore, we may convert these light measurements into physical quantities – the starformation rate (SFR) density  $\Sigma_{\text{SFR}}$  and stellar mass surface density  $\Sigma_{\text{M*}}$  - to give additional insight into the local properties of FRB burst sites. It is also informative to compare the local values to the global mean surface



FRB	Filter	Host Magnitude	$A_{\lambda}$	$\mu_{ m FRB}$	$\Sigma_{\rm SFR(FRB)}$
		(AB mag)	(mag)	$(\mu Jy \ arcsec^{-2})$	$(M_\odot{\rm yr^{-1}kpc^{-2}})$
180924	F300X	$23.478 \pm 0.058$	0.12	< 0.85	< 0.025
190102	F300X	> 27.200	1.40	< 2.58	< 0.062
190608	F300X	$19.765\pm0.014$	0.30	< 1.73	$0.005\pm0.007$
190711	F300X	$25.008 \pm 0.121$	0.88	< 1.61	< 0.130
190714	F300X	$23.072 \pm 0.053$	0.39	< 1.12	< 0.018
191001	F300X	$21.228 \pm 0.020$	0.18	< 0.88	< 0.014

Table 4. Derived Properties from UV Observations of Host Galaxies

NOTE—Magnitudes are not corrected for Galactic extinction in the direction of the FRB  $(A_{\lambda})$ . Limits correspond to  $5\sigma$  confidence.

Table 5. Derived Properties from IR Observations of Host Galaxies

FRB	Filter	$r_e$ (Isophotal)	$r_e~({\tt GALFIT})$	Host Magnitude	Limit	$A_{\lambda}$	$\mu_{ m FRB}$	$\Sigma_{M*(FRB)}$
		$(\mathrm{kpc})$	$(\mathrm{kpc})$	(AB mag)	(AB mag)	(mag)	$(\mu Jy \ arcsec^{-2})$	$(10^8  M_\odot  {\rm kpc}^{-2})$
121102	F160W	$0.66 \pm 0.03$	$2.05\pm0.11$	$23.435 \pm 0.055$	27.4	0.40	$14.41\pm0.20$	$0.129\pm0.002$
180916	F110W	$3.57\pm0.36$	$6.00\pm0.01$	$16.178 \pm 0.005$	25.4	0.88	$31.00\pm0.10$	$0.115\pm0.000$
180924	F160W	$2.75\pm0.10$	$2.82\pm0.53$	$19.349\pm0.002$	26.2	0.01	$20.97\pm0.10$	$0.821 \pm 0.010$
190102	F160W	$4.43\pm0.51$	$5.00\pm0.15$	$20.550\pm0.006$	27.1	0.11	$11.03\pm0.11$	$0.092\pm0.002$
190608	F160W	$2.84 \pm 0.23$	$7.37\pm0.06$	$16.693 \pm 0.001$	25.2	0.02	$17.78\pm0.10$	$0.339\pm0.001$
190711	F160W	$2.94\pm0.17$	$2.48\pm0.13$	$22.899 \pm 0.014$	27.6	0.07	$6.19 \pm 0.11$	$0.045\pm0.004$
190714	F160W	$3.94\pm0.05$	$3.85\pm0.03$	$18.896\pm0.002$	25.9	0.03	$17.55\pm0.10$	$1.708\pm0.018$
191001	F160W	$5.55\pm0.03$	$6.23 \pm 0.04$	$17.135 \pm 0.001$	24.8	0.01	$7.91\pm0.10$	$0.319 \pm 0.007$

NOTE—Magnitudes are not corrected for Galactic extinction in the direction of the FRB  $(A_{\lambda})$ . Limits correspond to  $5\sigma$  confidence.

densities of their host galaxies, to understand if the locations of FRBs are, for example, elevated or not in terms of these quantities.

To complete the photometric measurements, we create apertures with r = 3 pixels at each pixel within the FRB localization. These aperture sums are then weighted by a 2-D Gaussian probability distribution centered on the measured FRB position, the same as that used for the offset and fractional flux determinations in Sections 3.2-3.3 – with the resolution of the grid being limited by the image pixel scale. We then take the weighted average and divide by the area of the aperture to get an aperture sum per arcsec<sup>2</sup>. We compute the magnitude using WFC3 tabulated zeropoints, the corresponding flux in both filters, and the luminosity for the UV band (Table 4). Next, we use the UV luminosity-SFR relation from Kennicutt (1998) to convert UV measures into star formation rate and subsequently the star formation rate surface density per kpc<sup>2</sup>,  $\Sigma_{\rm SFR}$  at the burst site in each of the HST hosts. For the stellar mass surface density per kpc<sup>2</sup>,  $\Sigma_{\rm M*}$  we compute the ratio of FRB flux to total host flux and adopt the total stellar mass estimates from Heintz et al. (2020) to estimate the local value. We derive  $3\sigma$  upper limits on star formation rate densities in the same manner, relevant for hosts that are non-detections or have very low S/N at the FRB site.

In terms of  $\Sigma_{\text{SFR}}$  Figure 8 shows that most FRBs do not obviously occur in elevated regions of star formation with respect to the global values of their host galaxies (albeit most of the UV values are upper limits). The two FRBs with measurements in our sample, FRBs 121102



Figure 7. Residual images produced by GALFIT from the F160W host galaxy image set (and FRB 180916 for F110W). The North and East directions are indicated by the black arrows at the top left. The light distribution for all galaxies was modeled as a sum of two Sérsic profiles corresponding to a central core and an outer disk, except in the host galaxies of FRBs 180924 and 190711. In those two galaxies, a single Sérsic profile fit was used because a two component fit did not produce a significant improvement in the visual quality of the residuals. It is interesting to note that five of the eight FRB locations, marked by the dashed  $2\sigma$  localization ellipses (200 $\sigma$  for 180916 and  $30\sigma$  for 121102), are coincident with spiral structures in their respective hosts.





Figure 8. Comparisons of local FRB properties to global host properties. The points with black outlines are "repeaters". The triangles denote  $3\sigma$  upper limits on the SFR. *Left:* Star formation rate surface density at FRB location versus star formation rate surface density of the host. Compared to the 1:1 line, the sites of FRBs are not clearly elevated in star formation rate surface densities with respect to their hosts, and do not reach the surface densities of Galactic star-forming regions (Evans et al. 2009). *Right:* Average stellar mass surface density at the burst site versus average stellar mass surface density of the host. Most FRB locations track the 1:1 line, with a few deviating from this relation. The average value for Milky Way GMCs is marked shown as the black star (Lada & Dame 2020).

and 190608 lie above the 1:1 line, in concert with previous conclusions (Bassa et al. 2017; Chittidi et al. 2020). We do however find that FRB 191001 lies below this 1:1 line. This burst also has the highest offset and, as is apparent in the UV image, is offset from the UV-bright regions of its host.

For context, Galactic star-forming clouds such as the Lupus and Perseus clouds (Evans et al. 2009) are shown to be well above the 1:1 line in comparison to the Milky Way average as reported in Kennicutt & Evans (2012). The FRB locations, except that of FRB 121102, do not reach these levels of elevated star-formation.

One caveat is that the Galactic star forming clouds and their measurements are derived from pc-scale measurements as opposed to the kpc scales for FRB localizations. Ideally, we would like resolve down to scales that are considered "local" in studies of star formation in the Milky Way. This will require additional ~mas-level localizations and larger aperture space or ground-based observations (e.g. JWST, ALMA).

In terms of stellar mass surface density  $\Sigma_{M*}$  Figure 8 reinforces several of the conclusions from Heintz et al. (2020). For example, a continuum of characteristics between "repeaters" and "non-repeaters" arises when investigating the stellar mass of the host and the burst site. Like the SFR density, FRBs also do not clearly occur in regions of elevated stellar mass surface densities with respect to the global values of their hosts, and only a few FRBs deviate from the 1:1 relation. We use Milky Way Giant Molecular Clouds (GMCs; Lada & Dame 2020) as a point of comparison to put into context the characteristics of these burst sites with other sites of star-formation. The stellar mass surface density for these sites relative to their hosts is slightly above, but are not very disparate from that shown for the Galactic GMCs which are  $\approx 35 M_{\odot}$  parsec<sup>-2</sup> as concluded in Lada & Dame (2020).

## 4.3. Luminosity Constraints on Satellite or Background Galaxies

With the early association of FRB 121102 to a very faint host, the community has adopted a bias towards associating FRBs with low luminosity galaxies. The subsequent association of FRBs to brighter galaxies (e.g. Bannister et al. 2019) has therefore led some to question whether a fainter, true host galaxy lurks below. To place constraints on an alternate, apparently fainter host galaxy candidate at the FRB position, we use the GALFIT residual images (Figure 7), in which the elliptical components from the bright, putative host galaxy have been removed to derive point-source limiting magnitudes  $m_{\rm lim}$  at the FRB position. We then compute the residual flux value using a circular aperture of 0.5'' diameter, corresponding to  $\sim 2.5$  times the PSF FWHM. We compute the net standard deviation for all pixels within this aperture. We then take the larger of the flux measurement and five times the net standard deviation as the upper limit on any point source flux that can be detected from the residual images (5 $\sigma$  limit). We find limits of  $m_{\text{lim}} \gtrsim 24.8 - 27.6 \text{ AB mag}$  (see Table 5).

We convert each of the at-position limits to an IR luminosity as a function of redshift (Figure 9). First, we explore these limits in the context of a spatially coincident satellite galaxy at the same redshift as the putative (brighter) host galaxy (triangles in Figure 9). At these redshifts, the limits of  $L_{\rm IR} \lesssim (0.5 - 9.2) \times 10^7 L_{\odot}$ , are significantly deeper than the luminosity of any known FRB host, including FRB121102. This means that despite the presence of morphological features in the GALFIT residuals which preclude extremely deep limits, we can still rule out a galaxy with similar luminosity to the host of FRB 121102, which is now considered an outlier in terms of FRB host stellar mass and luminosity (Tendulkar et al. 2017; Li et al. 2019; Bhandari et al. 2020a; Heintz et al. 2020). Any underlying host would need to have an IR luminosity of  $\lesssim 0.02 - 0.31$  times that of the host of FRB 121102 if it was at the same redshift of the brighter host galaxy (Figure 9). For reference, we measure IR luminosities for the putative hosts of  $\approx 3.0 \times 10^8 - 1.1 \times 10^{11} L_{\odot}$  (set by FRB 121102 and FRB 191001, respectively).

It is also worthwhile to explore whether or not a lowluminosity host galaxy of the same luminosity as the host of FRB 121102 may reside at a *higher* redshift than the apparently brighter galaxy. In this case, we find that the redshift of any background galaxy must be at  $z \gtrsim 0.4$ (Figure 9). The exception is FRB 180916, which still has a meaningful constraint of  $z \gtrsim 0.25$ . We also calculate the upper limit on the redshift inferred from the Macquart relation for each FRB following the methods of Macquart et al. (2020), and assuming a Milky Way DM  $= 50 \text{ pc cm}^{-3}$  and a host  $\text{DM} = 50/(1 + z_{\text{FRB}}) \text{ pc cm}^{-3}$ . This analysis results in limits of  $z \leq 0.17 - 0.75$  (95%) confidence). These redshift limits provide an absolute upper bound on the allowed luminosity of an underlying host of  $\approx 4 \times 10^7 - 8 \times 10^8 L_{\odot}$ . In all cases except for FRBs 190714 and 191001, we can thus rule out an underlying background galaxy of similar luminosity to the host of FRB 121102. We note that raising the required luminosity of the host galaxy would only push the required redshift to a higher range. We therefore find that the presence of background galaxies at higher redshifts are not likely for these FRBs given the HST limits and constraints from the measured DMs of the FRBs.

## 5. DISCUSSION

Here we discuss the locations, luminosity limits, and morphological features revealed by HST imaging in the context of other transient populations with known progenitors, and implications for FRB progenitors.



Figure 9. Limits on the near-IR luminosity at the FRB positions (lines) as a function of redshift, derived from GALFIT residual images after a Sérsic component is removed. Filled triangles represent  $5\sigma$  limits at the redshift of the putative, brighter FRB host galaxy. The limits rule out satellite galaxies at the FRB positions at the same redshift as the bright host to deep limits. If instead a background galaxy exists with a host luminosity similar to FRB 121102, this would require redshifts larger than allowed by the measured DM (diamonds, 95% confidence) in all cases except FRB 190714 and 191001.

## 5.1. The Locations of FRBs with Respect to their Host Galaxies

The high angular resolution of the HST imaging enables the determination of effective radii and the precise locations for FRB events with respect to their host galaxies using a variety of measures. In general, locations have been used in a variety of transient studies as a major diagnostic to uncovering their progenitors (Prieto et al. 2008; Fong et al. 2010; Fong & Berger 2013; Blanchard et al. 2016; Lunnan et al. 2017; De et al. 2020; Audcent-Ross et al. 2020; Schulze et al. 2020), as well as the relation to the distribution of young stars (UV,  $H\alpha$ light) and stellar mass (IR light) in their host galaxies. As a means to deciphering the origins of FRBs, we consider comparison data sets from transients which span a wide range of progenitor systems, from those which originate from massive stars in which the populations follow the UV light and exponential disk profiles of their hosts (LGRBs, CCSNe, SLSNe; Fruchter et al. 2006; Lunnan et al. 2015); to those with older stellar progenitors associated with compact objects which are weakly correlated with the UV light of their hosts (SGRBs, Type Ia SNe; Fong & Berger 2013; Wang et al. 2013; AudcentRoss et al. 2020); to those with unknown progenitors and larger offsets from their hosts (Ca-rich transients<sup>5</sup>).

Comparative studies based on ground-based observations of FRBs have found that the spatial distribution of FRBs are inconsistent with the distribution of LGRBs and SLSNe, both of which originate from strippedenvelope massive stars (Fruchter et al. 2006; Lunnan et al. 2015), but are consistent with other transient types (Bhandari et al. 2020a; Heintz et al. 2020). Our results support these studies, where we find that the locations of FRBs as a population are clearly more extended than LGRBs and SLSNe in terms of physical offsets, with a median of 3.2 kpc (68% interval of [2.0, 8.7] kpc). However, the host galaxies of FRBs are on average larger in physical size (and also stellar mass; c.f., Bhandari et al. 2020a; Heintz et al. 2020) than the hosts of LGRBs and SLSNe (Blanchard et al. 2016; Lunnan et al. 2015), with a range of sizes,  $r_e \approx 0.7 - 5.6$  kpc. Due to the larger FRB host galaxy sizes among the transient populations, the differences in offsets becomes less significant when normalized by the size of their host galaxies: the hostnormalized offsets of FRBs has a median of 1.4  $r_e$  (68%) interval of  $[0.6, 2.1]r_e$ , and are only statistically distinct from Ca-rich transients. Finally, FRBs appear to occur at slightly larger host-normalized offsets than expected given an exponential disk profile.

In terms of their host galaxy rest-frame optical and NIR light distributions, the FRBs are on moderately fainter regions of their host galaxies (median  $F_{\rm F}=0.33$ ). These distributions serve as proxies for the distribution of stellar mass in their galaxies, and older, moderately massive to low-mass stars, respectively. As a population, FRBs are once again statistically distinct from LGRBs and SLSNe which on average occur on the brighter regions of their host galaxies (Fruchter et al. 2006; Blanchard et al. 2016). The locations of FRBs are also distinct from SGRBs which are very weakly correlated with stellar mass, a consequence of their compact object progenitors which experience kicks and moderate delay times (Fong et al. 2010; Fong & Berger 2013). Notably, unlike SGRBs, no FRBs in our sample occur on the faintest regions of their host galaxies (tempered by the FRB localization errors and the small sample).

The locations of FRBs are consistent with the *radial* distribution of their host rest-frame optical light (fraction of enclosed light), and are indistinguishable from the locations of CCSNe in this regard (c.f., Audcent-Ross et al. 2020). Indeed, the fact that the locations of

FRBs trace the 1:1 distribution of the radial distribution of their host light, coupled with the fact that their local stellar mass surface densities are representative of their global host galaxy values, is indicative that their locations are consistent with the stellar mass within their host galaxies. We further find that while two FRBs have elevated local star formation rate densities compared to their global host values, as a population we do not find any clear correlation between FRBs and regions of elevated local star formation rate densities.

In general, the host galaxies of known repeating FRBs tend to have bluer colors, lower stellar masses, and higher star forming rates than those of apparent non-repeating FRBs (Bhandari et al. 2020a; Heintz et al. 2020). This is most saliently highlighted in the star-forming low-mass host galaxy of the repeating FRB 121102 (Chatterjee et al. 2017), which is an outlier in most host galaxy properties. Here, we find that, in terms of the IR distributions, the three known repeaters in this sample (FRBs 121102, 180916 and 190711) span the full range of offsets (physical and host-normalized) populated by apparent non-repeaters, as well as fractional flux and enclosed flux. While there do not appear to be any obvious trends in these properties between known repeaters and apparent non-repeaters, we caution that the sample sizes considered here are small.

Overall, the locations of FRBs support the picture that if there is one dominant progenitor population, that they do not originate from massive stars which are stripped of H and/or He (the progenitors of enginedriven SNe such as LGRBs and H-poor SLSNe). We further find that their locations are inconsistent with compact object progenitors which experienced kicks or long delay times transients from significantly older stellar progenitors (SGRBs, Ca-rich transients). However, given the size of the current sample, it is still possible that a fraction of FRBs originate from one of these alternative progenitor channels. These conclusions overall support previous results based on host stellar population properties that LGRB/SLSNe progenitors are not significant contributors (Heintz et al. 2020; Li et al. 2019; Bhandari et al. 2020a). Furthermore, we cannot differentiate the population of FRBs from CCSNe or Type Ia SNe based on their locations, although FRBs do not clearly trace either of these populations in every measured quantity. Therefore, we find it less plausible that the main progenitor channel of FRBs are compact object progenitors such as neutron star mergers or neutron star-black hole mergers, although progenitors which invoke white dwarfs (e.g., accretion-induced collapse of a WD to a NS), which are expected to resemble the prop-

<sup>&</sup>lt;sup>5</sup> These are also referred to as Ca-strong transients or CASTs; Shen et al. (2019)

erties of Type Ia SNe (Margalit et al. 2019) could still play a role.

## 5.2. The Association of FRB Locations with Spiral Arm Structure

In addition to precise location information, the deep HST imaging presented here also enhances low surface brightness features and morphological structure. In particular, previous HST studies of two galaxies, those of FRBs 190608 and 180916, demonstrate that they both exhibit complex spiral arm structure (Chittidi et al. 2020; Tendulkar et al. 2020). Spiral structure was also apparent in ground-based imaging for the host galaxy of FRB 191001, and supported by extended, continuum radio emission indicative of star formation (Bhandari et al. 2020b). Here, we find an *additional* two FRB hosts with clear spiral arm structure; those of FRBs 180924 and 190714 (Figure 7), and further uncover a bar feature in the host galaxy of FRB 180924. With the exception of FRB 180916, all FRB spiral-arm hosts are associated with apparent non-repeaters. The two remaining known repeaters in the sample are FRBs 121102 and 190711; the former originates from a low-luminosity host, while the latter originates in a host at the high redshift end of our sample. Thus we do not consider the non-detection of spiral features from these hosts to be constraining or informative.

Overall the prevalence of clear spiral structure (5/8,or  $\approx 60\%$  in our sample) is consistent with the observed galaxy population (Willett et al. 2013). Furthermore, despite the larger offsets of FRBs, we find that the locations of all well-localized FRBs with hosts that exhibit spiral structure are consistent with major spiral arm features. It is important to note that the IR light profile is dominated by red supergiants, AGB stars, and low-mass stars, as opposed to young, massive O- and B-stars seen in H $\alpha$  and UV imaging. In particular, in accordance with the density wave theory of spiral structure, the IR spiral arms generally spatially lag the H $\alpha$  light (Pour-Imani et al. 2016), although significant enhancement in star formation in the vicinity of the IR spiral arms is expected (Seigar & James 2002). The signal-to-noise of the FRB UV images prevent such a constraint for FRB hosts.

The locations of transients with respect to spiral arm features, as well as offsets from regions of peak brightness within the spiral arms, can serve as a major clue for their progenitors. In particular, the offset from bright peaks serves as a proxy of the spatial drift from birth to explosion site, and can set a timescale for the lifetime of the progenitor. Indeed, SNe exhibit distributions of offsets from the peak of their spiral arms in accordance with their progenitor age, with stripped-envelope SNe (Type Ib/c) having smaller offsets from the peak than Type II or Type Ia SNe (Aramyan et al. 2016). If all FRBs originated from young magnetars, it is expected that their positions would generally correlate with the UV spiral arms of their hosts, and at small offsets from star-forming features (c.f., the distribution of Galactic magnetars; Olausen & Kaspi 2014). However, we find that while the FRB positions are consistent with spiral features, they are not on the brightest part of the spiral arms. Indeed, UV and H $\alpha$  studies of the known repeating FRBs 121102 and 180916 found clear offsets from the closest star-forming features of  $\approx 250$  pc (Bassa et al. 2017; Tendulkar et al. 2020). This is also in agreement with the results by Chittidi et al. (2020), who found from detailed analysis of the UV imaging of FRB 190608 that the FRB did not prefer the most active star-forming region in the galaxy.

Taken together, this supports a picture that FRBs do not originate from the youngest, most massive stars, in concert with previous, comparative results with other transients (Li et al. 2019; Heintz et al. 2020; Bhandari et al. 2020a). We also find that FRBs do not appear to reside in the inner bulges of their host galaxies, which are generally dominated by older, higher-metallicity stars in comparison to the spiral arms (Peletier & Balcells 1996). It is further worth noting that the main selection effect at play in FRB discovery is the difficulty of detecting highly-scattered FRB signals, where the signal is temporally broadened by multipath propagation in a dense, turbulent medium. Since such sites are preferentially associated with star formation, one might naively expect there to be additional observational challenges in detecting FRBs in spiral arms where the chance of the FRB sightline intersecting an enhanced region of turbulence is higher. However, the precise effects of discovering FRBs with respect to morphological structure is not well-quantified.

## 5.3. Luminosity Limits on Alternative Host Candidates

Finally, we remark on the presence of fainter, alternative host galaxy candidates at the positions of the FRBs. This question is in part motivated by the low-luminosity host galaxy of FRB 121102 (Tendulkar et al. 2017), which, coupled with the remaining FRB hosts suggests a broad host galaxy luminosity function spanning the full range of galaxies (Heintz et al. 2020). Here we have explored the presence of both satellite galaxies at the same redshift as the putative host, and background galaxies at higher redshifts. The relatively low redshift range of the population examined here,  $0.03 \leq z \leq 0.522$ , enables deep constraints even in the presence of strong morphological features. In both scenarios, we find it unlikely that the FRBs originated from an underlying galaxy. The exceptions are FRBs 190714 and 191001: in the former case, a galaxy of equal luminosity to FRB 121102 would approach the redshift limit, while in the latter case, the high DM allows a host with  $\approx 8$  times the luminosity of FRB 121102, albeit still on the faint-end slope of the galaxy luminosity function ( $\approx 8 \times 10^8 L_{\odot}$ ).

### 6. SUMMARY & CONCLUSIONS

In this paper we used high-resolution HST imaging to perform a detailed study on the locations of 8 FRBs and their environments, 6 of which are newly presented here. We used these data to place constraints on the spatial distributions (physical and host-normalized), in support of previous works based on ground-based imaging. We find a median host-normalized offset of  $1.4([0.6, 2.1]r_e; 68\%$  interval), and overall a distribution that lies between the more centrally concentrated LGRBs and SLSNe, and the extended SGRBs and Ca-rich transients. We also determine the distribution of FRBs with respect to their IR (rest-frame optical) host galaxy light (fractional flux, and radial distribution), showing that FRBs are consistent with tracing the stellar mass distribution of their host galaxies.

The sensitivity of HST additionally enables constraints on possible alternative host galaxy candidates; we find it improbable that there exists a satellite or background galaxy at the FRB locations, strengthening the associations with the brighter, putative hosts identified in ground-based imaging for this sample. We explore the FRB site properties in terms of star formation rate (near-UV) and stellar mass (IR) surface densities, finding that the locations are not particularly enhanced in either property compared to the global values of their hosts (although few measurements exist for the star formation rate densities). Finally, we find that 5/8 FRB host galaxies in the sample have spiral arm features, and that these FRBs are consistent with the locations of those spiral arms (albeit inconsistent with locations on the *brightest* peaks of these spiral arms). If there is a dominant progenitor population among this tested distribution, we thus do not find strong support for a connection to the most massive (stripped-envelope) stars, or events which require kicks and long delay times.

The promise of sub-arcsecond localized FRBs in solving the progenitor question is being realized, in part, with the first population studies of their local and host galaxy environments. Such precisely-localized FRBs are and will continue to be detected at growing rates. As the number of secure associations continues to increase, we will be able to make significant progress toward understanding their progenitors, as well as connecting the properties of FRBs to those of their host galaxies. Larger sample sizes of FRBs with high-resolution imaging, enabled by space-based facilities such as HST and JWST will also be crucial in understanding if there is a single, dominant progenitor channel or multiple contributing channels.

Software: Photutils (Bradley et al. 2020); GALFIT (Peng et al. 2002); Source Extractor (Bertin & Arnouts 1996); IRAF (Tody 1986); NumPy (Oliphant 2006); Astropy (Price-Whelan et al. 2018); Matplotlib (Hunter 2007); SciPy (Virtanen et al. 2020).

## Facilities: HST:WFC3

## ACKNOWLEDGEMENTS

We acknowledge Nia Imara and Ben Margalit for helpful discussions. A.G.M. acknowledges support by the National Science Foundation Graduate Research Fellowship under Grant No. 1842400. The Fast and Fortunate for FRB Follow-up team acknowledges support from NSF grants AST-1911140 and AST-1910471. W.F. acknowledges support by the National Science Foundation under grant Nos. AST-1814782 and AST-190935. N.T. acknowledges support by FONDECYT grant 11191217. C.K.D. acknowledges the support of the CSIRO Postgraduate Scholarship - Astronomy and Space (47417). A.T.D. is the recipient of an ARC Future Fellowship (FT150100415). This research is based on observations made with the NASA/ESA Hubble Space Telescope obtained from the Space Telescope Science Institute, which is operated by the Association of Universities for Research in Astronomy, Inc., under NASA contract NAS 5–26555. These observations are associated with programs #15878, 16080, 14890 and 16072.

## REFERENCES

Aggarwal, K., Budavari, T., Deller, A., et al. 2021, ApJ,

submitted, xx

Anderson, J. 2016, Supplemental Dither Patterns for WFC3/IR, Space Telescope WFC Instrument Science Report

- Anderson, J., MacKenty, J., Baggett, S., & Noeske, K. 2012, The Efficacy of Post-Flashing for Mitigating CTE-Losses in WFC3/UVIS Images, Tech. rep., Baltimore, MD: STScI
- Aramyan, L. S., Hakobyan, A. A., Petrosian, A. R., et al. 2016, MNRAS, 459, 3130, doi: 10.1093/mnras/stw873
- Audcent-Ross, F. M., Meurer, G. R., Audcent, J. R., et al. 2020, MNRAS, 492, 848, doi: 10.1093/mnras/stz3282
- Avila, R. J., Hack, W., Cara, M., et al. 2015, in Astronomical Society of the Pacific Conference Series, Vol. 495, Astronomical Data Analysis Software an Systems XXIV (ADASS XXIV), ed. A. R. Taylor & E. Rosolowsky, 281. https://arxiv.org/abs/1411.5605
- Bannister, K. W., Deller, A. T., Phillips, C., et al. 2019, Science, 365, 565, doi: 10.1126/science.aaw5903
- Bassa, C. G., Tendulkar, S. P., Adams, E. A. K., et al. 2017, ApJL, 843, L8, doi: 10.3847/2041-8213/aa7a0c
- Bertin, E., & Arnouts, S. 1996, A&AS, 117, 393
- Bhandari, S., Sadler, E. M., Prochaska, J. X., et al. 2020a, ApJL, 895, L37, doi: 10.3847/2041-8213/ab672e
- Bhandari, S., Bannister, K. W., Lenc, E., et al. 2020b, ApJL, 901, L20, doi: 10.3847/2041-8213/abb462
- Blanchard, P. K., Berger, E., & Fong, W.-f. 2016, ApJ, 817, 144, doi: 10.3847/0004-637X/817/2/144
- Bochenek, C. D., Ravi, V., Belov, K. V., et al. 2020, Nature, 587, 59, doi: 10.1038/s41586-020-2872-x
- Bradley, L., Sipőcz, B., Robitaille, T., et al. 2020, astropy/photutils: 1.0.0, 1.0.0, Zenodo, doi: 10.5281/zenodo.4044744
- Chatterjee, S., Law, C. J., Wharton, R. S., et al. 2017, Nature, 541, 58, doi: 10.1038/nature20797
- CHIME/FRB Collaboration, Amiri, M., Bandura, K., et al. 2018, ApJ, 863, 48, doi: 10.3847/1538-4357/aad188
- CHIME/FRB Collaboration, Andersen, B. C., Bandura,
  K. M., et al. 2020, Nature, 587, 54,
  doi: 10.1038/s41586-020-2863-y
- Chittidi, J. S., Simha, S., Mannings, A., et al. 2020, arXiv e-prints, arXiv:2005.13158. https://arxiv.org/abs/2005.13158
- Cordes, J. M., & Chatterjee, S. 2019, ARA&A, 57, 417, doi: 10.1146/annurev-astro-091918-104501
- Day, C. K., Deller, A. T., Shannon, R. M., et al. 2020, MNRAS, in prep.
- De, K., Kasliwal, M. M., Tzanidakis, A., et al. 2020, arXiv e-prints, arXiv:2004.09029.
  - https://arxiv.org/abs/2004.09029
- Eftekhari, T., & Berger, E. 2017, ApJ, 849, 162, doi: 10.3847/1538-4357/aa90b9

- Eftekhari, T., Berger, E., Margalit, B., Metzger, B. D., &
  Williams, P. K. G. 2020, ApJ, 895, 98,
  doi: 10.3847/1538-4357/ab9015
- Evans, Neal J., I., Dunham, M. M., Jørgensen, J. K., et al. 2009, ApJS, 181, 321, doi: 10.1088/0067-0049/181/2/321
- Fong, W., & Berger, E. 2013, ApJ, 776, 18, doi: 10.1088/0004-637X/776/1/18
- Fong, W., Berger, E., & Fox, D. B. 2010, ApJ, 708, 9, doi: 10.1088/0004-637X/708/1/9
- Fruchter, A. S., Levan, A. J., Strolger, L., et al. 2006, Nature, 441, 463, doi: 10.1038/nature04787
- Heintz, K. E., Prochaska, J. X., Simha, S., et al. 2020, arXiv e-prints, arXiv:2009.10747. https://arxiv.org/abs/2009.10747
- Hunter, J. D. 2007, Computing in Science & Engineering, 9, 90, doi: 10.1109/MCSE.2007.55
- Kennicutt, Robert C., J. 1998, ARA&A, 36, 189, doi: 10.1146/annurev.astro.36.1.189
- Kennicutt, R. C., & Evans, N. J. 2012, ARA&A, 50, 531, doi: 10.1146/annurev-astro-081811-125610
- Kumar, P., Shannon, R. M., Flynn, C., et al. 2020, arXiv e-prints, arXiv:2009.01214. https://arxiv.org/abs/2009.01214
- Lada, C. J., & Dame, T. M. 2020, ApJ, 898, 3, doi: 10.3847/1538-4357/ab9bfb
- Law, C. J., Butler, B. J., Prochaska, J. X., et al. 2020, ApJ, 899, 161, doi: 10.3847/1538-4357/aba4ac
- Li, Y., Zhang, B., Nagamine, K., & Shi, J. 2019, ApJL, 884, L26, doi: 10.3847/2041-8213/ab3e41
- Lorimer, D. R., Bailes, M., McLaughlin, M. A., Narkevic, D. J., & Crawford, F. 2007, Science, 318, 777, doi: 10.1126/science.1147532
- Lunnan, R., Chornock, R., Berger, E., et al. 2015, ApJ, 804, 90, doi: 10.1088/0004-637X/804/2/90
- Lunnan, R., Kasliwal, M. M., Cao, Y., et al. 2017, ApJ, 836, 60, doi: 10.3847/1538-4357/836/1/60
- Macquart, J.-P., Bailes, M., Bhat, N. D. R., et al. 2010, Publications of the Astronomical Society of Australia, 27, 272, doi: 10.1071/AS09082
- Macquart, J. P., Prochaska, J. X., McQuinn, M., et al. 2020, Nature, 581, 391, doi: 10.1038/s41586-020-2300-2
- Marcote, B., Paragi, Z., Hessels, J. W. T., et al. 2017, ApJL, 834, L8, doi: 10.3847/2041-8213/834/2/L8
- Marcote, B., Nimmo, K., Hessels, J. W. T., et al. 2020, Nature, 577, 190, doi: 10.1038/s41586-019-1866-z
- Margalit, B., Berger, E., & Metzger, B. D. 2019, ApJ, 886, 110, doi: 10.3847/1538-4357/ab4c31
- Margalit, B., & Metzger, B. D. 2018, ApJL, 868, L4, doi: 10.3847/2041-8213/aaedad

- Nicholl, M., Williams, P. K. G., Berger, E., et al. 2017, ApJ, 843, 84, doi: 10.3847/1538-4357/aa794d
- Olausen, S. A., & Kaspi, V. M. 2014, ApJS, 212, 6, doi: 10.1088/0067-0049/212/1/6
- Oliphant, T. E. 2006, A guide to NumPy, Vol. 1 (Trelgol Publishing USA)
- Palmerio, J. T., Vergani, S. D., Salvaterra, R., et al. 2019, A&A, 623, A26, doi: 10.1051/0004-6361/201834179
- Peletier, R. F., & Balcells, M. 1996, AJ, 111, 2238, doi: 10.1086/117958
- Peng, C. Y., Ho, L. C., Impey, C. D., & Rix, H.-W. 2002, AJ, 124, 266, doi: 10.1086/340952
- Perley, D. A., Quimby, R. M., Yan, L., et al. 2016, ApJ, 830, 13, doi: 10.3847/0004-637X/830/1/13
- Petroff, E., Hessels, J. W. T., & Lorimer, D. R. 2019, A&A Rv, 27, 4, doi: 10.1007/s00159-019-0116-6
- Planck Collaboration, Ade, P. A. R., Aghanim, N., et al. 2016, A&A, 594, A13, doi: 10.1051/0004-6361/201525830
- Platts, E., Weltman, A., Walters, A., et al. 2019, PhR, 821, 1, doi: 10.1016/j.physrep.2019.06.003
- Pour-Imani, H., Kennefick, D., Kennefick, J., et al. 2016, ApJL, 827, L2, doi: 10.3847/2041-8205/827/1/L2
- Price-Whelan, A. M., Sipőcz, B. M., Günther, H. M., et al. 2018, AJ, 156, 123, doi: 10.3847/1538-3881/aabc4f
- Prieto, J. L., Stanek, K. Z., & Beacom, J. F. 2008, ApJ, 673, 999, doi: 10.1086/524654
- Rafelski, M., Teplitz, H. I., Gardner, J. P., et al. 2015, AJ, 150, 31, doi: 10.1088/0004-6256/150/1/31
- Safarzadeh, M., Prochaska, J. X., Heintz, K. E., & Fong, W.-f. 2020, arXiv e-prints, arXiv:2009.11735. https://arxiv.org/abs/2009.11735

- Schulze, S., Yaron, O., Sollerman, J., et al. 2020, arXiv e-prints, arXiv:2008.05988. https://arxiv.org/abs/2008.05988
- Seigar, M. S., & James, P. A. 2002, MNRAS, 337, 1113, doi: 10.1046/j.1365-8711.2002.06007.x
- Shen, K. J., Quataert, E., & Pakmor, R. 2019, ApJ, 887, 180, doi: 10.3847/1538-4357/ab5370
- Spitler, L. G., Scholz, P., Hessels, J. W. T., et al. 2016, Nature, 531, 202, doi: 10.1038/nature17168
- Svensson, K. M., Levan, A. J., Tanvir, N. R., Fruchter, A. S., & Strolger, L. G. 2010, MNRAS, 405, 57, doi: 10.1111/j.1365-2966.2010.16442.x
- Tendulkar, S. P., Bassa, C. G., Cordes, J. M., et al. 2017, ApJL, 834, L7, doi: 10.3847/2041-8213/834/2/L7
- Tendulkar, S. P., Gil de Paz, A., Kirichenko, A. Y., et al. 2020, arXiv e-prints, arXiv:2011.03257. https://arxiv.org/abs/2011.03257
- Thornton, D., Stappers, B., Bailes, M., et al. 2013, Science, 341, 53, doi: 10.1126/science.1236789
- Tody, D. 1986, in Society of Photo-Optical Instrumentation Engineers (SPIE) Conference Series, Vol. 627, Instrumentation in astronomy VI, ed. D. L. Crawford, 733, doi: 10.1117/12.968154
- Uddin, S. A., Burns, C. R., Phillips, M. M., et al. 2020, The Astrophysical Journal, 901, 143, doi: 10.3847/1538-4357/abafb7
- Virtanen, P., Gommers, R., Oliphant, T. E., et al. 2020, Nature Methods, 17, 261, doi: https://doi.org/10.1038/s41592-019-0686-2
- Wang, X., Wang, L., Filippenko, A. V., Zhang, T., & Zhao, X. 2013, Science, 340, 170, doi: 10.1126/science.1231502
- Willett, K. W., Lintott, C. J., Bamford, S. P., et al. 2013, MNRAS, 435, 2835, doi: 10.1093/mnras/stt1458
- Woosley, S. E. 1993, ApJ, 405, 273, doi: 10.1086/172359