

## THE ACS VIRGO CLUSTER SURVEY. XV. THE FORMATION EFFICIENCIES OF GLOBULAR CLUSTERS IN EARLY-TYPE GALAXIES: THE EFFECTS OF MASS AND ENVIRONMENT<sup>1</sup>

ERIC W. PENG<sup>2,3</sup>

Herzberg Institute of Astrophysics, National Research Council of Canada, Victoria,  
BC V9E 2E7, Canada; eric.peng@nrc-cnrc.gc.ca

ANDRÉS JORDÁN<sup>4,5,6</sup>

Harvard-Smithsonian Center for Astrophysics, Cambridge, MA 02138

PATRICK CÔTÉ

Herzberg Institute of Astrophysics, National Research Council of Canada, Victoria, BC V9E 2E7, Canada

MARIANNE TAKAMIYA

Department of Physics and Astronomy, University of Hawaii, Hilo, HI 96720

MICHAEL J. WEST<sup>7,8</sup>

European Southern Observatory, Alonso de Cordova 3107, Vitacura, Santiago, Chile

JOHN P. BLAKESLEE<sup>9</sup>

Herzberg Institute of Astrophysics, National Research Council of Canada,  
Victoria, BC V9E 2E7, Canada

CHIN-WEI CHEN<sup>10</sup>

Institute for Astronomy, National Central University Taiwan, Chung-Li 32054, Taiwan

LAURA FERRARESE

Herzberg Institute of Astrophysics, National Research Council of Canada, Victoria, BC V9E 2E7, Canada

SIMONA MEI<sup>11</sup>

GEPI, Observatoire de Paris, Section de Meudon, 92195 Meudon Cedex, France

JOHN L. TONRY

Institute for Astronomy, University of Hawaii, Honolulu, HI 96822

AND

ANDREW A. WEST

Department of Astronomy, University of California, Berkeley, CA 94720

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### ABSTRACT

The fraction of stellar mass contained in globular clusters (GCs), also measured by number as the specific frequency, is a fundamental quantity that reflects both a galaxy's early star formation and its entire merging history. We present specific frequencies, luminosities, and mass fractions for the globular cluster systems of 100 early-type galaxies in the ACS Virgo Cluster Survey, the largest homogeneous catalog of its kind. We find the following: (1) GC mass fractions can be high in both giants and dwarfs but are universally low in galaxies with intermediate luminosities. (2) The behavior of specific frequency across galaxy mass is dominated by the blue GCs. (3) GC fractions of low-mass galaxies exhibit a dependence on environment. Nearly all dwarf galaxies with high GC fractions are within 1 Mpc of the cD galaxy M87, presenting the first strong evidence that GC formation in dwarfs is biased toward dense environments. (4) GC formation in central dwarfs is biased because their stars form earliest and most intensely. Comparisons to the Millennium Simulation show that central dwarfs have older stellar populations and form more stars at higher star formation rates (SFRs) and SFR surface densities. The SFR surface density in simulated dwarfs peaks before the total SFR, naturally producing GC populations that are older and more metal-poor than the field stars. (5) Dwarfs within  $\sim 40$  kpc of the giant ellipticals M87 and M49 are red and have few or no GCs, suggesting that they have been tidally stripped and

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<sup>2</sup> Space Telescope Science Institute, Baltimore, MD 21228; peng@stsci.edu.

<sup>3</sup> Department of Astronomy, Peking University, Beijing 100871, China.

<sup>4</sup> Clay Fellow.

<sup>5</sup> European Southern Observatory, 85748 Garching bei München, Germany.

<sup>6</sup> Departamento de Astronomía y Astrofísica, Pontificia Universidad Católica de Chile, Casilla 306, Santiago 22, Chile.

<sup>7</sup> Gemini Observatory, Casilla 603, La Serena, Chile.

<sup>8</sup> Department of Physics and Astronomy, University of Hawaii, Hilo, HI 96720.

<sup>9</sup> Department of Physics and Astronomy, Washington State University, Pullman, WA 99164-2814.

<sup>10</sup> Herzberg Institute of Astrophysics, National Research Council of Canada, Victoria, BC V9E 2E7, Canada.

<sup>11</sup> University of Paris 7 Denis Diderot, 75205 Paris Cedex 13, France.

have contributed their GCs to the halos of their giant neighbors. The central dwarfs with high GC mass fractions are thus likely to be the survivors most similar to the protogalaxies that assembled the rich M87 globular cluster system.

*Subject headings:* galaxies: dwarf — galaxies: elliptical and lenticular, cD — galaxies: evolution — galaxies: halos — galaxies: star clusters — globular clusters: general

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## 1. INTRODUCTION

Globular clusters (GCs) constitute a small fraction of the stellar mass in galaxies, but their ubiquity, relative simplicity, and old ages make them the most prominent representatives of a bygone epoch of galaxy formation. GCs are made of stars that are among the oldest in galaxies, and they can be observed at large distances (e.g., Blakeslee et al. 2003b). These old star clusters are thus unique, both intrinsically and observationally, for understanding the early, intense star-forming episodes that mark galaxy formation.

In the local universe, we see massive star clusters forming wherever there are high star formation rate (SFR) surface densities (Larsen & Richtler 2000), providing a connection that suggests that the properties of star cluster populations (age, metallicity, mass) should scale quite closely with field stars formed in the same events. The properties of globular cluster systems do in fact correlate strongly with the properties of the field stars of their host galaxies. The mean metallicities of GC systems have long been known to scale with the metallicity of their host (van den Bergh 1975; Brodie & Huchra 1991), and the mean metallicities of both the metal-rich and metal-poor subpopulations also correlate with the luminosity and mass of the host galaxy (Larsen et al. 2001; Peng et al. 2006a, hereafter Paper IX, and references therein). However, if GC systems directly followed the underlying field light in every way, they might be less interesting. For instance, although the metallicities of GC systems may track those of galaxies, they are consistently offset to lower values by  $\sim 0.5$ – $0.8$  dex in  $[\text{Fe}/\text{H}]$  (e.g., Jordán et al. 2004a; Lotz et al. 2004). Most conspicuously, even the most massive and metal-rich galaxies have GC systems dominated by metal-poor star clusters ( $[\text{Fe}/\text{H}] \lesssim -1$ ). This suggests a disconnect between the formation of “halo” stellar populations and the bulk of the galaxy.

One of the most studied aspects of this GC-galaxy duality concerns the specific frequency of globular clusters, or the number of GCs per unit stellar luminosity. Specific frequency,  $S_N$ , was introduced by Harris & van den Bergh (1981) and is defined as the number of GCs normalized to a galaxy luminosity of  $M_V = -15$ . The purpose of studying  $S_N$  across galaxies of different masses, morphologies, and environments is, in the words of that initial paper, “to investigate whether there is in fact a ‘universal’ and uniform capability for globular cluster formation.” This simple quantity and similar ones related to it turn out to be extremely interesting galaxy diagnostics. It appears true that for galaxies above a certain mass there is a nearly universal capability to form globular clusters, but this process is not uniform across all galaxies, at least as seen in comparison to the field stars. Specific frequencies of spiral galaxies like the Milky Way are generally 0.5–1 (Goudfrooij et al. 2003; Rhode & Zepf 2004; Chandar et al. 2004), although they have a mean of  $\sim 4$  when normalized only to bulge luminosity (Côté et al. 2000). The specific frequencies of massive ellipticals are 2–6, and those of some cD galaxies such as M87 (NGC 4486) can be well in excess of 10. These trends are apparent even when the number of GCs is normalized to stellar mass as opposed to stellar luminosity (Rhode et al. 2005). Dwarf elliptical galaxies (dEs), whose GC systems are predominantly metal-poor, can also have high  $S_N$  similar to those of giant ellip-

ticals (Durrell et al. 1996a; Miller et al. 1998; Lotz et al. 2004; Miller & Lotz 2007; Puzia & Sharina 2008), as can some dwarf irregulars (Seth et al. 2004), suggesting the possibility that the halos of large galaxies were formed mainly through the accretion of dwarflike objects (Searle & Zinn 1978; Côté et al. 1998, 2000).

A central question in the study of GC systems is, how do we understand different GC fractions in the context of galaxy assembly? The formation of globular cluster systems has been particularly tied to the formation of massive elliptical galaxies, in which GCs are often present in large numbers and where GCs are most easily observed. The mergers and accretion events expected during the hierarchical assembly of these galaxies must also be able to form their GC systems. Observations and simulations of elliptical galaxy formation are creating a picture in which the stars form early and quickly (e.g., Kodama et al. 1998), mimicking the traditional “monolithic collapse” scenario, but where the assembly of these stars into a single galaxy continues until late times through largely dissipationless mergers (De Lucia et al. 2006; De Lucia & Blaizot 2007) and star formation at late times is suppressed by energy feedback (Springel et al. 2005; Croton et al. 2006).

Intertwined with the issue of galaxy formation is that of the formation efficiency of the GCs themselves: why do globular clusters form with different efficiencies with respect to their light in different galaxies? Blakeslee et al. (1997) and Blakeslee (1999) studied the GC systems of brightest cluster galaxies (BCGs) in Abell galaxy clusters and found that the number of GCs scaled with the velocity dispersion of the galaxy cluster rather than with the luminosity of the BCGs, suggesting that GC formation is closely linked to the total mass of the system, i.e.,  $S_N \propto M/L$ . In a similar vein, McLaughlin (1999a) examined the high  $S_N$  in M87 and found that the large number of GCs it possessed was not anomalous when normalized to the total baryonic mass (including the hot X-ray gas) rather than just to the stellar mass. McLaughlin (1999a) defined a “universal” GC formation efficiency of  $\hat{\epsilon} = 0.26\%$ , where  $\hat{\epsilon}$  is the fraction of the baryonic mass that ends up in globular clusters. Kravtsov & Gnedin (2005) studied the formation of the GC system in a high-resolution hydrodynamic simulation and also found that the mass in GCs was directly proportional to the total halo mass of the galaxy. If total mass drives GC formation, then it is the variation in converting baryons into field stars that drives trends in specific frequency.

The connection between GCs and galactic mass (baryonic or total) makes it tempting to try to explain them using simulations of dark matter and galaxy assembly. Beasley et al. (2002) simulated the color distributions of GCs using semianalytic models of galaxy formation, while the aforementioned simulations of Kravtsov & Gnedin (2005) formed a Milky Way mass galaxy in detail. Moore et al. (2006) identified metal-poor GCs with early collapsing dark matter peaks in cosmological  $N$ -body simulations. An empirical connection between the total mass of a system and the mass contained in its globular clusters does not, however, explain why the star formation histories of the GCs and the field should be different.

One clue is that in the local universe, high mass fractions in massive star clusters occur in regions of high star formation surface density (Larsen & Richtler 2000). Thus, one explanation for high- $S_N$  galaxies might be that they formed more of their stars in

high-efficiency events. Given that the star formation rate in galaxies was most intense at very early times ( $z \gtrsim 2$ ), variations in  $S_N$  could result from different times of formation, especially when coupled with a sharp star formation cutoff at reionization (Santos 2003). In this scenario, with  $S_N$  a function of formation time, GC formation is biased toward the earliest collapsing halos that can create a large fraction of their stars at high efficiency before reionization. Low-mass halos in dense environments collapse earlier (Gao et al. 2005; Diemand et al. 2005), and these are also the ones most susceptible to heating through photoionization and feedback. One expectation of this scenario is that  $S_N$  will be “biased” toward dense environments (West 1993), especially in low-mass galaxies. The simulations of Moore et al. (2006) suggested that GCs and satellites would be more highly clustered in dense regions for higher redshifts of reionization, and the same idea can be applied to  $S_N$  at fixed reionization redshift.

Previous observational studies have shown some evidence of an environmental dependence for  $S_N$ . West (1993) showed that the mean  $S_N$  of elliptical galaxies correlates with the local galaxy density. Blakeslee (1997, 1999) found that  $S_N$  in BCGs scaled with properties that reflect cluster density (cluster velocity dispersion, X-ray temperature, and X-ray luminosity), and that cluster galaxies closer to their cluster’s X-ray center have higher  $S_N$ . West et al. (1995) also found a correlation between the  $S_N$  of BCGs and the X-ray luminosity of the cluster, interpreting it as evidence for a population of intergalactic globular clusters. On the other hand, Lotz et al. (2004) and Miller & Lotz (2007) did not find an obvious correlation between  $S_N$  and clustercentric radius for dEs in the Virgo and Fornax Clusters, although the number of GCs in their galaxies was small and  $S_N$  errors were quite large. Also, Spitler et al. (2008) recently found that the relatively isolated elliptical NGC 821 has an  $S_N$  comparable to cluster galaxies of like luminosities.

The specific frequencies (or formation efficiencies) of globular cluster systems are clearly a fundamental property of galaxies. However, accurately measuring  $S_N$  is traditionally fraught with uncertainty despite the fact that at its most basic level it requires only simple counting of GCs. In practice, observations need to be deep enough to observe past the mean of the GC luminosity function (GCLF), and GC selection needs to be efficient enough so that contaminants do not overwhelm the GCs in lower mass galaxies. The total magnitude and distance of the galaxy also need to be known to establish its luminosity (or mass). Studies of larger galaxies also benefit from wide-field coverage so that enough of the GC system is sampled to minimize extrapolation errors. It is because of these limitations that homogeneous surveys of  $S_N$  across a wide range of galaxy mass are difficult to conduct. To put together a complete picture of globular cluster system and galaxy formation, it is important to study galaxies at all masses in the same way: at low masses we study the survivors of hierarchical assembly, and at high masses we study its final products.

With this in mind, we have undertaken a careful study of the formation efficiencies for the GC systems of 100 early-type galaxies in the ACS Virgo Cluster Survey (ACSVCS; Côté et al. 2004, hereafter Paper I). The deep, high-resolution, relatively wide-field imaging provides the most complete census to date of GCs in early-type galaxies over a large range in galaxy luminosity ( $-22 < M_B < -15$ ). In this paper we present the specific frequencies and other related quantities of the ACSVCS galaxies and quantify their trends as a function of host galaxy properties.

## 2. OBSERVATIONS AND DATA

The ACS Virgo Cluster Survey is a large program to image 100 early-type galaxies in the Virgo Cluster with the *Hubble Space Telescope* (*HST*) Advanced Camera for Surveys (ACS; Ford et al.

1998). The survey is described in detail in Paper I, and the data reduction techniques are outlined in Jordán et al. (2004b, hereafter Paper III), but we briefly summarize them here. We obtained images in the F475W ( $g$ ) and F850LP ( $z$ ) filters of galaxies selected to be early type and with confirmed cluster membership in the Virgo Cluster Catalog (VCC) of Binggeli et al. (1985). Our sample excludes the southern extension, has  $B_T < 16$ , and the galaxies are morphologically classified as E, S0, dE, dE,N, dS0, or dS0,N. This early-type galaxy sample is magnitude limited for the brightest 26 galaxies ( $B_T < 12.15$  or  $M_B < -19.10$ ) and contains a representative sample of galaxies for  $M_B < -15.11$ . The color distributions of the GC systems were presented in Paper IX, and their size distributions were presented in Jordán et al. (2005, hereafter Paper X). We have measured distances for 84 galaxies from the method of surface brightness fluctuations (SBFs; Mei et al. 2005a, 2005b, 2007, hereafter Papers V, IV, and XIII, respectively). For galaxies where an SBF distance could not be measured, we adopt a distance to the Virgo Cluster of  $D = 16.5$  Mpc with a distance modulus of  $31.09 \pm 0.03$  mag from Tonry et al. (2001), corrected by the final results of the Key Project distances (Freedman et al. 2001; see also discussion in Paper V). The surface brightness profiles, total magnitudes, and colors of the sample galaxies in  $g$  and  $z$  were described in Ferrarese et al. (2006b, hereafter Paper VI), and the properties of the GCLFs for 89 of our sample galaxies were presented in Jordán et al. (2006) and Jordán et al. (2007, hereafter Paper XII). The galaxy nuclear properties were presented in Côté et al. (2006, hereafter Paper VIII), and their connection to supermassive black holes was presented in Ferrarese et al. (2006a). These data have also been analyzed for ultracompact dwarf galaxies (Haşegan et al. 2005, hereafter Paper VII), diffuse star clusters (Peng et al. 2006b, hereafter Paper XI), the connection between GCs and low-mass X-ray binaries (Jordán et al. 2004c, hereafter Paper II; Sivakoff et al. 2007), and color-magnitude relations in GC systems (Mieske et al. 2006). Together, these data create the best opportunity to date to study the formation efficiencies of GC systems in early-type galaxies, and this paper refers often to the quantities measured in the preceding papers.

### 2.1. Data Reduction and Control Fields

Each galaxy was imaged with the Wide Field Channel (WFC) of the ACS. We reduced the ACS WFC images using a dedicated pipeline described in Paper III (see also Blakeslee et al. 2003a). We produced the science images by combining and cleaning them of cosmic rays using the Pyraf routine `multidrizzle` (Koekemoer et al. 2002). We then subtracted a model of the galaxy light and used the source detection program SExtractor (Bertin & Arnouts 1996) to detect and mask sources and remove residual background. Our final object detection includes estimates of both the image noise and noise due to SBFs (ignoring the latter results in many false detections in the bright central regions of the galaxy), and objects are only included in the final catalog if they are detected in both filters. After rejecting very bright or elongated objects to eliminate obvious foreground stars and background galaxies, and passing our catalog through a generous color cut, we use the program KINGPHOT (Paper X) to measure magnitudes and King model parameters. KINGPHOT fits King (1966) model surface brightness profiles convolved with the filter- and spatially dependent point-spread function (PSF). Magnitudes and colors are corrected for foreground extinction using the reddening maps of Schlegel et al. (1998) and extinction ratios for a G2 star (Paper II; Sirianni et al. 2005). For the purposes of this paper, whenever we refer to  $g$  and  $z$ , we mean the *HST* ACS magnitudes  $g_{475}$  and  $z_{850}$ . Magnitudes in the Sloan Digital Sky Survey (SDSS) system will

be designated explicitly, e.g., as in  $g_{SDSS}$ . The low redshifts of the galaxies being studied ( $-575$  to  $2284 \text{ km s}^{-1}$ ) mean that  $K$ -corrections are negligible and thus we do not  $K$ -correct any of the magnitudes presented in this paper.

In any study of extragalactic star clusters, it is important to quantify and correct for background contamination. To that end, we have also reduced 16 blank, high-latitude control fields taken from the ACS Pure Parallel Program (GO-9488, PI: Ratnatunga; and GO-9575, PI: Sparks). For each galaxy, we have “customized” the control sample to mimic the spatially varying detection efficiency that is a function of the surface brightness of the unresolved galaxy light. Details of this procedure are given in Papers IX and XI.

We select probable globular clusters using their sizes and magnitudes. We can assign a probability that any object is a GC based on position in the  $r_h$ - $z$  diagram and the locus of contaminants in the same plane. Objects with probabilities greater than 0.5 are included in our GC sample, although the exact choice of the cutoff value does not significantly affect our results. We describe the details of this selection in Paper IX and A. Jordán et al. (2008, in preparation).

## 2.2. WFPC2 Parallel Fields

In addition to ACS WFC imaging of our target galaxies, we have also acquired parallel imaging of 100 “blank” fields using the Wide Field Planetary Camera 2 (WFPC2) in the F606W (wide  $V$ ) and F814W ( $I$ ) filters. The separation between ACS and WFPC2 in the *HST* focal plane is approximately  $5.8'$ , which is 29 kpc at the mean distance to the Virgo Cluster (16.5 Mpc). For most of the ACSVCS galaxies, this distance puts the WFPC2 field well outside of the target galaxy, and thus one motivation for taking them is to search for intergalactic GCs (e.g., West et al. 1995; Jordán et al. 2003; Williams et al. 2007). In this paper we use the WFPC2 parallel images instead to constrain the total number of GCs in the larger, more luminous galaxies, where the GC system is still detectable 29 kpc from the center. The WFPC2 images were reduced using a modified form of the PyRAF pipeline written by Alasdair Allan.<sup>12</sup> We implemented the cosmic-ray cleaning algorithm LACOSMIC (van Dokkum 2001) and a geometric distortion correction (Bagget et al. 2002) in the PyRAF pipeline. Sources were selected from a SExtractor catalog having magnitudes and colors of GCs. The sizes and magnitudes of GC candidates were also measured with KINGPHOT. To account for sample contamination, we selected 10 blank WFPC2 fields from the *HST* archive of similar depth and observed with the same filters as our WFPC2 data. We analyzed them in the same manner as our data, using WFPC2 PSFs provided by P. B. Stetson. We find that there are on average  $5.8 \pm 2.5$  contaminant sources per WFPC2 field.

## 3. CALCULATING SPECIFIC FREQUENCIES, LUMINOSITIES, AND MASSES

### 3.1. Total Numbers of GCs

The determination of specific frequency requires knowledge of the total numbers of GCs, the total apparent magnitude of the galaxy (traditionally in the  $V$  band), and the distance. For the distances, we use the SBF-determined distances presented in Paper XIII, using their polynomial calibration. The measurement of the total number of GCs is in principle a simple task of counting but is not quite so straightforward in practice. Our images are deep enough that we generally detect the brightest  $\sim 90\%$  of the

GCLF, limiting the uncertainties from extrapolations in luminosity. Nevertheless, a full count of the GCs within the ACS field of view involves knowing our level of completeness, which is a function of the GC magnitude, size, and the local background flux. We also need to know the form of the GCLF so that we can estimate the number of GCs missed.

Fortunately, both the level of completeness and the GCLF can be determined for our data. We empirically derived the completeness of our data as a function of magnitude, background flux, and GC size, using simulations of nearly 5 million GCs placed in actual ACSVCS images for galaxies of different surface brightnesses. Using these data, we can nonparametrically estimate the expected detection probability of any GC in the survey. In addition, Paper XII was able to determine the form of the GCLF for 89 galaxies in the survey. For galaxies where we were unable to measure the form of the GCLF, we assume the mean to be the same as that of the cD galaxy VCC 1316 and a Gaussian distribution with sigma derived from equation (2) in Jordán et al. (2006). In this paper we use the Gaussian parameterization of the GCLF.

Most of our galaxies are small enough that the ACS WFC field of view is sufficient to encompass all of the galaxy’s GC system. In these cases, we determine the total number of GCs by counting GC candidates within our imaged field, correcting for incompleteness using the known mean and sigma of the GCLF. Calculating the completeness correction, however, is not straightforward because the surface brightness of the galaxy varies spatially. If, for instance, there were strong radial gradients in the GCLF or size distribution, we could be over- or undercorrecting for unseen GCs. Our previous work on these galaxies (Paper X; Paper XII) shows that this is not likely to be the case. We use bright GCs ( $z < 22.5$ ), those that are bright enough as to be complete over the entire range of galaxy surface brightness and GC sizes, to sample the “true” distribution of background surface brightnesses and GC sizes. We then make the assumption that the fainter GCs sample the same distributions to derive a mean completeness correction for the entire population.

In the more luminous galaxies, the ACS WFC field of view is insufficient to image the entire GC system. To correct for this, we determine the radial spatial density profile of the GC system, fit a Sérsic profile, and integrate over all radii to estimate the total number of GCs. For the largest galaxies, we supplemented the density profiles at larger radii with data from three sources: (1) ACSVCS imaging of companion galaxies whose GC systems are dominated by the halo population of the neighboring giant, (2) our WFPC2 parallel imaging of nearby halo fields (M. Takamiya et al. 2008, in preparation), and (3) ground-based data from the literature. In cases where there are no wide-field ground-based data, our additional pencil beam *HST* WFC and WFPC2 observations still do a comparably good job because the spatial resolution of *HST* minimizes the noise from background contamination. We use the same method described above for completeness correction, except we do so in bins of galactocentric radius instead of for the whole field of view.

Errors for  $N_{GC}$  were determined using a Monte Carlo technique where we introduced the appropriate random Poisson noise to our density profiles, then fitted and integrated the Sérsic (1968) profile. After doing this 1000 times, we used the distribution of estimated  $N_{GC}$  to determine the intervals that contained the closest 68% of the measurements to the mean. In Figure 1 we show the combined radial profile data for VCC 1316 (M87/NGC 4486, hereafter M87), which includes the central ACS pointing, four ACS pointings targeted at nearby companions, WFPC2 pointings, and the ground-based data of McLaughlin (1999a). All are in good

<sup>12</sup> See [http://www.astro.ex.ac.uk/people/aa/pages/computing/pyraf\\_pipeline.html](http://www.astro.ex.ac.uk/people/aa/pages/computing/pyraf_pipeline.html).

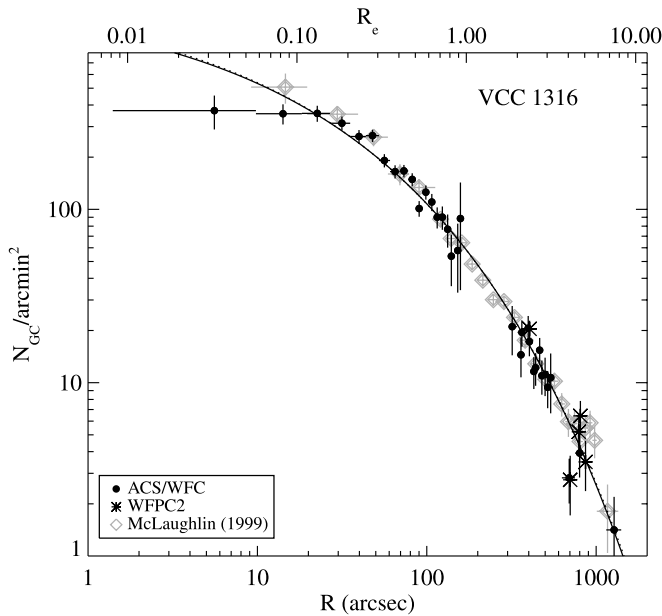


FIG. 1.—Surface density of globular clusters in VCC 1316 (M87/NGC 4486) as a function of projected galactocentric radius. We use data from three sources to fit the spatial density profile: ACS pointings of the central field and companion fields (*filled circles*), WFPC2 parallel pointings in adjacent blank fields (*asterisks*), and ground-based data from McLaughlin (1999a) (*diamonds*). All data are in good agreement and are well fitted by a Sérsic profile (*solid line*). We integrate the Sérsic profile to obtain the total number of GCs. [See the electronic edition of the *Journal* for a color version of this figure.]

agreement and are well fitted by a Sérsic profile, except perhaps in the very central regions, which are difficult to measure (especially from the ground) and do not contribute large numbers of GCs. The spatial density profiles of the GC systems of ACSVCS galaxies will be presented in E. W. Peng et al. (2008, in preparation), which will contain a more detailed description of the techniques. Below we outline our use of supplementary data in the larger galaxies.

*VCC 1226 (M49/NGC 4472, hereafter M49).*—For the most luminous galaxy in the Virgo Cluster, we used the GC counts in the fields of the neighboring compact galaxies VCC 1192 and VCC 1199, excluding a region of  $4R_e$  around the galaxies themselves, where  $R_e$  is the effective radius. We also used four WFPC2 parallel fields in the halo and the ground-based density profiles of McLaughlin (1999a) and Rhode & Zepf (2001). All these data are in excellent agreement.

*VCC 1316 (M87/NGC 4486).*—Four ACSVCS galaxies are near enough to this cluster cD that the GCs observed in their images are dominated by those of the giant: VCC 1327 (NGC 4486A), VCC 1297 (NGC 4486B), VCC 1279, and VCC 1250. VCC 1250 also appears to have ongoing star formation. As described above for M49, we masked a region of  $4R_e$  around these nearby neighbors and used the remaining GCs to constrain the density profile of the larger galaxy. In addition, we used five WFPC2 parallel fields and the ground-based density profiles from McLaughlin (1999a) (see Fig. 1).

*VCC 881 (NGC 4406).*—We supplemented the profile with two WFPC2 fields and the ground-based density profile of Rhode & Zepf (2004).

*VCC 763 (NGC 4374).*—We supplemented the profile with two WFPC2 fields and the ground-based density profile of Gómez & Richtler (2004).

*VCC 798, VCC 731, VCC 1535, VCC 1903, VCC 1632, VCC 1231, VCC 2095, VCC 1154, VCC 1062, VCC 1030, and*

*VCC 1664.*—For these galaxies, we extended the measured density profile using between one and three WFPC2 fields per galaxy, which were at galactocentric radii between  $5'$  and  $17'$ .

When we compare the total number of GCs derived using the integrated GC radial density profiles against the total number counted within the ACS WFC field of view, the numbers converge for galaxies with  $M_B > -18$ . Because integrated radial density profiles are more uncertain and not necessary in low-luminosity galaxies, we have adopted  $M_B = -18$  as a cutoff; for galaxies brighter than this, we use the Sérsic-integrated GC number counts, and for galaxies fainter than this, we use the corrected number of GCs directly counted within the ACS WFC image.

VCC 1938 has a close companion dE, VCC 1941, whose GC system complicates an accurate count of the larger galaxy's GC system. We mask an  $R = 70''$  region around VCC 1941 and count only GCs in the ACS WFC field of view, so our count of the VCC 1938 GC system is likely to be a lower limit.

All GC counts are corrected for foreground and background contamination using control fields (except in the eight cases described below). In many other studies, the background is taken from an annulus around the target galaxy. The advantages of using separate control fields are as follows: (1) We can sample a much larger area of sky to determine the mean surface density of contaminants, thus greatly reducing the Poisson errors introduced in contaminant subtraction. (2) We can use the full field of view for the target galaxy to measure as much of the GC system as we can. (3) We can still study systems where the GC system fills the entire field of view. (4) We can naturally incorporate spatially varying completeness as a function of position within the galaxy. The only potential disadvantage is the addition of cosmic variance into the error.

The advantage in using multiple control fields over a local background approach is most clear for the dwarf galaxies that have few GCs, where we need to minimize the error in the mean expected background. We determine the error in the background by measuring the number of GC-like objects selected in each of the 16 custom control fields and taking the standard deviation of these counts. This number combines both Poisson errors and cosmic variance and takes into account the varying selection and completeness from galaxy to galaxy. We find that the contribution of cosmic variance to the errors is on the order of the Poisson noise or less and that the use of control fields is superior to the use of a local background.

For eight galaxies, VCC 1327, VCC 1297, VCC 1279, VCC 1250, VCC 1185, VCC 1192, VCC 1199, and VCC 1178, control fields are insufficient, and we measure the total number of GCs by counting candidates in an  $R = 70''$  aperture and subtracting a local background. The first five are nearby neighbors to the giant elliptical M87 and so some (and sometimes all) GCs detected in their vicinity belong to the halos of the giants. The next three are neighbors of M49. The size of this aperture was chosen to fit the ACS WFC field of view and extends  $11R_e$ ,  $30R_e$ ,  $4R_e$ ,  $4R_e$ ,  $10R_e$ ,  $20R_e$ ,  $11R_e$ , and  $5R_e$  in radius for these galaxies, respectively. These numbers are, in principle, lower limits to the total numbers of GCs in each system.

In total, we present a homogeneous catalog of GC number counts for 100 ACSVCS galaxies. We have taken great care in this paper to produce our best estimate for the total number of GCs in the ACSVCS galaxies. Paper IX listed numbers of GCs used in their analysis of color distributions, but *we emphasize that those numbers were not corrected for field of view or completeness* and are best used only to evaluate the signal-to-noise ratio of the color distributions. Forbes (2005) and Miller & Lotz (2007) both derive  $S_N$  values from Paper IX (the Forbes [2005]

paper uses them exclusively), and so the appropriate warnings apply.

### 3.2. Total Luminosities and Masses of GCs

Another quantity of interest is the total luminosity (or stellar mass) in GCs. We take a straightforward empirical approach to measuring the total luminosity. For each galaxy, we add up the  $z$ -band luminosity in observed GCs down to 1 mag fainter than the mean of the GCLF, with adjustments for contaminants and completeness. For galaxies with  $M_B < -18$ , we apply an aperture correction to the total luminosity in GCs, which is the ratio of the total number of GCs determined from the integrated radial profile to the total number observed in the ACS WFC image. At this depth, we are directly counting 84% of the GCs in a Gaussian GCLF but are sampling 99% of the luminosity in GCs. Thus, completeness corrections are small, and it has been advocated (e.g., Harris 1991) that the luminosity in GCs is a more robust quantity than the number of GCs. To obtain the total mass in GCs, we assume that each GC is a 13 Gyr simple stellar population (SSP; Chaboyer et al. 1996) and use its  $(g - z)$  color in conjunction with the Bruzual & Charlot (2003, hereafter BC03) models with a Chabrier (2003) initial mass function (IMF) to obtain a mass-to-light ratio in the  $z$  bandpass. The range of  $\mathcal{M}/L_z$  variation with  $[\text{Fe}/\text{H}]$  for globular clusters is no more than  $\pm 20\%$ .

### 3.3. Galaxy Magnitudes: A Consistent Catalog of $M_V$ And $M_z$

Globular cluster formation efficiencies can be measured against the total luminosity or mass of a galaxy. Specific frequency is typically calculated using the absolute  $V$  magnitude. The VCC lists  $M_B$  for all of our sample galaxies, but these magnitudes were not derived from CCD photometry. We do not have  $V$  imaging for our galaxies, so we predict  $V$  using optical colors measured from SDSS imaging. We use  $(ugriz)_{\text{SDSS}}$  photometry measured directly from images in the SDSS Data Release 5 (Adelman-McCarthy et al. 2007). In many cases the large sizes of the galaxies compared to the size of the SDSS CCDs required careful stitching of neighboring runs and camera columns to produce flat images with matching sky (West et al. 2007). Using these specially prepared images, we measured total magnitudes in the five SDSS bands for all the ACSVCS galaxies using a growth curve analysis. Details of this analysis and the full catalog of magnitudes and colors will appear in C.-W. Chen et al. (2008, in preparation).

For the purposes of this paper, we only use the SDSS photometry for two objectives, deriving a  $V$  magnitude and supplementing the total  $z$  magnitudes given in Paper VI derived from ACS imaging. For determining each galaxy's  $V$  magnitude, we fit SSPs from BC03 to the four optical colors,  $(u - g)_{\text{SDSS}}$ ,  $(g - r)_{\text{SDSS}}$ ,  $(r - i)_{\text{SDSS}}$ , and  $(i - z)_{\text{SDSS}}$ , and use the  $g_{\text{SDSS}} - V$  color of the best-fit model. Across the entire sample, the  $g_{\text{SDSS}} - V$  color ranges from 0.18 to 0.48 with a mean of 0.37. In practice, this approach is extremely robust because  $g_{\text{SDSS}} - V$  is almost entirely a function of  $(g - r)_{\text{SDSS}}$ . We subtract this color from the measured  $g_{\text{SDSS}}$  magnitude to obtain  $V$ . The mean relationship between  $V$  and  $(g - r)_{\text{SDSS}}$  matches the empirical relation of Blanton & Roweis (2007), except that our magnitudes are fainter in the mean by 0.03 mag.

Because redder wavelengths are a better tracer of stellar mass than the traditionally used  $B$  and  $V$ , we often use the total integrated  $z$  luminosities of galaxies in this paper. For most of these galaxies, we use the  $z$  magnitudes measured from the ACS images as presented in Paper VI. However, for the brightest 10 galaxies, the ACS field of view was substantially smaller than the extent of the galaxy, and we prefer to use the wide-field SDSS photometry (after introducing a small  $[<0.04 \text{ mag}]$  color-dependent

correction between the SDSS and the ACS photometric systems). We also use the SDSS photometry for VCC 1030, VCC 575, and VCC 1512, which have suspect integrated magnitudes from their ACS surface brightness profiles (see notes on these galaxies in Paper VI). Otherwise, the two independent measures of the total luminosity are in good agreement. We choose to use the ACS magnitudes for the remaining 87 galaxies because they have higher signal-to-noise ratio than the SDSS photometry, which were taken from the ground with much higher sky backgrounds and shorter exposures.

### 3.4. Stellar Mass

Stellar mass is a better basis for comparison when studying galaxies of different morphological types and star formation histories. The  $V$ -band light in star-forming galaxies is affected significantly by young stars and is not the best tracer of the total mass. Using the  $(g - z)$  colors from Paper VI,  $J - K_s$  colors from the Two Micron All Sky Survey (2MASS; Skrutskie et al. 2006), and the model SSPs from BC03 with a Chabrier (2003) IMF, we obtain mean mass-to-light ratios in the  $z$  bandpass ( $\mathcal{M}/L_z$ ) from which we can derive total stellar masses. In the near-infrared, we measure the  $J$  and  $K_s$  galaxy magnitudes in images taken by 2MASS and supplement them with those in the 2MASS Extended Source Catalog (XSC) and Large Galaxy Atlas (Jarrett et al. 2003). For nine galaxies with unreliable  $K_s$  photometry from 2MASS (those not included in the XSC), we use only the  $(g - z)$  color assuming a 10 Gyr SSP to determine  $\mathcal{M}/L_z$ . We assume a younger age than for GCs because dEs are measured to have younger ages than GCs or massive ellipticals (Geha et al. 2003). Although  $\mathcal{M}/L$  is related to both age and metallicity, it is easier to determine  $\mathcal{M}/L$  than age or metallicity individually. Most of the sample has  $\mathcal{M}/L_z$  consistent with old stellar populations (age  $> 5$  Gyr), although a few galaxies are dE/dI transition objects and have noticeably younger mean ages and correspondingly lower  $\mathcal{M}/L_z$  (the most extreme example being VCC 1499). For one of these galaxies, VCC 1030, we derive a low mean age, high metallicity, and a correspondingly low  $\mathcal{M}/L_z$ . While this galaxy may be interacting, the colors may also be suspect due to a large-scale central dust disk. Until we get more data, we consider  $\mathcal{M}/L$  for this galaxy uncertain.

Fortunately, the  $z$  bandpass is less sensitive to recent star formation than bluer bandpasses, and the range of  $\mathcal{M}/L_z$  is only 0.4–2.1, a dynamic range that is a factor of  $\sim 3$  smaller than that for  $\mathcal{M}/L_B$ . The error in  $\mathcal{M}/L_z$  is dominated by the quality of the infrared photometry. We estimate errors in  $\mathcal{M}/L_z$  with a Monte Carlo procedure where the  $J - K_s$  color is perturbed by a random amount drawn from a Gaussian distribution with sigma equal to the claimed photometric error. The  $\mathcal{M}/L_z$  is then recalculated in this fashion 1000 times and the error is the half-width of the middle 68% of the distribution. The typical error in  $\mathcal{M}/L_z$  is  $\pm 0.4$ , or  $\sim 25\%$ . These errors are propagated forward into mass-normalized quantities, although any possible variations in the IMF are not included.

All of the global galaxy properties discussed above are listed in Table 1, and all the quantities related to the globular cluster systems are presented in Table 2.

## 4. RESULTS

### 4.1. $S_N$ versus Galaxy Magnitude

The historical definition of specific frequency,  $S_N$ , is the number of GCs per unit  $M_V = -15$  of galaxy luminosity, or

$$S_N = N_{\text{GC}} \times 10^{0.4(M_V + 15)} \quad (1)$$

TABLE 1  
GLOBAL PROPERTIES OF ACS VIRGO CLUSTER SURVEY GALAXIES

Number (1)	VCC (2)	$M_V$ (3)	$M_z$ (4)	$R_p$ (Mpc) (5)	$R_{3D}$ (Mpc) (6)	$L_{z,*}$ ( $10^9 L_\odot$ ) (7)	$M_*$ ( $10^9 M_\odot$ ) (8)
1.....	1226	-22.90	-23.87	1.27	$1.28 \pm 0.39$	263.04	$531.84 \pm 110.48$
2.....	1316	-22.66	-23.86	0.00	0.00	188.40	$302.27 \pm 79.13$
3.....	1978	-22.41	-23.42	0.94	$1.04 \pm 0.46$	167.62	$339.32 \pm 50.28$
4.....	881	-22.54	-23.53	0.36	$1.49 \pm 0.42$	181.42	$289.63 \pm 59.87$
5.....	798	-22.34	-23.25	1.71	$2.09 \pm 0.42$	132.84	$186.50 \pm 43.84$
6.....	763	-22.29	-23.20	0.43	$2.09 \pm 0.44$	141.95	$236.49 \pm 61.04$
7.....	731	-22.31	-23.37	1.53	$6.76 \pm 0.54$	141.76	$226.21 \pm 52.45$
8.....	1535	-21.38	-22.40	1.37	...	57.88	$77.21 \pm 16.21$
9.....	1903	-21.18	-22.19	0.82	$1.87 \pm 0.35$	50.26	$83.87 \pm 19.10$
10.....	1632	-21.36	-22.33	0.34	$0.63 \pm 0.38$	60.33	$95.35 \pm 16.89$
11.....	1231	-20.70	-21.58	0.31	$1.30 \pm 0.36$	32.37	$53.46 \pm 12.30$
12.....	2095	-20.78	-20.95	1.59	...	33.74	$52.75 \pm 10.12$
13.....	1154	-20.85	-21.79	0.47	$0.66 \pm 0.38$	36.86	$77.95 \pm 13.64$
14.....	1062	-20.56	-21.34	0.77	$1.42 \pm 0.36$	28.65	$52.69 \pm 8.88$
15.....	2092	-20.46	-21.64	1.54	$1.61 \pm 0.38$	25.55	$47.32 \pm 7.66$
16.....	369	-20.19	-20.41	0.79	$1.04 \pm 0.37$	20.17	$32.46 \pm 7.67$
17.....	759	-20.44	-21.41	0.46	$0.60 \pm 0.40$	25.74	$45.15 \pm 10.04$
18.....	1692	-20.26	-21.04	1.54	$1.64 \pm 0.40$	21.85	$33.97 \pm 6.55$
19.....	1030	-20.45	-21.38	0.30	$0.33 \pm 0.39$	22.63	$13.32 \pm 9.05$
20.....	2000	-19.60	-20.66	1.03	$2.03 \pm 0.35$	11.71	$23.76 \pm 4.10$
21.....	685	-20.11	-21.04	1.33	...	19.12	$31.23 \pm 6.69$
22.....	1664	-19.90	-20.97	0.48	$0.77 \pm 0.38$	15.39	$26.00 \pm 6.15$
23.....	654	-19.88	-20.59	1.35	...	14.22	$22.74 \pm 4.12$
24.....	944	-19.80	-20.66	0.86	$1.05 \pm 0.38$	13.94	$29.14 \pm 4.88$
25.....	1938	-19.93	-20.86	0.89	$1.00 \pm 0.40$	15.71	$24.86 \pm 5.50$
26.....	1279	-19.78	-20.72	0.04	$0.47 \pm 0.40$	13.42	$22.03 \pm 4.03$
27.....	1720	-19.63	-20.68	0.94	$0.97 \pm 0.38$	11.55	$20.21 \pm 5.31$
28.....	355	-19.42	-20.41	1.07	$1.55 \pm 0.36$	10.06	$15.78 \pm 3.62$
29.....	1619	-19.36	-20.19	0.33	$1.51 \pm 0.36$	8.65	$17.24 \pm 3.12$
30.....	1883	-19.68	-20.73	1.65	$1.69 \pm 0.38$	11.19	$16.66 \pm 3.92$
31.....	1242	-19.38	-20.29	0.49	$1.42 \pm 0.43$	9.07	$15.27 \pm 3.17$
32.....	784	-19.31	-20.26	1.00	$1.30 \pm 0.37$	8.48	$16.79 \pm 3.14$
33.....	1537	-18.99	-19.95	0.39	$1.18 \pm 0.36$	6.18	$10.24 \pm 2.35$
34.....	778	-19.56	-20.03	0.79	$1.17 \pm 0.41$	10.73	$18.18 \pm 4.19$
35.....	1321	-18.80	-19.93	1.26	$1.99 \pm 0.35$	4.86	$6.98 \pm 1.75$
36.....	828	-19.15	-20.07	0.38	$1.42 \pm 0.42$	7.51	$13.69 \pm 3.00$
37.....	1250	-18.98	-20.10	0.06	$0.79 \pm 0.57$	5.23	$3.68 \pm 2.04$
38.....	1630	-19.07	-20.06	0.34	$0.51 \pm 0.38$	7.03	$11.61 \pm 2.39$
39.....	1146	-18.93	-19.64	0.28	$0.60 \pm 0.38$	5.71	$8.69 \pm 2.17$
40.....	1025	-19.58	-20.57	1.24	$5.98 \pm 0.53$	10.90	$21.39 \pm 4.25$
41.....	1303	-18.84	-19.80	0.97	$0.98 \pm 0.39$	5.55	$10.39 \pm 2.33$
42.....	1913	-18.87	-19.74	1.58	$1.67 \pm 0.40$	5.67	$10.74 \pm 2.72$
43.....	1327	-19.10	-19.81	0.04	$1.45 \pm 0.51$	8.36	$17.24 \pm 2.34$
44.....	1125	-19.07	-19.44	0.24	...	6.80	$8.06 \pm 3.40$
45.....	1475	-18.56	-19.42	1.13	$1.15 \pm 0.38$	4.07	$7.73 \pm 1.50$
46.....	1178	-18.35	-19.35	1.22	$1.48 \pm 0.37$	3.64	$7.09 \pm 1.42$
47.....	1283	-18.65	-19.70	0.34	$0.86 \pm 0.41$	4.83	$9.08 \pm 1.79$
48.....	1261	-18.42	-19.21	0.47	$1.44 \pm 0.51$	3.14	$4.87 \pm 1.69$
49.....	698	-18.78	-19.71	0.58	$1.80 \pm 0.43$	5.03	$9.52 \pm 2.56$
50.....	1422	-17.97	-18.73	0.62	$1.60 \pm 0.43$	2.18	$3.82 \pm 1.35$
51.....	2048	-17.85	-18.57	1.32	...	1.89	$2.94 \pm 1.08$
52.....	1871	-17.31	-18.32	0.79	$1.43 \pm 0.44$	1.36	$2.26 \pm 0.58$
53.....	9	-18.04	-18.80	1.57	$1.64 \pm 0.65$	2.17	$3.08 \pm 0.61$
54.....	575	-18.42	-19.33	1.34	$5.30 \pm 0.61$	3.56	$5.00 \pm 1.75$
55.....	1910	-17.39	-18.33	0.82	$1.11 \pm 0.37$	1.43	$2.10 \pm 0.81$
56.....	1049	-16.69	-17.34	1.26	$1.43 \pm 0.53$	0.60	$0.52 \pm 0.40$
57.....	856	-17.57	-18.25	0.76	$0.76 \pm 0.47$	1.47	$2.22 \pm 0.85$
58.....	140	-17.51	-18.27	1.24	$1.30 \pm 0.46$	1.37	$2.33 \pm 0.71$
59.....	1355	-17.51	-18.16	0.50	$0.52 \pm 0.63$	1.29	$1.82 \pm 0.35$
60.....	1087	-17.79	-18.64	0.25	$0.34 \pm 0.46$	1.84	$3.29 \pm 1.07$

TABLE 1—*Continued*

Number (1)	VCC (2)	$M_V$ (3)	$M_z$ (4)	$R_p$ (Mpc) (5)	$R_{3D}$ (Mpc) (6)	$L_{z,*}$ ( $10^9 L_\odot$ ) (7)	$M_*$ ( $10^9 M_\odot$ ) (8)
61.....	1297	-17.67	-18.75	0.04	$0.23 \pm 0.46$	1.95	$3.96 \pm 0.62$
62.....	1861	-17.60	-18.34	0.80	$1.15 \pm 0.45$	1.63	$2.88 \pm 1.00$
63.....	543	-17.41	-18.19	0.91	$1.44 \pm 0.44$	1.27	$2.19 \pm 0.65$
64.....	1431	-17.39	-18.25	0.34	$0.90 \pm 0.45$	1.37	$2.20 \pm 0.74$
65.....	1528	-17.16	-18.04	0.34	$0.70 \pm 0.45$	1.06	$1.63 \pm 0.48$
66.....	1695	-17.49	-18.32	0.43	$0.46 \pm 0.54$	1.30	$1.69 \pm 0.78$
67.....	1833	-17.13	-17.89	1.22	$1.33 \pm 0.45$	0.97	$0.86 \pm 0.48$
68.....	437	-17.82	-18.69	1.62	$1.65 \pm 0.48$	1.92	$2.80 \pm 1.25$
69.....	2019	-17.36	-18.18	1.08	$1.13 \pm 0.48$	1.17	$1.02 \pm 0.72$
70.....	33	-16.39	-17.01	1.48	$2.19 \pm 0.57$	0.46	$0.43 \pm 0.25$
71.....	200	-17.12	-17.80	1.02	$1.70 \pm 0.59$	0.95	$1.34 \pm 0.49$
72.....	571	-17.32	-18.29	1.40	$7.27 \pm 1.01$	1.07	$0.98 \pm 0.65$
73.....	21	-16.83	-17.38	1.58	...	0.58	$0.56 \pm 0.37$
74.....	1488	-16.78	-17.34	0.88	...	0.57	$0.41 \pm 0.38$
75.....	1779	-16.90	-17.38	0.89	...	0.63	$1.02 \pm 0.38$
76.....	1895	-16.60	-17.29	1.16	$1.47 \pm 0.37$	0.57	$0.76 \pm 0.39$
77.....	1499	-16.53	-16.95	0.22	...	0.34	$0.14 \pm 0.19$
78.....	1545	-16.91	-17.74	0.26	$0.30 \pm 0.54$	0.87	$1.41 \pm 0.42$
79.....	1192	-16.86	-18.14	1.27	...	0.94	$1.85 \pm 0.62$
80.....	1857	-16.64	-17.32	0.90	...	0.53	$0.73 \pm 0.14$
81.....	1075	-16.78	-17.55	0.63	$0.87 \pm 0.68$	0.73	$1.11 \pm 0.47$
82.....	1948	-16.06	-16.78	0.99	...	0.35	$0.50 \pm 0.09$
83.....	1627	-16.46	-17.39	0.34	$2.19 \pm 0.34$	0.59	$1.03 \pm 0.32$
84.....	1440	-16.86	-17.66	0.88	$1.21 \pm 0.45$	0.78	$1.19 \pm 0.44$
85.....	230	-16.21	-16.96	0.96	$1.41 \pm 0.84$	0.43	$0.69 \pm 0.29$
86.....	2050	-16.36	-17.15	1.16	$1.47 \pm 0.52$	0.47	$0.29 \pm 0.26$
87.....	1993	-16.30	-17.02	0.95	$1.03 \pm 0.38$	0.48	$0.75 \pm 0.27$
88.....	751	-16.97	-17.92	1.72	$2.05 \pm 0.44$	0.92	$1.41 \pm 0.51$
89.....	1828	-16.78	-17.47	0.67	$0.67 \pm 0.55$	0.72	$1.26 \pm 0.47$
90.....	538	-16.24	-17.10	1.62	$6.47 \pm 0.86$	0.44	$0.63 \pm 0.26$
91.....	1407	-16.72	-17.43	0.17	$0.23 \pm 0.39$	0.68	$1.24 \pm 0.42$
92.....	1886	-16.25	-16.93	0.76	...	0.36	$0.51 \pm 0.10$
93.....	1199	-15.47	-16.94	1.25	...	0.29	$0.58 \pm 0.16$
94.....	1743	-16.33	-16.95	0.84	$1.27 \pm 0.75$	0.45	$0.40 \pm 0.29$
95.....	1539	-16.05	-17.12	0.25	$0.26 \pm 0.87$	0.37	$0.52 \pm 0.11$
96.....	1185	-16.77	-17.37	0.10	$0.10 \pm 0.79$	0.78	$0.89 \pm 0.51$
97.....	1826	-16.03	-16.71	0.98	$1.11 \pm 0.61$	0.35	$0.61 \pm 0.22$
98.....	1512	-16.25	-16.54	0.38	$1.50 \pm 0.42$	0.43	$0.61 \pm 0.12$
99.....	1489	-15.61	-16.41	0.45	...	0.22	$0.32 \pm 0.06$
100.....	1661	-15.81	-17.40	0.70	$1.26 \pm 1.20$	0.29	$0.43 \pm 0.08$

NOTES.—Col. (1): Running number, sorted by increasing  $B_T$  magnitude. Col. (2): Number in Virgo Cluster Catalog. Col. (3): Absolute  $V$  magnitude. Col. (4): Absolute  $z$  magnitude. Col. (5): Projected distance from M87 (VCC 1316). Col. (6): Three-dimensional distance from M87 (VCC 1316), using the polynomial calibration from Paper XIII. Col. (7): Stellar  $z$  luminosity. Col. (8): Stellar mass.

(Harris & van den Bergh 1981). This number is approximately unity for the Milky Way. One of the motivations for calculating  $S_N$  is to understand whether GC formation scales with the bulk of star formation in the same way across galaxy types and masses. In Figure 2 we show the behavior of  $S_N$  as a function of  $M_V$  in our sample of 100 ACSVCS galaxies (with binned values listed in Table 3). The luminous early-type galaxies, mostly giant ellipticals, are well known to have  $S_N \sim 2-5$ . The cD galaxy, M87, has  $S_N = 12.6 \pm 0.8$ , which is consistent with other measures of its  $S_N$  (e.g., Harris et al. 1998). Early-type galaxies of intermediate luminosity ( $-22 < M_V < -18$ ) have a nearly uniformly low  $\langle S_N \rangle \sim 1.5$ . Galaxies in this luminosity range also have a tendency to be lenticular, with the VCC classifying 70% of these ACSVCS intermediate-luminosity galaxies as S0, E/S0, or S0/E. By contrast, the fainter galaxies ( $M_V > -18$ ) exhibit a

wide range of  $S_N$ , with values as low as zero and as high as those of M87.<sup>13</sup>

This trend has been hinted at previously from studies of giant ellipticals, lenticulars, and dwarf ellipticals. In Figure 3 we show the ACSVCS sample combined with data compiled from the literature. At the high-luminosity end, we take  $S_N$  values from the compilation of Ashman & Zepf (1998), using ones that were reliably determined from CCD data (*gray diamonds*). These include estimates from Kissler-Patig et al. (1996, 1997), Dirsch et al. (2003a, 2003b, 2005), Forbes et al. (1996), Rhode & Zepf (2004),

<sup>13</sup> Strader et al. (2006) analyzed a subset of these data and claimed a bimodal distribution of  $S_N$  in the dEs. A histogram of the  $S_N$  distribution shows that it is better described as strongly peaked with a tail to higher  $S_N$ , especially when one includes dS0s.



TABLE 2  
GC FORMATION EFFICIENCIES IN ACS VIRGO CLUSTER SURVEY GALAXIES

Number (1)	VCC (2)	$N_{GC}$ (3)	$S_N$ (4)	$S_{N,z}$ (5)	$T$ (6)	$S_L$ (7)	$S_{M}$ (8)	$S_{N,z,blue}$ (9)	$S_{N,z,red}$ (10)	$T_{blue}$ (11)	$T_{red}$ (12)	$f_{red}$ (13)
1.....	1226	7813 ± 830	5.40 ± 0.57	2.20 ± 0.23	14.7 ± 3.4	1.23 ± 0.01	0.87 ± 0.18	1.57 ± 0.19	0.64 ± 0.14	10.4 ± 2.5	4.3 ± 1.3	0.29
2.....	1316	14660 ± 891	12.59 ± 0.77	4.19 ± 0.25	48.5 ± 13.0	2.33 ± 0.02	2.83 ± 0.74	3.04 ± 0.21	1.14 ± 0.15	35.2 ± 9.5	13.2 ± 3.9	0.27
3.....	1978	4745 ± 1099	5.16 ± 1.20	2.03 ± 0.47	14.0 ± 3.8	1.15 ± 0.02	0.84 ± 0.12	1.28 ± 0.38	0.75 ± 0.27	8.8 ± 2.9	5.2 ± 2.0	0.37
4.....	881	2660 ± 129	2.57 ± 0.12	1.03 ± 0.05	9.2 ± 1.9	0.56 ± 0.01	0.49 ± 0.10	0.87 ± 0.04	0.16 ± 0.03	7.7 ± 1.6	1.5 ± 0.4	0.16
5.....	798	1110 ± 181	1.29 ± 0.21	0.56 ± 0.09	6.0 ± 1.7	0.20 ± 0.01	0.21 ± 0.05	0.33 ± 0.07	0.23 ± 0.05	3.5 ± 1.1	2.4 ± 0.8	0.41
6.....	763	4301 ± 1201	5.20 ± 1.45	2.26 ± 0.63	18.2 ± 6.9	1.11 ± 0.03	0.89 ± 0.23	2.01 ± 0.51	0.24 ± 0.37	16.2 ± 5.9	2.0 ± 3.0	0.11
7.....	731	3246 ± 598	3.86 ± 0.71	1.45 ± 0.27	14.3 ± 4.2	0.85 ± 0.02	0.86 ± 0.20	0.80 ± 0.22	0.65 ± 0.16	7.9 ± 2.8	6.5 ± 2.2	0.45
8.....	1535	388 ± 117	1.09 ± 0.33	0.42 ± 0.13	5.0 ± 1.8	0.20 ± 0.00	0.25 ± 0.05	0.26 ± 0.10	0.16 ± 0.07	3.1 ± 1.4	1.9 ± 1.0	0.38
9.....	1903	803 ± 355	2.70 ± 1.19	1.07 ± 0.47	9.6 ± 4.8	0.68 ± 0.01	0.61 ± 0.14	0.55 ± 0.38	0.52 ± 0.27	4.9 ± 3.6	4.7 ± 2.7	0.49
10.....	1632	984 ± 198	2.82 ± 0.57	1.15 ± 0.23	10.3 ± 2.8	0.65 ± 0.01	0.61 ± 0.11	0.69 ± 0.19	0.46 ± 0.14	6.2 ± 2.0	4.1 ± 1.4	0.40
11.....	1231	376 ± 97	1.98 ± 0.51	0.88 ± 0.23	7.0 ± 2.4	0.32 ± 0.01	0.28 ± 0.06	0.50 ± 0.18	0.38 ± 0.13	4.0 ± 1.7	3.0 ± 1.3	0.43
12.....	2095	211 ± 34	1.03 ± 0.17	0.88 ± 0.14	4.0 ± 1.0	0.16 ± 0.00	0.14 ± 0.03	0.52 ± 0.11	0.35 ± 0.08	2.4 ± 0.7	1.6 ± 0.5	0.40
13.....	1154	218 ± 28	1.00 ± 0.13	0.42 ± 0.05	2.8 ± 0.6	0.12 ± 0.00	0.09 ± 0.02	0.22 ± 0.04	0.20 ± 0.03	1.5 ± 0.4	1.3 ± 0.3	0.48
14.....	1062	178 ± 30	1.06 ± 0.18	0.52 ± 0.09	3.4 ± 0.8	0.21 ± 0.00	0.16 ± 0.03	0.52 ± 0.07	0.00 ± 0.05	3.4 ± 0.7	0.0 ± 0.3	0.00
15.....	2092	103 ± 17	0.68 ± 0.11	0.23 ± 0.04	2.2 ± 0.5	0.08 ± 0.00	0.07 ± 0.01	0.13 ± 0.03	0.09 ± 0.02	1.3 ± 0.4	0.9 ± 0.3	0.39
16.....	369	179 ± 17	1.51 ± 0.14	1.23 ± 0.12	5.5 ± 1.4	0.28 ± 0.01	0.25 ± 0.06	0.58 ± 0.09	0.65 ± 0.08	2.6 ± 0.7	2.9 ± 0.8	0.53
17.....	759	200 ± 41	1.34 ± 0.27	0.54 ± 0.11	4.4 ± 1.3	0.24 ± 0.00	0.20 ± 0.05	0.35 ± 0.09	0.19 ± 0.07	2.9 ± 1.0	1.6 ± 0.6	0.35
18.....	1692	139 ± 23	1.09 ± 0.18	0.53 ± 0.09	4.1 ± 1.0	0.16 ± 0.01	0.16 ± 0.03	0.34 ± 0.07	0.19 ± 0.05	2.6 ± 0.7	1.5 ± 0.5	0.36
19 <sup>a</sup> .....	1030	345 ± 80	2.27 ± 0.53	0.97 ± 0.22	25.9 ± 18.6	0.40 ± 0.01	1.06 ± 0.72	0.73 ± 0.18	0.24 ± 0.13	19.5 ± 14.1	6.4 ± 5.6	0.25
20.....	2000	205 ± 28	2.97 ± 0.41	1.12 ± 0.15	8.6 ± 1.9	0.53 ± 0.01	0.37 ± 0.06	0.97 ± 0.13	0.15 ± 0.08	7.5 ± 1.6	1.2 ± 0.7	0.13
21.....	685	196 ± 60	1.78 ± 0.54	0.75 ± 0.23	6.3 ± 2.3	0.40 ± 0.01	0.36 ± 0.08	0.54 ± 0.19	0.22 ± 0.13	4.5 ± 1.8	1.8 ± 1.2	0.29
22.....	1664	213 ± 31	2.35 ± 0.34	0.87 ± 0.13	8.2 ± 2.3	0.42 ± 0.01	0.37 ± 0.09	0.54 ± 0.10	0.33 ± 0.07	5.1 ± 1.5	3.1 ± 1.0	0.38
23.....	654	45 ± 16	0.50 ± 0.18	0.26 ± 0.09	2.0 ± 0.8	0.05 ± 0.00	0.05 ± 0.01	0.23 ± 0.08	0.04 ± 0.05	1.7 ± 0.7	0.3 ± 0.4	0.15
24.....	944	72 ± 10	0.87 ± 0.12	0.39 ± 0.05	2.5 ± 0.5	0.11 ± 0.01	0.08 ± 0.01	0.25 ± 0.04	0.14 ± 0.03	1.6 ± 0.4	0.9 ± 0.3	0.36
25 <sup>b</sup> .....	1938	59.8 ± 9.2	0.64 ± 0.10	0.27 ± 0.04	2.4 ± 0.6	0.12 ± 0.01	0.11 ± 0.02	0.22 ± 0.04	0.05 ± 0.02	1.9 ± 0.5	0.5 ± 0.2	0.19
26 <sup>c</sup> .....	1279	58 ± 11	0.72 ± 0.15	0.30 ± 0.06	2.7 ± 0.7	0.13 ± 0.01	0.11 ± 0.02	0.21 ± 0.05	0.09 ± 0.04	1.8 ± 0.6	0.8 ± 0.3	0.30
27.....	1720	62 ± 13	0.87 ± 0.18	0.33 ± 0.07	3.1 ± 1.0	0.12 ± 0.00	0.10 ± 0.03	0.22 ± 0.06	0.12 ± 0.04	2.0 ± 0.7	1.1 ± 0.5	0.36
28.....	355	100 ± 31	1.70 ± 0.53	0.69 ± 0.21	6.3 ± 2.4	0.11 ± 0.02	0.11 ± 0.03	0.53 ± 0.17	0.16 ± 0.12	4.9 ± 2.0	1.4 ± 1.2	0.23
29.....	1619	84 ± 19	1.52 ± 0.34	0.70 ± 0.16	4.9 ± 1.4	0.14 ± 0.01	0.10 ± 0.02	0.52 ± 0.13	0.19 ± 0.09	3.6 ± 1.1	1.3 ± 0.7	0.27
30.....	1883	83 ± 25	1.11 ± 0.34	0.42 ± 0.13	5.0 ± 1.9	0.16 ± 0.00	0.16 ± 0.04	0.29 ± 0.10	0.13 ± 0.07	3.4 ± 1.5	1.6 ± 0.9	0.31
31.....	1242	116 ± 24	2.05 ± 0.42	0.88 ± 0.18	7.6 ± 2.2	0.31 ± 0.01	0.26 ± 0.05	0.57 ± 0.15	0.31 ± 0.11	4.9 ± 1.6	2.7 ± 1.1	0.35
32.....	784	50 ± 14	0.95 ± 0.27	0.39 ± 0.11	3.0 ± 1.0	0.07 ± 0.01	0.05 ± 0.01	0.23 ± 0.09	0.17 ± 0.07	1.7 ± 0.7	1.3 ± 0.6	0.44
33.....	1537	31.4 ± 7.2	0.80 ± 0.18	0.33 ± 0.08	3.1 ± 1.0	0.00 ± 0.00	0.00 ± 0.00	0.29 ± 0.07	0.04 ± 0.03	2.7 ± 0.9	0.4 ± 0.3	0.12
34.....	778	74 ± 32	1.11 ± 0.48	0.72 ± 0.31	4.1 ± 2.0	0.14 ± 0.01	0.12 ± 0.03	0.51 ± 0.25	0.21 ± 0.18	2.9 ± 1.6	1.2 ± 1.1	0.29
35.....	1321	31.0 ± 9.0	0.94 ± 0.27	0.33 ± 0.10	4.4 ± 1.7	0.09 ± 0.01	0.08 ± 0.02	0.24 ± 0.08	0.09 ± 0.05	3.3 ± 1.3	1.2 ± 0.8	0.27

TABLE 2—Continued

Number (1)	VCC (2)	$N_{GC}$ (3)	$S_V$ (4)	$S_{N,z}$ (5)	$T$ (6)	$S_L$ (7)	$S_M$ (8)	$S_{N,z,blue}$ (9)	$S_{N,z,red}$ (10)	$T_{blue}$ (11)	$T_{red}$ (12)	$f_{red}$ (13)
36.....	828	69.5 ± 9.8	1.52 ± 0.21	0.65 ± 0.09	5.1 ± 1.3	0.00 ± 0.00	0.00 ± 0.00	0.50 ± 0.08	0.15 ± 0.05	3.9 ± 1.1	1.2 ± 0.4	0.23
37 <sup>e</sup> .....	1250	20.1 ± 7.3	0.52 ± 0.19	0.18 ± 0.07	5.5 ± 3.6	0.06 ± 0.01	0.14 ± 0.08	0.17 ± 0.06	0.01 ± 0.03	5.0 ± 3.3	0.4 ± 1.0	0.06
38.....	1630	47 ± 11	1.11 ± 0.26	0.44 ± 0.10	4.0 ± 1.3	0.12 ± 0.02	0.11 ± 0.03	0.26 ± 0.08	0.18 ± 0.06	2.4 ± 0.9	1.7 ± 0.7	0.41
39.....	1146	72 ± 12	1.93 ± 0.32	1.00 ± 0.17	8.3 ± 2.5	0.39 ± 0.02	0.40 ± 0.10	0.30 ± 0.12	0.70 ± 0.12	2.5 ± 1.1	5.8 ± 1.8	0.70
40.....	1025	141 ± 34	2.08 ± 0.50	0.84 ± 0.20	6.6 ± 2.1	0.25 ± 0.07	0.19 ± 0.08	0.75 ± 0.17	0.09 ± 0.11	5.9 ± 1.8	0.7 ± 0.9	0.11
41.....	1303	72 ± 18	2.10 ± 0.53	0.87 ± 0.22	6.9 ± 2.3	0.24 ± 0.03	0.18 ± 0.05	0.82 ± 0.18	0.05 ± 0.11	6.5 ± 2.1	0.4 ± 0.9	0.06
42.....	1913	71 ± 14	2.02 ± 0.40	0.90 ± 0.18	6.6 ± 2.1	0.31 ± 0.01	0.23 ± 0.06	0.74 ± 0.15	0.16 ± 0.09	5.4 ± 1.8	1.2 ± 0.8	0.18
43 <sup>e</sup> .....	1327	11 ± 12	0.26 ± 0.29	0.14 ± 0.15	0.7 ± 0.7	0.07 ± 0.00	0.05 ± 0.01	0.09 ± 0.12	0.05 ± 0.09	0.4 ± 0.6	0.2 ± 0.4	0.36
44.....	1125	52.3 ± 8.5	1.23 ± 0.20	0.88 ± 0.14	6.5 ± 2.9	0.16 ± 0.01	0.20 ± 0.08	0.81 ± 0.13	0.07 ± 0.06	6.0 ± 2.7	0.5 ± 0.5	0.08
45.....	1475	81 ± 10	3.05 ± 0.40	1.38 ± 0.18	10.5 ± 2.5	0.33 ± 0.03	0.23 ± 0.05	1.33 ± 0.17	0.06 ± 0.06	10.1 ± 2.4	0.4 ± 0.5	0.04
46 <sup>e</sup> .....	1178	25.3 ± 9.2	1.16 ± 0.42	0.46 ± 0.17	3.6 ± 1.5	0.12 ± 0.01	0.09 ± 0.02	0.31 ± 0.14	0.15 ± 0.10	2.4 ± 1.2	1.2 ± 0.8	0.33
47.....	1283	58.6 ± 9.3	2.02 ± 0.32	0.77 ± 0.12	6.5 ± 1.6	0.22 ± 0.03	0.16 ± 0.04	0.59 ± 0.10	0.18 ± 0.06	5.0 ± 1.3	1.5 ± 0.6	0.23
48.....	1261	35.1 ± 7.6	1.51 ± 0.33	0.72 ± 0.16	7.2 ± 2.9	0.29 ± 0.02	0.28 ± 0.10	0.60 ± 0.14	0.12 ± 0.08	6.0 ± 2.5	1.2 ± 0.9	0.17
49.....	698	114 ± 12	3.53 ± 0.38	1.50 ± 0.16	12.0 ± 3.5	0.56 ± 0.01	0.41 ± 0.11	1.23 ± 0.14	0.27 ± 0.07	9.8 ± 2.9	2.2 ± 0.8	0.18
50.....	1422	24.9 ± 6.0	1.62 ± 0.39	0.80 ± 0.19	6.5 ± 2.8	0.21 ± 0.02	0.17 ± 0.06	0.54 ± 0.16	0.26 ± 0.11	4.4 ± 2.0	2.1 ± 1.2	0.32
51.....	2048	17.2 ± 5.4	1.25 ± 0.39	0.64 ± 0.20	5.9 ± 2.8	0.19 ± 0.02	0.19 ± 0.07	0.49 ± 0.17	0.15 ± 0.11	4.5 ± 2.3	1.3 ± 1.1	0.23
52.....	1871	10.4 ± 5.0	1.24 ± 0.60	0.49 ± 0.24	4.6 ± 2.5	0.21 ± 0.07	0.17 ± 0.07	0.48 ± 0.21	0.01 ± 0.11	4.5 ± 2.3	0.1 ± 1.0	0.02
53.....	9	25.7 ± 6.4	1.57 ± 0.39	0.78 ± 0.19	8.3 ± 2.6	0.21 ± 0.02	0.20 ± 0.05	0.68 ± 0.17	0.09 ± 0.09	7.3 ± 2.4	1.0 ± 1.0	0.12
54.....	575	18.0 ± 6.1	0.77 ± 0.26	0.33 ± 0.11	3.6 ± 1.7	0.00 ± 0.06	-0.00 ± 0.06	0.26 ± 0.10	0.07 ± 0.06	2.8 ± 1.4	0.8 ± 0.7	0.21
55.....	1910	48.7 ± 8.4	5.38 ± 0.93	2.27 ± 0.39	23.2 ± 9.8	1.09 ± 0.04	1.12 ± 0.43	1.47 ± 0.32	0.79 ± 0.23	15.1 ± 6.7	8.1 ± 3.9	0.35
56.....	1049	8.4 ± 4.5	1.77 ± 0.95	0.97 ± 0.52	16.2 ± 15.1	0.05 ± 0.03	0.09 ± 0.08	0.75 ± 0.44	0.22 ± 0.28	12.4 ± 12.0	3.7 ± 5.5	0.23
57.....	856	43.4 ± 7.9	4.07 ± 0.74	2.17 ± 0.39	19.5 ± 8.3	0.55 ± 0.03	0.56 ± 0.22	1.80 ± 0.35	0.37 ± 0.19	16.2 ± 6.9	3.3 ± 2.1	0.17
58.....	140	21.3 ± 6.1	2.10 ± 0.60	1.05 ± 0.30	9.2 ± 3.8	0.18 ± 0.03	0.16 ± 0.06	0.79 ± 0.25	0.26 ± 0.16	6.9 ± 3.1	2.3 ± 1.6	0.25
59.....	1355	10.8 ± 5.6	1.07 ± 0.56	0.59 ± 0.30	5.9 ± 3.3	0.05 ± 0.01	0.07 ± 0.02	0.53 ± 0.26	0.06 ± 0.15	5.3 ± 2.8	0.6 ± 1.6	0.10
60.....	1087	66.0 ± 9.5	5.07 ± 0.73	2.31 ± 0.33	20.0 ± 7.1	0.87 ± 0.03	0.73 ± 0.24	2.17 ± 0.31	0.14 ± 0.12	18.8 ± 6.7	1.2 ± 1.1	0.06
61 <sup>e</sup> .....	1297	4 ± 11	0.42 ± 1.00	0.16 ± 0.37	1.2 ± 3.0	0.15 ± 0.02	0.12 ± 0.02	0.11 ± 0.30	0.05 ± 0.22	0.9 ± 2.4	0.4 ± 1.7	0.31
62.....	1861	37.6 ± 7.4	3.44 ± 0.68	1.73 ± 0.34	13.0 ± 5.2	0.69 ± 0.07	0.49 ± 0.18	1.26 ± 0.29	0.47 ± 0.18	9.5 ± 3.9	3.5 ± 1.8	0.27
63.....	543	18.1 ± 5.5	1.97 ± 0.60	0.96 ± 0.29	8.3 ± 3.5	0.15 ± 0.03	0.14 ± 0.05	0.96 ± 0.27	0.00 ± 0.11	8.3 ± 3.4	0.0 ± 0.9	0.00
64.....	1431	60.6 ± 9.3	6.72 ± 1.03	3.03 ± 0.46	27.5 ± 10.2	0.78 ± 0.04	0.70 ± 0.24	2.57 ± 0.41	0.45 ± 0.21	23.4 ± 8.7	4.1 ± 2.4	0.15
65.....	1528	40.7 ± 7.6	5.57 ± 1.04	2.46 ± 0.46	24.9 ± 8.7	0.66 ± 0.04	0.69 ± 0.21	2.24 ± 0.42	0.22 ± 0.19	22.7 ± 7.9	2.2 ± 2.0	0.09
66.....	1695	14.4 ± 5.7	1.45 ± 0.58	0.68 ± 0.27	8.5 ± 5.2	0.09 ± 0.03	0.10 ± 0.06	0.48 ± 0.22	0.20 ± 0.15	6.0 ± 3.9	2.5 ± 2.2	0.29
67.....	1833	18.1 ± 5.5	2.55 ± 0.77	1.26 ± 0.38	21.1 ± 13.5	0.15 ± 0.04	0.25 ± 0.16	0.87 ± 0.32	0.39 ± 0.22	14.6 ± 9.8	6.5 ± 5.2	0.31
68.....	437	42.1 ± 7.9	3.13 ± 0.59	1.40 ± 0.26	15.0 ± 7.2	0.46 ± 0.04	0.34 ± 0.15	1.32 ± 0.24	0.08 ± 0.10	14.1 ± 6.8	0.9 ± 1.2	0.06
69.....	2019	23.9 ± 6.1	2.72 ± 0.69	1.28 ± 0.33	23.5 ± 17.8	0.72 ± 0.05	1.19 ± 0.85	1.28 ± 0.31	0.00 ± 0.11	23.5 ± 17.6	0.0 ± 2.1	0.00
70.....	33	2.2 ± 4.2	0.61 ± 1.17	0.34 ± 0.66	5.1 ± 10.1	-0.13 ± 0.07	-0.21 ± 0.16	0.22 ± 0.53	0.12 ± 0.38	3.3 ± 8.1	1.8 ± 5.7	0.35

TABLE 2—Continued

Number (1)	VCC (2)	$N_{GC}$ (3)	$S_V$ (4)	$S_{N,z}$ (5)	$T$ (6)	$S_L$ (7)	$S_M$ (8)	$S_{N,z,blue}$ (9)	$S_{N,z,red}$ (10)	$T_{blue}$ (11)	$T_{red}$ (12)	$f_{red}$ (13)
71.....	200	15.5 ± 5.8	2.21 ± 0.83	1.18 ± 0.44	11.5 ± 6.0	0.05 ± 0.05	0.05 ± 0.06	1.18 ± 0.40	0.00 ± 0.19	11.5 ± 5.8	0.0 ± 1.8	0.00
72.....	571	10.9 ± 5.6	1.29 ± 0.66	0.53 ± 0.27	11.2 ± 9.4	0.14 ± 0.07	0.17 ± 0.20	0.53 ± 0.24	0.00 ± 0.13	11.2 ± 9.0	0.0 ± 2.7	0.00
73.....	21	20.7 ± 6.5	3.85 ± 1.21	2.32 ± 0.73	37.0 ± 27.2	0.67 ± 0.08	0.97 ± 0.65	2.32 ± 0.66	0.00 ± 0.30	37.0 ± 26.78	0.0 ± 4.8	0.00
74.....	1488	7.3 ± 4.2	1.42 ± 0.82	0.85 ± 0.49	17.9 ± 19.5	0.03 ± 0.04	0.08 ± 0.12	0.84 ± 0.44	0.01 ± 0.22	17.7 ± 18.8	0.2 ± 4.6	0.01
75.....	1779	2.1 ± 3.7	0.37 ± 0.64	0.23 ± 0.41	2.1 ± 3.7	-0.11 ± 0.05	-0.13 ± 0.08	0.23 ± 0.35	0.00 ± 0.22	2.1 ± 3.2	0.0 ± 2.0	0.00
76.....	1895	6.3 ± 4.3	1.45 ± 0.99	0.76 ± 0.52	8.3 ± 7.2	0.07 ± 0.06	0.05 ± 0.07	0.65 ± 0.44	0.11 ± 0.27	7.1 ± 6.1	1.3 ± 3.1	0.14
77.....	1499	19.0 ± 6.6	4.65 ± 1.62	3.14 ± 1.09	133.8 ± 185.1	0.98 ± 0.14	4.07 ± 5.47	2.67 ± 0.94	0.47 ± 0.55	113.7 ± 157.5	20.1 ± 35.7	0.15
78.....	1545	54.2 ± 8.8	9.37 ± 1.52	4.34 ± 0.71	38.5 ± 13.0	1.02 ± 0.07	0.76 ± 0.23	3.91 ± 0.64	0.43 ± 0.29	34.6 ± 11.8	3.8 ± 2.8	0.10
79 <sup>c</sup> .....	1192	-6 ± 13	-1.12 ± 2.42	-0.34 ± 0.74	-3.4 ± 7.3	-0.39 ± 0.05	-0.31 ± 0.11	-0.21 ± 0.61	-0.13 ± 0.45	-2.1 ± 5.9	-1.3 ± 4.4	0.00
80.....	1857	10.8 ± 5.9	2.39 ± 1.31	1.28 ± 0.70	14.8 ± 8.6	0.32 ± 0.09	0.29 ± 0.12	0.78 ± 0.56	0.50 ± 0.42	9.1 ± 6.7	5.8 ± 5.0	0.39
81.....	1075	16.5 ± 5.2	3.19 ± 1.01	1.58 ± 0.50	14.8 ± 7.8	0.33 ± 0.06	0.31 ± 0.14	1.58 ± 0.46	0.00 ± 0.18	14.8 ± 7.6	0.0 ± 1.7	0.00
82.....	1948	2.5 ± 3.0	0.94 ± 1.13	0.48 ± 0.58	5.0 ± 6.1	-0.28 ± 0.17	-0.22 ± 0.12	0.48 ± 0.50	0.00 ± 0.29	5.0 ± 5.3	0.0 ± 3.0	0.00
83.....	1627	3.6 ± 3.7	0.94 ± 0.97	0.40 ± 0.41	3.5 ± 3.8	1.21 ± 0.13	1.05 ± 0.34	0.40 ± 0.36	0.00 ± 0.21	3.5 ± 3.3	0.0 ± 1.8	0.00
84.....	1440	26.7 ± 6.8	4.81 ± 1.22	2.30 ± 0.58	22.5 ± 10.1	0.29 ± 0.05	0.31 ± 0.12	1.93 ± 0.51	0.37 ± 0.28	18.9 ± 8.6	3.6 ± 3.1	0.16
85.....	230	28.7 ± 6.7	9.39 ± 2.19	4.73 ± 1.10	41.4 ± 20.0	5.45 ± 1.17	3.68 ± 1.73	4.64 ± 1.03	0.09 ± 0.41	40.5 ± 19.4	0.8 ± 3.6	0.02
86.....	2050	9.2 ± 4.3	2.63 ± 1.23	1.27 ± 0.60	31.5 ± 32.0	0.21 ± 0.06	0.43 ± 0.41	1.27 ± 0.54	0.00 ± 0.25	31.5 ± 31.4	0.0 ± 6.1	0.00
87.....	1993	-1.6 ± 2.1	-0.48 ± 0.63	-0.25 ± 0.33	-2.1 ± 2.9	-0.07 ± 0.08	0.29 ± 0.13	-0.03 ± 0.30	-0.22 ± 0.29	-0.3 ± 2.6	-1.9 ± 2.6	0.00
88.....	751	9.2 ± 4.3	1.49 ± 0.70	0.62 ± 0.29	6.5 ± 3.9	0.13 ± 0.04	0.11 ± 0.06	0.62 ± 0.27	0.00 ± 0.12	6.5 ± 3.6	0.0 ± 1.3	0.00
89.....	1828	20.4 ± 5.8	3.94 ± 1.12	2.10 ± 0.60	16.3 ± 7.7	0.95 ± 0.08	0.78 ± 0.30	2.10 ± 0.56	0.00 ± 0.22	16.3 ± 7.5	0.0 ± 1.7	0.00
90.....	538	0.6 ± 3.5	0.19 ± 1.11	0.09 ± 0.51	0.9 ± 5.5	-0.31 ± 0.22	-0.25 ± 0.19	0.09 ± 0.42	0.00 ± 0.29	0.9 ± 4.6	0.0 ± 3.1	0.00
91.....	1407	49.7 ± 8.6	10.18 ± 1.76	5.30 ± 0.92	40.2 ± 15.5	0.69 ± 0.07	0.50 ± 0.18	4.61 ± 0.82	0.69 ± 0.41	35.0 ± 13.5	5.2 ± 3.6	0.13
92.....	1886	3.9 ± 2.6	1.24 ± 0.82	0.66 ± 0.44	7.7 ± 5.3	0.29 ± 0.06	0.41 ± 0.11	0.66 ± 0.41	0.00 ± 0.17	7.7 ± 5.0	0.0 ± 1.9	0.00
93 <sup>c</sup> .....	1199	-9 ± 14	-6.24 ± 9.23	-1.60 ± 2.37	-16.6 ± 24.9	-3.39 ± 0.24	-2.55 ± 0.74	-0.90 ± 1.93	-0.70 ± 1.45	-9.3 ± 20.1	-7.3 ± 15.2	0.00
94.....	1743	9.8 ± 6.4	2.88 ± 1.88	1.63 ± 1.06	24.6 ± 23.9	-0.03 ± 0.29	-0.13 ± 0.55	1.11 ± 0.87	0.52 ± 0.62	16.7 ± 17.8	7.9 ± 10.9	0.32
95.....	1539	31.0 ± 7.0	11.83 ± 2.67	4.40 ± 0.99	59.4 ± 18.7	2.41 ± 0.33	1.73 ± 0.40	4.40 ± 0.93	0.00 ± 0.35	59.4 ± 18.1	0.0 ± 4.7	0.00
96 <sup>c</sup> .....	1185	14.0 ± 5.7	2.73 ± 1.11	1.58 ± 1.11	15.7 ± 11.0	0.20 ± 0.04	0.26 ± 0.16	1.45 ± 0.56	0.13 ± 0.31	14.4 ± 9.9	1.3 ± 3.1	0.08
97.....	1826	4.5 ± 3.9	1.74 ± 1.51	0.93 ± 0.81	7.3 ± 6.9	0.08 ± 0.11	-0.01 ± 0.12	0.93 ± 0.71	0.00 ± 0.40	7.3 ± 6.1	0.0 ± 3.1	0.00
98.....	1512	4.7 ± 3.8	1.48 ± 1.20	1.14 ± 0.92	7.7 ± 6.5	-0.40 ± 0.58	-0.42 ± 0.42	1.14 ± 0.81	0.00 ± 0.44	7.7 ± 5.7	0.0 ± 3.0	0.00
99.....	1489	11.7 ± 4.8	6.69 ± 2.74	3.19 ± 1.31	36.9 ± 16.7	0.46 ± 0.16	0.34 ± 0.14	2.20 ± 1.08	0.99 ± 0.75	25.4 ± 13.3	11.4 ± 8.9	0.31
100.....	1661	10.2 ± 4.7	4.83 ± 2.23	1.12 ± 0.51	23.9 ± 12.0	0.08 ± 0.08	0.06 ± 0.14	1.12 ± 0.46	0.00 ± 0.22	23.9 ± 11.0	0.0 ± 4.7	0.00

NOTES.—Col. (1): Running number, sorted by increasing  $B_r$  magnitude. Col. (2): Number in Virgo Cluster Catalog. Col. (3): Total number of GCs. Col. (4): Specific frequency. Col. (5): Specific frequency in  $z$  bandpass. Col. (6):  $N_{GC}$  normalized to stellar mass of  $10^9 M_{\odot}$ . Col. (7): Percentage of galaxy  $z$  luminosity in GCs. Col. (8): Percentage of galaxy stellar mass in GCs. Col. (9): Specific frequency in  $z$  for blue GCs. Col. (10): Specific frequency in  $z$  for red GCs. Col. (11):  $T$  for blue GCs. Col. (12):  $T$  for red GCs. Col. (13): Fraction of red GCs.

<sup>a</sup> Mass-to-light ratio (i.e., optical-IR colors) suspect.

<sup>b</sup> Quantities possibly affected by nearby dE VCC 1941, which has been masked.

<sup>c</sup> Quantities calculated using an  $R = 70''$  aperture around galaxy.

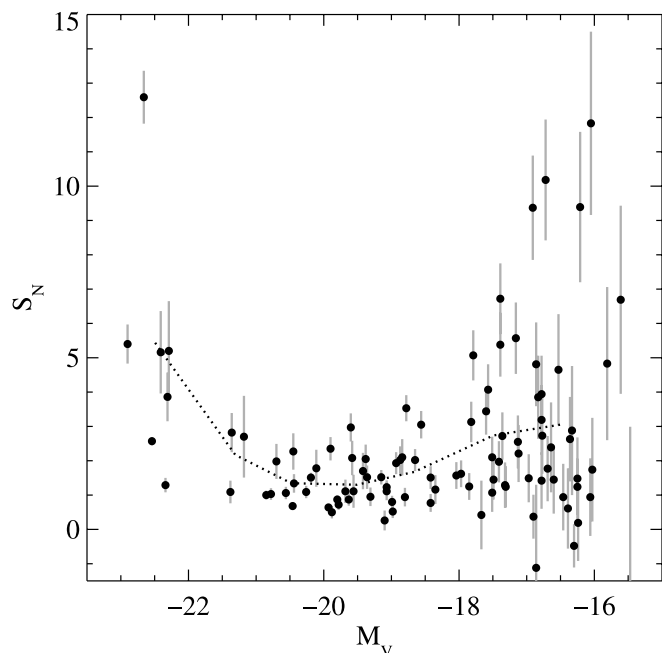


FIG. 2.— $S_N$  vs. galaxy  $M_V$  for 100 ACSVCS galaxies. M87, the giant elliptical with  $S_N \sim 13$ , is a well-known outlier. Luminous early-type galaxies have higher  $S_N$  than intermediate-luminosity early-type galaxies by a factor of  $\sim 2$ – $3$ . Galaxies with intermediate luminosities ( $-20.5 < M_V < -18$ ) generally have  $S_N \sim 1.5$ . Early-type dwarf galaxies have a large spread in  $S_N$ , with some having zero GCs, and others having among the highest measured  $S_N$  in our sample. The dotted line shows the mean trend (including M87), whose values are listed in Table 3.

Zepf et al. (1995), and Harris et al. (2004). At the faint end, we include data from three sources that extend the range in galaxy luminosity by many magnitudes. We include the seven dEs studied by Durrell et al. (1996a, 1996b) that do not overlap with our sample, the Virgo and Fornax dEs in the *HST* WFPC2 study of Miller et al. (1998) and Lotz et al. (2004), and the five Local Group dwarfs listed in Lotz et al. (2004).

At the high-luminosity end, the literature values have large scatter, possibly due to their heterogeneous nature. One of the benefits of the ACSVCS is that relative distances between Virgo Cluster galaxies are both small and measured in a homogeneous fashion. At low luminosities, the studies are mainly in Virgo and Fornax so relative distances are less of a problem. At luminosities fainter than the ACSVCS, the trend of higher  $S_N$  coupled with a larger range of  $S_N$  continues to the limit of the data. The ACSVCS sample not only has smaller errors, it fills the important luminosity regime between giants and dwarfs. Kundu & Whitmore (2001a, 2001b) studied galaxies in the magnitude range  $-23 < M_B < -16$ ,

TABLE 3  
SPECIFIC FREQUENCY IN BINS OF  $M_V$

$M_V$ Range	$\langle M_V \rangle$	$S_N$
$(-24, -22)^a$ .....	-22.5	4.0
$(-24, -22)$ .....	-22.5	5.4
$(-22, -21)$ .....	-21.3	2.2
$(-21, -20)$ .....	-20.5	1.3
$(-20, -19)$ .....	-19.5	1.3
$(-19, -18)$ .....	-18.7	1.7
$(-18, -17)$ .....	-17.5	2.7
$(-17, -15)$ .....	-16.4	3.1

<sup>a</sup> Not including VCC 1316 (M87).

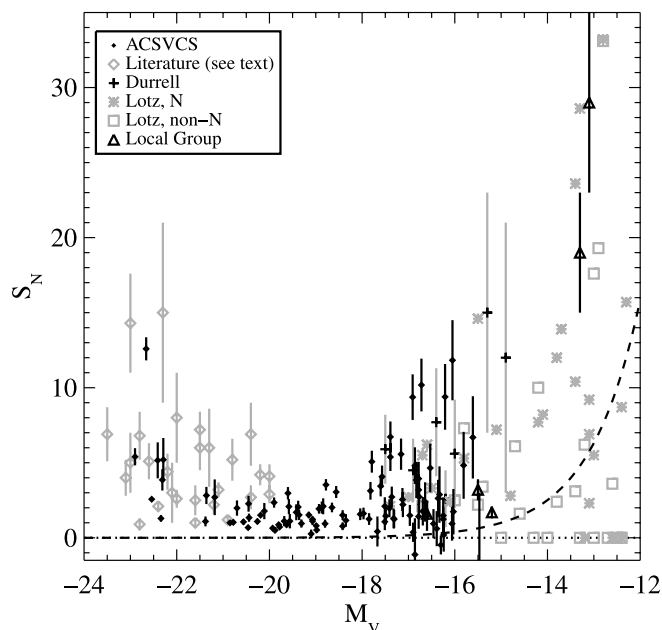


FIG. 3.— $S_N$  vs. galaxy  $M_V$  for 100 ACSVCS galaxies (black filled diamonds), and early-type galaxy data from the literature (references in the text). Literature values for  $S_N$  follow and extend trends visible in the ACSVCS galaxies. In particular, dwarf galaxy  $S_N$  values from the work of Durrell et al. (1996a) and Lotz et al. (2004) show that fainter dwarf ellipticals can have an even larger range of specific frequency. For clarity, error bars are not plotted for the *HST* WFPC2 dEs, but the uncertainty can be very large. The dashed line shows the  $S_N$  value if a galaxy at that magnitude had one GC. [See the electronic edition of the *Journal* for a color version of this figure.]

similar to the ACSVCS, and found that ellipticals generally had higher  $S_N$  than lenticulars, but they were only able to measure “local”  $S_N$  within the WFPC2 field of view.

Ultimately, studies of GC specific frequency are about how GC formation scales with galaxy mass, using  $M_V$  as a proxy. However, a redder bandpass more faithfully traces stellar mass. In Figure 4 and Table 4 we introduce  $S_{N,z}$ , which is defined identically to  $S_N$  except that the number of GCs is normalized to an absolute  $z$  magnitude of  $M_z = -15$  and plotted against  $M_z$ .  $S_{N,z}$  values are 1.5–3 times smaller than  $S_N$  because early-type galaxies are red (and thus more luminous in  $z$ ). The specific frequencies of giant ellipticals are adjusted more than those of the dwarfs because massive galaxies are redder in color. While the trends in this figure are the same as those in Figure 2, one interesting difference is that the highest  $S_{N,z}$  values for the dwarfs now equal or exceed that of M87, with the change being due to M87’s redder color. For the rest of the paper we use quantities based on  $M_z$ .

#### 4.2. $S_{N,z}$ and Bulge Luminosity

In most scenarios of GC system formation, the GCs are associated with the stellar spheroids (either halo or bulge) and not with the formation of present-day disks. If this is the case, then perhaps normalizing GC numbers to total spheroid luminosity would be more fundamental than normalizing to total luminosity. In our ACSVCS sample, over half of our galaxies are morphologically classified in the VCC as either E/S0, S0/E, S0, or dS0. Most of these galaxies are at intermediate luminosities and make up a substantial fraction of the galaxies that have low  $S_{N,z}$ . Could it be the case that these low  $S_{N,z}$  values are due to the inclusion of a stellar disk in the total luminosity? Moreover, the scatter in the  $S_{N,z}$  of the intermediate-luminosity galaxies in our sample exceeds the scatter expected from the errors, implying that a parameter such as bulge fraction might be important.

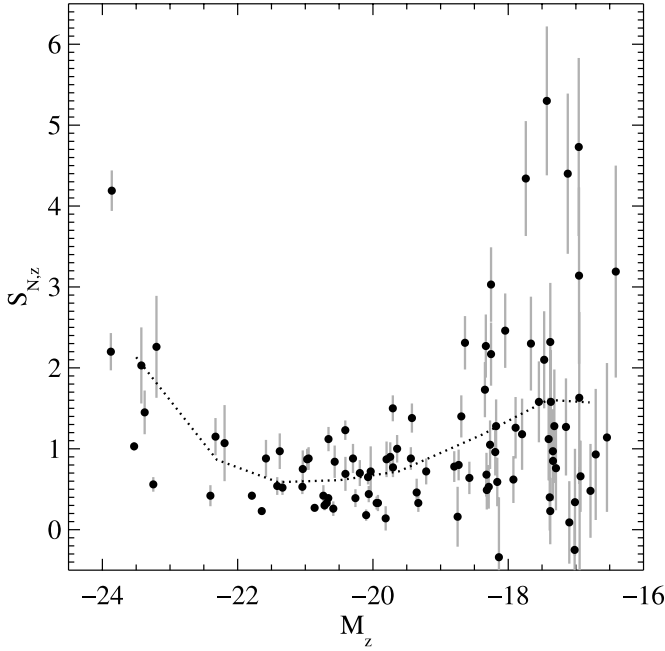


FIG. 4.— $S_{N,z}$  vs. galaxy  $M_z$ . This figure shows similar quantities to those in Fig. 2, except the number of GCs is normalized to the galaxy absolute magnitude in the  $z$  band, with  $M_z$  plotted along the  $x$ -axis. This far-red bandpass is a better tracer of the total stellar mass than the traditionally used  $V$  band. The dotted line shows the mean trend with values in Table 4.

We can test whether normalizing by bulge luminosity is more fundamental with quantitative bulge-disk decompositions (C.-W. Chen et al. 2008, in preparation). In this paper, however, we use the morphological T types given in the Third Reference Catalog of Bright Galaxies (RC3; de Vaucouleurs et al. 1991). These T types are correlated with the ratio of bulge to total luminosity using the relation of Simien & de Vaucouleurs (1986). We can do this for 55 galaxies in the ACSVCS sample, and we show how their  $S_{N,z}$  values change in Figure 5. This figure shows  $S_{N,z}$  as derived using total luminosity compared to  $S_{N,z}$  as derived from bulge luminosity alone with the two values connected by an arrow. For some galaxies, their specific frequencies do increase to values typical of the luminous ellipticals ( $S_{N,z} \sim 1-2$ ), but for others whose disks are less prominent, their  $S_{N,z}$  stay relatively unchanged. For two galaxies, their RC3 morphological type values imply that they have little or no bulge, which give them very high  $S_{N,z}$ .

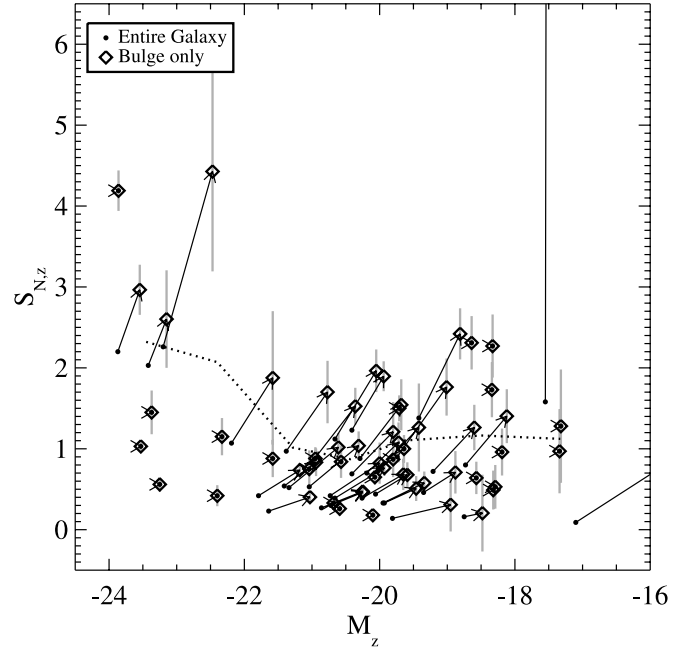


FIG. 5.— $S_{N,z}$  vs. galaxy  $M_z$  for ACSVCS galaxies with VCC morphological types of E/S0, S0/E, S0, and dS0, and which have morphological T types in the RC3. We plot the total  $S_{N,z}$  (same as Fig. 4), as well as an arrow connecting to the galaxy's  $S_{N,z}$  normalized to bulge luminosity alone. Some galaxies have  $S_{N,z}$  similar to those in luminous ellipticals, but others are relatively unchanged.

Based on the current morphological classifications, it is not apparent that GCs form at a constant efficiency with respect to bulge luminosity. Normalizing to bulge luminosity increases the scatter in  $S_{N,z}$  in the intermediate-luminosity range. However, a more thorough analysis based on modern bulge-to-disk decompositions is necessary before we can reach any strong conclusions.

#### 4.3. Normalizing to Stellar Mass

Using the stellar masses for the ACSVCS galaxies, we calculate the  $T$  parameter introduced by Zepf & Ashman (1993), which is the number of GCs ( $N_{GC}$ ) per  $10^9 M_{\odot}$ ,

$$T = N_{GC} / (\mathcal{M}_{G*} / 10^9 M_{\odot}), \quad (2)$$

where  $\mathcal{M}_{G*}$  is the stellar mass of the galaxy. The advantage of using  $T$  instead of  $S_N$  is that it allows comparisons across galaxies

TABLE 4  
BINS OF  $M_z$

$M_z$ Range	$\langle M_z \rangle$	$S_{N,z}$	$T$	$S_L$	$S_M$	$S_{N,z,blue}$	$S_{N,z,red}$	$f_{red}$	$S_{N,z,close}^a$	$S_{N,z,far}^b$
(-25, -23) <sup>c</sup> .....	-23.4	1.6	13.2	0.90	0.74	1.17	0.47	0.29	1.72	1.57
(-25, -23) .....	-23.5	2.1	18.2	1.12	1.04	1.53	0.60	0.28	2.55	1.57
(-23, -22) .....	-22.3	0.9	8.5	0.51	0.50	0.49	0.37	0.43	1.11	0.43
(-22, -21) .....	-21.4	0.6	4.1	0.21	0.17	0.40	0.19	0.32	0.65	0.45
(-21, -20) .....	-20.5	0.6	4.7	0.18	0.16	0.43	0.19	0.31	0.57	0.68
(-20, -19) .....	-19.6	0.7	5.8	0.21	0.17	0.57	0.15	0.21	0.75	0.67
(-19, -18) .....	-18.4	1.2	11.4	0.38	0.34	0.99	0.18	0.16	1.41	0.80
(-18, -17) .....	-17.4	1.6	17.8	0.36	0.37	1.46	0.15	0.09	2.04	1.02
(-17, -16) .....	-16.8	1.6	17.3	0.51	0.35	1.36	0.21	0.13	2.03	...

<sup>a</sup> Within 1 Mpc of VCC 1316 (M87), in projection.

<sup>b</sup> Outside 1 Mpc of VCC 1316 (M87), in projection.

<sup>c</sup> Not including VCC 1316 (M87).

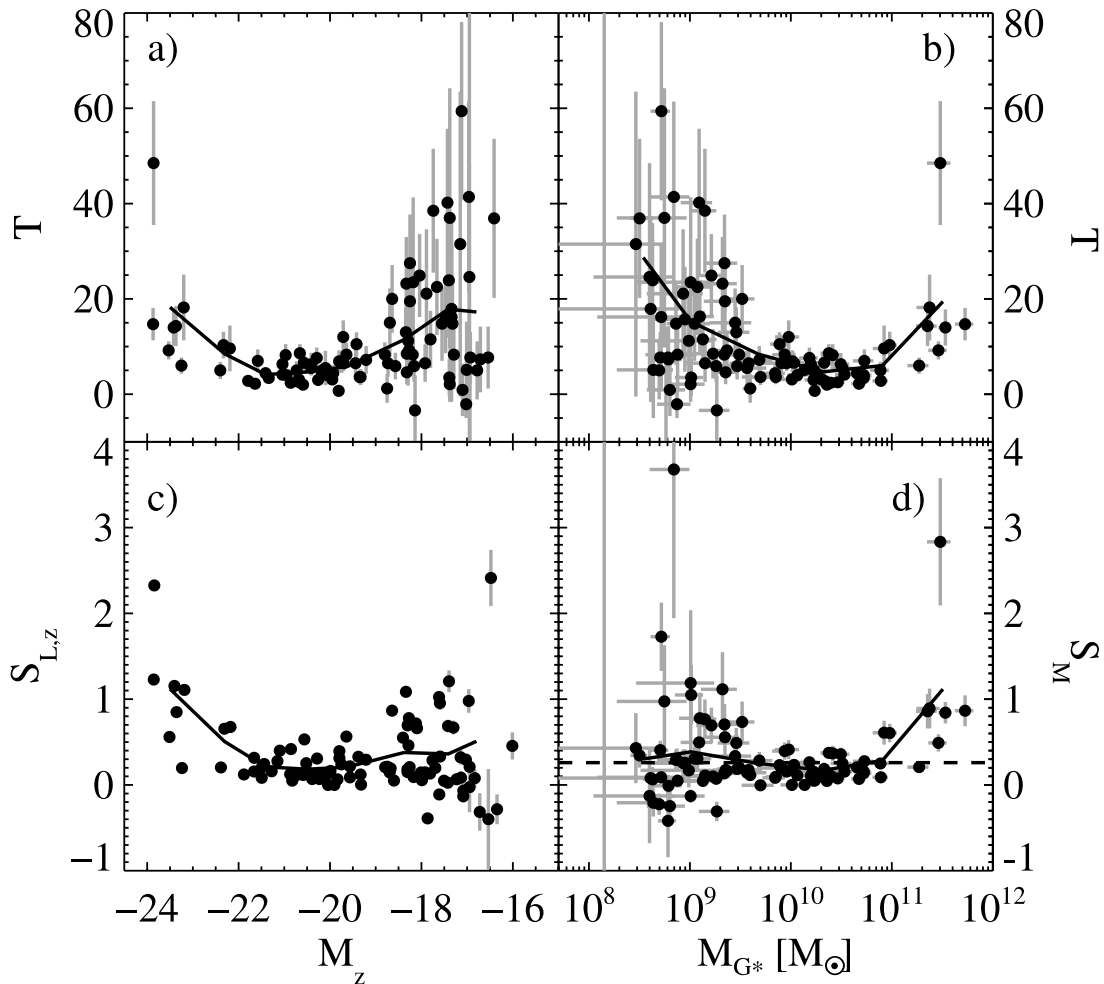


FIG. 6.— $T_N$  vs. (a)  $M_z$  and (b) galaxy stellar mass.  $T_N$  is the number of GCs per  $10^9 M_\odot$ . In (c) and (d) we show the specific  $z$  luminosity,  $S_{L,z}$ , vs.  $M_z$  and the specific mass,  $S_M$ , vs. stellar mass. The dashed line in (d) marks  $S_M = 0.26$ , the “universal” GC formation efficiency from McLaughlin (1999a). Because the ACSVCS sample is all early-type galaxies,  $M/L$  does not vary much, and hence in all four panels we see trends similar to those for  $S_N$  and  $S_{N,z}$  (see text for exceptions). The difference between dwarfs and giants in  $S_L$  and  $S_M$  is not as large as it is in  $S_N$  or  $T_N$ . This reflects the changing GCLF across galaxy luminosity (see Fig. 7). In all figures, the solid line shows the mean trend; values in Table 5.

with different mass-to-light ratios. In Figures 6a and 6b (and corresponding Table 5) we show  $T$  plotted against  $M_z$  and  $M_{G*}$ . Although the errors are larger than in Figure 4, the same trends are evident. Our values for  $T$  are higher than those in previous studies, such as Rhode & Zepf (2004). This is due to differences in the mass-to-light ratios used. Previous studies have assumed  $M/L_V = 10$  for elliptical galaxies, whereas we have estimated  $M/L$  using galaxy colors and the BC03 models with a Chabrier

TABLE 5  
BINS OF  $M_{G*}$

$M_*$ Range <sup>a</sup>	$\langle M_* \rangle$	$T$	$S_L$	$S_M$
(0.1, 0.5).....	0.3	18.2	1.12	0.30
(0.5, 2.2).....	1.0	8.5	0.51	0.39
(2.2, 10.0).....	4.8	4.1	0.21	0.25
(10.0, 46.4).....	21.1	4.7	0.18	0.17
(46.4, 215.4).....	80.8	5.8	0.21	0.29
(215.4, 1000.0).....	321.0	11.4	0.38	1.12
(215.4, 1000.0) <sup>b</sup> .....	324.7	14.0	1.00	0.80

<sup>a</sup>  $M_*/10^9 M_\odot$ .

<sup>b</sup> Not including VCC 1316 (M87).

(2003) IMF. Even accounting for the different bandpasses used, our  $M/L$  are systematically lower. For example, for M49, the most luminous elliptical in the sample,  $(V - z) = 0.97$  and its  $M/L_z = 2.0$  translates to  $M/L_V = 5$ , which is a factor of 2 lower than the canonical Zepf & Ashman (1993) values. We feel that although the  $M/L$  values in Zepf & Ashman (1993) have been valuable as standard conversions, they are too high, and that our values of  $M/L$  are more reasonable given recent dynamical mass measurements of elliptical galaxies (Kronawitter et al. 2000; Cappellari et al. 2006). The use of a Salpeter (1955) IMF increases  $M/L$  by a factor of 1.8 but is not observationally motivated for stars below  $1 M_\odot$  and would produce stellar  $M/L$  values higher than the total  $M/L$  measured in some early-type galaxies. For the early-type galaxies in our sample,  $M/L_V$  ranges from 1 to 5, although some galaxies with obvious star formation have  $M/L_V < 1$ . We reiterate, however, that using a redder bandpass is better, and all our masses are derived using  $M/L_z$ ; we discuss the  $V$  bandpass only for comparing with previous work.

For the purposes of this paper, however, the absolute scale of  $T$  is not as important as the relative scale within the sample. Figure 6b shows that  $T$  spans a wide range at masses  $M_{G*} < 4 \times 10^9 M_\odot$ , stays constant until  $10^{11} M_\odot$ , and then increases again at higher mass.

#### 4.4. Specific Luminosity and Mass

If we want to know what fraction of a galaxy’s luminosity (or mass) is in GCs, we can total it up directly, as described in § 3.2. We use the definition of specific luminosity,  $S_L$ , presented in Harris (1991) except that we use the  $z$  bandpass instead of  $V$ :

$$S_{L,z} = 100L_{GC,z}/L_{\text{galaxy},z}, \quad (3)$$

where  $L_{GC,z}$  and  $L_{\text{galaxy},z}$  are the total  $z$ -band luminosities of the GCs and the galaxy, respectively. This quantity has two advantages over  $S_N$  in that it is independent of distance and the completeness corrections for unobserved faint GCs are extremely small. We also plot the specific mass,  $S_M$ , defined as

$$S_M = 100M_{GC}/M_{G*}, \quad (4)$$

where  $M_{GC}$  is the total stellar mass in GCs, calculated as described in § 3.2.

Figure 6c shows the specific luminosity  $S_L$  of GCs plotted against  $M_z$ , and Figure 6d shows  $S_M$  versus  $M_{G*}$ . The dashed line marks the 0.26% “universal” GC formation efficiency from McLaughlin (1999a) and is a reasonable description of the GC mass fraction in intermediate-mass galaxies. For the ACSVCS galaxies with  $-22 < M_z < -19$ ,  $\langle S_L \rangle = 0.20$  and  $\sigma_{S_L} = 0.14$ . Similarly, for galaxies with  $M_{G*} = (0.4-6) \times 10^{10} M_\odot$ ,  $\langle S_M \rangle = 0.17$  and  $\sigma_{S_M} = 0.11$ . That the mass fraction is slightly lower than the luminosity fraction can be explained by the bluer colors and hence lower  $M/L$  of the GCs as compared to their host galaxies.

The major difference in Figure 6 between the top two panels showing  $T$  and the bottom two showing  $S_{L,z}$  and  $S_M$  is the comparison between the giants and dwarfs. While the dwarfs have very high number fractions, their luminosity and mass fractions are substantially lower compared to the giants. Why should this be the case? The reason for this has to do with the GCLF. If the GCLF was constant across all galaxies (i.e., same mean  $\mu$  and width  $\sigma$ ), then  $S_L$  would mirror  $T$ . However, Jordán et al. (2006) and Paper XII showed that the GCLF varies as a function of galaxy mass in the sense that less massive galaxies host GC systems with fainter  $\mu$  and smaller  $\sigma$ . This has the effect of lowering the mean luminosity (and mass) of GCs in the dwarf galaxies. We illustrate this in Figure 7, where we plot the mean  $z$  luminosity of GCs in each galaxy,  $\langle L_{GC,z} \rangle$ , against galaxy luminosity (in this case  $M_B$  because that is how we presented the data in Jordán et al.). It is clear that a constant GCLF does not describe the data well. The dotted line shows the effect of varying  $\mu$ , and the dashed line shows the relation predicted by the combination of fainter  $\mu$  and a narrower  $\sigma$  in the dwarfs. Although the scatter is large, the expected change in the GCLF can account for the changes in mean GC luminosity that we see. This trend is actually more pronounced in mass than in luminosity because GCs in dwarfs are more metal-poor and have lower  $M/L$  than those in giants.

The result is that some of the more extreme  $S_N$  or  $T$  values seen in the dwarfs are somewhat less extreme when expressed as a mass fraction. Nevertheless, many low-luminosity galaxies still have  $S_L$  and  $S_M$  that are significantly higher than those in the intermediate- $L$  galaxies. Because the global trends are similar in  $S_N$ ,  $T$ ,  $S_L$ , and  $S_M$ , we refer to these quantities collectively as “GC fractions.”

#### 4.5. Specific Frequencies of Red and Blue GCs

Our previous study of GC color and metallicity distributions in the ACSVCS galaxies shows that they are, on average, either

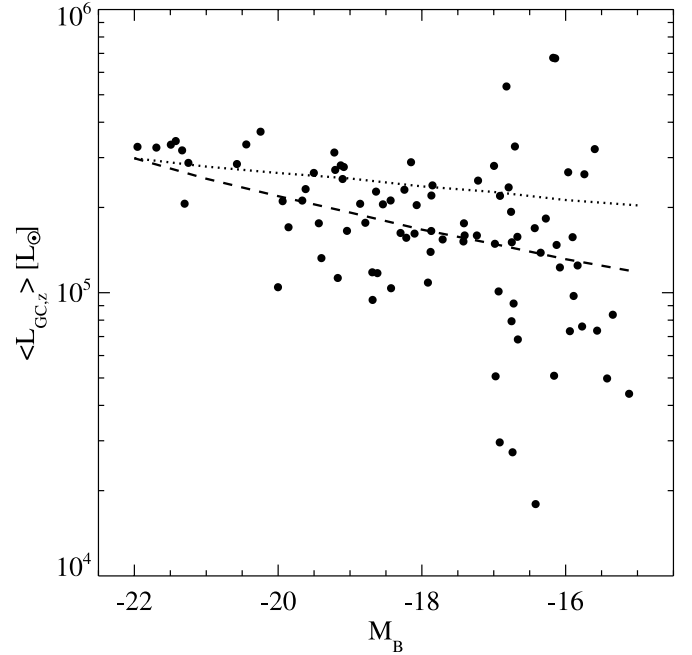


FIG. 7.—Mean  $z$  luminosity of GCs,  $\langle L_{GC,z} \rangle$ , in ACSVCS galaxies vs.  $M_z$ . If the GCLF was constant across all galaxies,  $\langle L_{GC,z} \rangle$  would also be constant. The effect of a fainter GCLF turnover in dwarf galaxies (dotted line) only partially explains the lower mean luminosities. A combination of fainter turnovers and narrower GCLFs (dashed line) reproduces the observed trend.

bimodal or asymmetric across the entire luminosity range of the sample, and we use the products of that analysis (Paper IX) in the current study. The blue (metal-poor) and red (metal-rich) are believed to trace either different epochs of formation or different progenitor halos (however, see Yoon et al. [2006] and Cantiello & Blakeslee [2007] for arguments that the bimodality may be an observational consequence of a nonlinear metallicity-color relation). It is an interesting question to ask how the specific frequencies of each GC population scale with galaxy properties. For example, Rhode & Zepf (2004) showed that the mass-normalized number of blue GCs,  $T_{\text{blue}}$ , increases as one goes from spirals to elliptical galaxies, although their spirals were less massive than ellipticals so the trend could also have been one in galaxy mass. Assuming that mergers of spirals only produce new red GCs, they argued that the GC systems of ellipticals cannot be formed purely by spiral mergers (see Ashman & Zepf 1992).

We determine the fraction of blue and red GCs in a hybrid approach similar to what we do for total numbers. For the brightest 21 galaxies in the sample (a complete magnitude-limited sample, as ranked by  $B_T$  in Paper I) there are sufficient numbers of GCs that the KMM two-Gaussian fits to the  $(g-z)$  distribution performed in Paper IX are reliable enough that we can use the “dip” between the two Gaussians (the color at which a GC is equally likely to belong to the blue or red GC distribution) as the dividing color between the two populations. For the remaining galaxies, we assume a fixed dividing color of  $(g-z) = 1.16$ . Although the colors of the individual peaks vary as a function of galaxy luminosity, the dip color is relatively invariant (see Fig. 5 in Paper IX).

The brighter galaxies may have better defined color distributions, but we only observe a fraction of their entire GC system, and the red-to-blue ratio is observed to decrease as a function of galactocentric radius. We correct for this bias in the brightest 14 galaxies by fitting Sérsic models to the surface density profiles of the red GCs separately and integrating the best model to obtain their total numbers. We fit to profiles derived from our ACS and

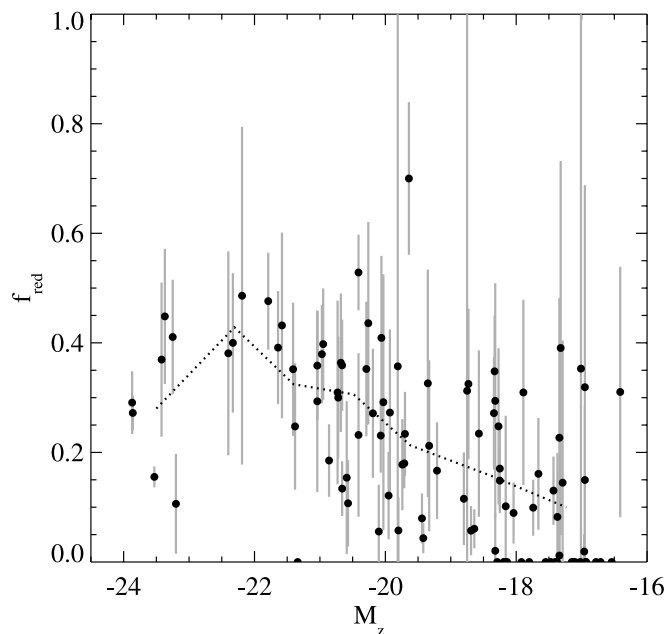


FIG. 8.—Fraction of red GCs ( $f_{\text{red}}$ ) vs.  $M_z$ . The dotted line represents the mean trend (Table 4). More luminous galaxies have higher fractions of red GCs up to  $M_z \sim -22$ , but at higher luminosities there is a flattening or turnover in  $f_{\text{red}}$ , and the most luminous galaxies do not have the highest red GC fractions.

WFPC2 data, as described in § 3.1. In the most luminous galaxies such as M49 and M87, the fraction of red GCs within the ACS WFC is as high as 0.6, but the red fraction of the entire GC system is more like 0.3. For the remainder of the galaxies, we find that the ACS WFC encompasses nearly all of the red GC system and use our corrected counts to determine the red GC fraction.

Figure 8 shows the fraction of red GCs as a function of  $M_z$ . The fraction of red GCs,  $f_{\text{red}}$ , generally increases with galaxy luminosity, going from  $\sim 0.1$  to  $\sim 0.5$ , but at  $M_z \lesssim -22$  the trend in  $f_{\text{red}}$  appears to either flatten or perhaps turn over. This is in contrast to the results of Paper IX, which only quantified the red GC fraction within the ACS WFC images and thus were biased to detecting the more centrally concentrated red GCs. With the proper aperture corrections, we can see that the most luminous galaxies in the sample do not have increasingly higher fractions of red GCs. Is this because there are fewer red GCs in these galaxies or more blue GCs?

Figure 9 shows the  $S_{N,z}$  for blue and red GCs as a function of  $M_z$ . The trend in the blue GCs mirrors the overall trend seen in Figure 4, with the massive and dwarf galaxies having the highest  $S_{N,z}$ . This is not surprising on the faint end since most of the GCs in the fainter galaxies are blue. Even for the massive galaxies, the high  $S_{N,z}$  values are dominated by blue GCs. However, the specific frequencies of red GCs also exhibit an increase for  $M_z < -21$ , especially for the cD galaxy M87, which is a  $4\sigma$  outlier. The elevated  $S_{N,z,\text{blue}}$  and  $S_{N,z,\text{red}}$  for the most massive galaxies suggest that massive galaxies are not underproducing red GCs, which is one possible interpretation of Figure 8. Instead, these massive galaxies have more red GCs, but even more blue GCs.

#### 4.6. $S_{N,z}$ and Nucleation

Observations by the *HST* WFPC2 of dEs in the Virgo and Fornax Clusters (Miller et al. 1998; Lotz et al. 2004; Miller & Lotz 2007) have found that dEs with stellar nuclei have a higher mean specific frequency, possibly implying that a higher past star formation efficiency resulted in the formation of both nuclei and GCs. The dEs in our sample are more luminous than the ones

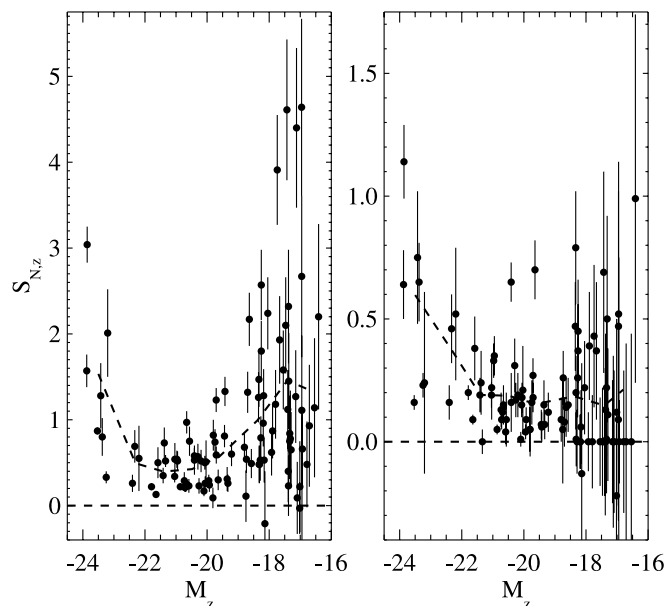


FIG. 9.— $S_{N,z}$  vs.  $M_z$  for blue (left) and red (right) globular clusters. Both red and blue GCs show enhanced  $S_{N,z}$  in massive galaxies, but the variation in  $S_{N,z}$  across galaxy mass is dominated by the blue GCs. This is true even in the massive galaxies. The dashed lines show the mean trends with values listed in Table 4. [See the electronic edition of the *Journal* for a color version of this figure.]

studied in the WFPC2 snapshot survey, and almost all of the galaxies are nucleated (Paper VIII). Therefore, it is difficult to test if there is any correlation between  $S_{N,z}$  and nucleation. Only four dwarf galaxies are definitely nonnucleated (Type II in Table 1 of Paper VIII): VCC 1049, VCC 1833, VCC 1499, and VCC 1512. Of these galaxies, VCC 1499 and VCC 1512 are dE/dI transition objects with young stars, and VCC 1049 has bluer colors toward the center. The  $S_{N,z}$  values of these galaxies range from 0.97 (VCC 1049) to 3.14 (VCC 1499) and do not appear different from the rest of the sample. When normalized to stellar mass, however, the young stellar populations have lower  $\mathcal{M}/L$  and thus much higher  $T$  and  $S_{\mathcal{M}}$  values.

#### 4.7. $S_{N,z}$ and Environment

One of the more intriguing questions presented by plots such as Figure 4 is the nature of  $S_{N,z}$  in the low-luminosity galaxies ( $M_z > -19$ ), which we loosely refer to as “dwarfs,” although some of the galaxies in our sample are on the more massive end of the spectrum of dwarf galaxies. Our relatively small errors show that the large spread in specific frequency for dwarf galaxies is not simply due to the expected observational scatter. Given our errors, the observed distribution of  $S_{N,z}$  is 2.4 times broader than we would expect if all dwarfs had a single  $S_{N,z}$ . This implies that there is at least one other parameter besides galaxy mass that governs the GC formation efficiency.

In this section we investigate the relationship between specific frequency and the galaxy’s environment, specifically, its distance from the cluster center, which is taken to be the location of the cD galaxy, M87 (Binggeli et al. 1985). Environment clearly plays a role in galaxy evolution, as is evidenced by the morphology-density relation, and this can be through gravitational and hydrodynamic processes such as tidal and ram pressure stripping, or through initial conditions where halos in denser regions collapse earlier. Dwarf galaxies, being the most vulnerable, are more likely to express the effects of their environment.

Figure 10 shows  $S_{N,z}$  for dwarf galaxies against their projected distance from the cluster center ( $R_p$ ), ignoring VCC 571 and VCC



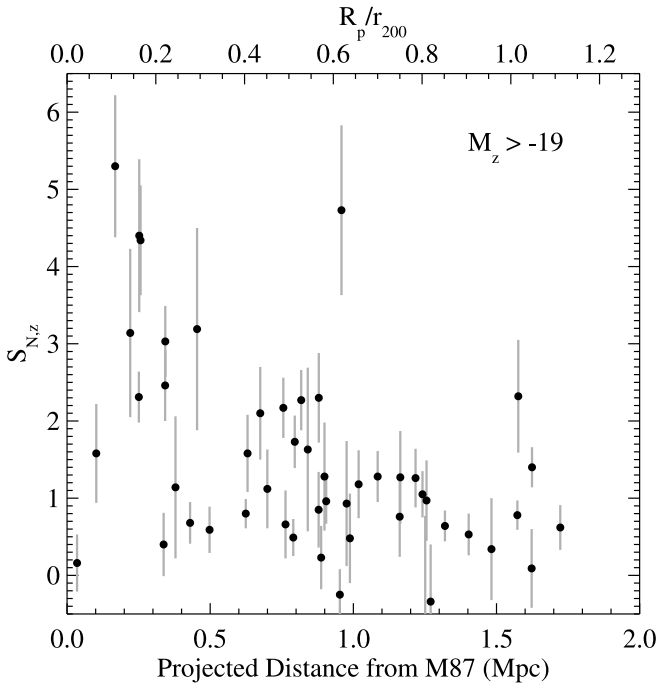


FIG. 10.— $S_{N,z}$  vs. projected clustercentric distance for low-luminosity galaxies ( $M_z > -19$ ), excluding VCC 571 and VCC 538, which are known to be  $\sim 6$ – $7$  Mpc behind the Virgo core (Paper XIII). The center of the Virgo Cluster is taken to be the location of the cD galaxy, M87, and the top axis shows  $R_p/r_{200}$ , where  $r_{200} = 1.55$  Mpc. There is a notable trend in specific frequency with clustercentric radius. All but one galaxy with  $S_{N,z} > 2$  are within  $R_p \sim 1$  Mpc of the cD.

538, which are known from their SBF distances to lie 6–7 Mpc behind the cluster core (Paper XIII). There is a remarkably clear correlation in which dwarfs closer to M87 have higher  $S_{N,z}$ . Of the 14 galaxies with  $S_{N,z} > 2$ , 13 lie within a projected radius of 1 Mpc (the exception being VCC 21, a possible dE/dI transition galaxy). We note that this is roughly half of  $r_{200}$  for the Virgo A subcluster centered on M87 (McLaughlin 1999b; Côté et al. 2001), where  $r_{200} = 1.55$  Mpc is the radius at which the mass density of M87 is 200 times the critical density. This value has been recalculated for  $D_{M87} = 16.5$  Mpc using the McLaughlin (1999a) model (D. McLaughlin 2007, private communication). When we plot  $S_{N,z}$  against three-dimensional (3D) clustercentric distance (Fig. 11), we see the same trend although the error in the line-of-sight distance introduces more scatter. Tables 6 and 7 list the values for the mean trends in these figures.

The M87 globular cluster system is one of the most extreme in the local supercluster, and it is possible that some of the M87 halo GCs, or a population of intracluster GCs, are contaminating the GC systems of the nearby dwarfs. In fact, the five galaxies closest to the cD, VCC 1297, VCC 1327, VCC 1279, VCC 1250, and VCC 1185, have been treated differently in our analyses in this paper because it is obvious that some to all of the GCs detected are part of the M87 halo. One of the telltale signatures of this is whether the GCs are uniformly distributed across the ACS WFC image rather than centrally concentrated around the targeted galaxy. Figure 12 shows the locations of GC candidates around the four dEs with the highest  $S_{N,z}$ . In all cases, the star clusters are concentrated toward the center of the galaxy and not in a uniform spatial distribution. An extension of the M87 GC radial density profile using a Sérsic model fit to the ACSVCS data and the ground-based data of McLaughlin (1999a) (see Fig. 1) predicts only one to three M87 GCs over the entire ACS WFC field of view at the distances of

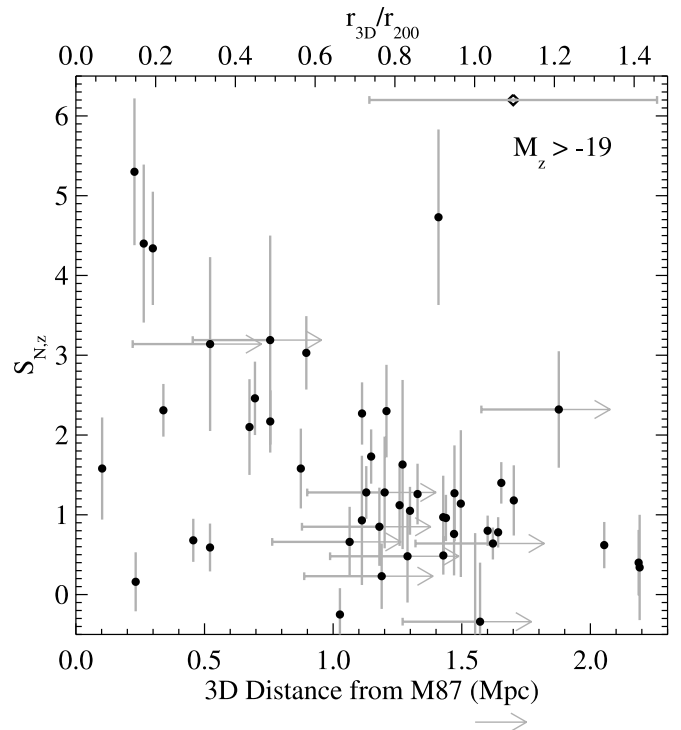


FIG. 11.— $S_N$  vs. 3D clustercentric distance for low-luminosity galaxies ( $M_z > -19$ ). The top axis shows  $R_{3D}/r_{200}$ , where  $r_{200} = 1.55$  Mpc. The error bar at top right shows the mean distance error. Points with arrows do not have SBF distances so we have used their projected distances and added 0.3 Mpc [the median ( $R_{3D} - R_p$ ) for the rest of the sample]. The left error bars are at their projected radii and are thus lower limits on their 3D radii.

VCC 1407, VCC 1545, and VCC 1539. The GC radial profile from Tamura et al. (2006) gives slightly higher values, with an expected  $9_{-6}^{+3}$  M87 GCs in the ACS field at the projected distance of VCC 1407, a dE that has 50 GCs. Given the difficulties in measuring the M87 GC density from ground-based data, especially given uncertainties in background subtraction, these independent estimates are consistent. Extrapolations of a de Vaucouleurs profile fit to the Tamura et al. (2006) data predict roughly two M87 GCs in the ACS fields for VCC 1545 and VCC 1539. We also have a parallel WFPC2 observation at a distance of  $40'$  from M87 in which we find a number of GCs consistent with zero. Finally, as these galaxies have 31–54 GCs each, this contamination from M87 GCs is expected to be at most 6%–18%.

The general trend of dEs having higher  $S_{N,z}$  is reversed for the two dwarfs (VCC 1297 and VCC 1185) with  $R_p \lesssim 100$  kpc. These galaxies have low intrinsic numbers of GCs, and we hypothesize that they have had their GCs stripped from them by M87. Both VCC 1297 and the more luminous VCC 1327 are within  $R_p \approx 40$  kpc and have GC numbers consistent with zero. The galaxies VCC 1192 and VCC 1199, which are at comparable distances from M49, also have undetectable intrinsic GC systems.

Figure 13 shows the positions of the ACSVCS galaxies on the sky, color- and size-coded by  $S_{N,z}$ , with open diamonds representing giants and filled diamonds for dwarfs. The high- $S_{N,z}$  dwarfs are clustered around M87, and there is no similar effect around the brightest galaxy in the cluster, M49, nor around VCC 1978 (M60/NGC 4649). Interesting exceptions are the aforementioned galaxies within the immediate vicinities of the giants ( $R \lesssim 40$  kpc), which have few or no GCs and may have been tidally stripped of their GC systems. The specific frequencies of the more massive galaxies are generally uniform, as shown in Figure 4, and thus are relatively unaffected by their distance from M87.

TABLE 6  
BINS OF  $R_p$  FOR GALAXIES WITH  $M_z > -19$

$R_p$ Range (Mpc)	$\langle R_p \rangle$ (Mpc)	$S_{N,z}$	$T$	$S_{N,z,\text{blue}}$	$S_{N,z,\text{red}}$
(0.00, 0.15).....	0.07	0.47	3.9	0.40	0.07
(0.15, 0.30).....	0.23	3.46	33.3	3.17	0.29
(0.30, 0.50).....	0.40	1.58	15.7	1.34	0.23
(0.50, 1.00).....	0.83	1.38	13.0	1.14	0.23
(1.00, 1.50).....	1.24	0.70	8.4	0.59	0.11
(1.50, 2.00).....	1.62	1.05	11.6	0.99	0.06

We can replot Figure 4, except divided into two samples, one inside and one outside a projected clustercentric distance of 1 Mpc. Figure 14 shows the difference between the two samples. Galaxies within 1 Mpc mirror the trends seen in Figure 4, with both giant and dwarf galaxies having high  $S_{N,z}$ . However, the sample of galaxies outside of 1 Mpc entirely lacks high- $S_{N,z}$  dwarfs. The division is not perfect, as there are still low- $S_{N,z}$  dwarfs within 1 Mpc (some of which have larger 3D distances), but the lack of high- $S_{N,z}$  dwarfs on the outskirts of the cluster or around the other massive ellipticals is very clear. The one possible exception to the general trend is VCC 230, which lies at  $R_p = 0.96$  Mpc and has an  $S_{N,z}$  higher by a factor of 2 than other galaxies at comparable distances.

The GCs in dwarf galaxies are predominantly from the blue (metal-poor) subpopulation, but many in the ACSVCS sample also possess small numbers of red GCs. Are both of these subpopulations affected by environment? In Figure 15 we plot  $S_{N,z}$  in dwarf galaxies against  $R_p$  for the blue and red GCs separately. As expected, the blue GCs mirror the overall trends since they dominate the GC budget. For the red GCs, although the errors are large, there is a slight hint that dwarfs with smaller clustercentric distances tend to have elevated red GC specific frequencies. The mean  $S_{N,z,\text{red}}$  of dwarfs within 1 Mpc is 0.2, compared to 0.1 for those outside.

Do the dwarfs with higher  $S_{N,z}$  exhibit any intrinsic characteristics that differentiate them from the others? Figure 16 shows the  $(g-z)-M_z$  color-magnitude diagram for the ACSVCS galaxies, singling out the dwarfs with both low and high  $S_{N,z}$ , and dividing the sample at  $S_{N,z} = 2$ . The top histogram plots the  $M_z$  distribution of the total, high- $S_{N,z}$ , and low- $S_{N,z}$  samples and shows how the luminosity distributions of the two groups are nearly identical. On the right histogram we plot the color distribution with the median colors marked by the dashed lines. Although the high- $S_{N,z}$  dwarfs formally have redder colors,  $\Delta(g-z) = 0.05$  mag, the difference is not significant. We have also investigated possible differences in their structural properties (Sérsic  $n$ ,  $r_e$ ), as well as in the colors of their globular clusters, and for none of these properties are high- $S_{N,z}$  dwarfs significantly different from low- $S_{N,z}$  dwarfs.

## 5. DISCUSSION

### 5.1. Mass Dependence of the GC Mass Fraction

Globular cluster specific frequency and its related quantities are clearly dependent on the mass of the host galaxy. However, GC fraction does *not* vary monotonically with galaxy mass, making it unlike most other properties measured for the ACSVCS galaxies, e.g., colors, structural parameters (Paper VI), core deficit and excess (Côté et al. 2007), GC mean colors (Paper IX), GC mean sizes (Paper X), and GCLF parameters (Paper XII). These papers (particularly Papers VI and VIII) show that there is no clear dis-

TABLE 7  
BINS OF  $R_{3D}$  FOR GALAXIES WITH  $M_z > -19$

$R_{3D}$ Range (Mpc)	$\langle R_{3D} \rangle$ (Mpc)	$S_{N,z}$	$T$	$S_{N,z,\text{blue}}$	$S_{N,z,\text{red}}$
(0.00, 0.25).....	0.19	1.38	11.3	1.19	0.18
(0.25, 0.50).....	0.34	2.39	23.9	2.19	0.20
(0.50, 1.00).....	0.71	2.13	20.8	1.87	0.26
(1.00, 1.50).....	1.27	1.27	13.5	1.05	0.22
(1.50, 2.00).....	1.65	0.78	7.7	0.68	0.10
(2.00, 2.50).....	2.14	0.50	5.2	0.47	0.03

inction between “dwarfs” and “giants” and that early-type galaxy properties have a smooth, monotonic dependence on galaxy mass. For GC fractions, however, a transition does appear to exist.

It is possible that GC fraction is telling us more about the variable formation efficiency of stars in the field than it is about the formation of GCs themselves. Blakeslee et al. (1997), Blakeslee (1999), and McLaughlin (1999a) have shown that, at least on the high-mass end, the number of GCs appears to scale directly to the baryonic, or even the total, mass. Kravtsov & Gnedin (2005) show a similar relation in their simulations of GC formation. High  $S_N$  values may then be interpreted as a lower fraction of baryons that form field stars.

The mismatch between total mass and stellar mass across the galaxy mass function is now well known in the study of galaxy formation. Observationally, dynamical mass estimates of galaxies across a wide range of mass require high mass-to-light ratios for both high- (Côté et al. 2001, 2003) and low-mass galaxies (Mateo 1998), with low  $\mathcal{M}/L$  measured for galaxies around  $L^*$  (Romanowsky et al. 2003; Peng et al. 2004; Napolitano et al. 2005). Mass-to-light ratios for ensembles of galaxies derived from weak lensing have also produced similar trends in  $\mathcal{M}/L$  (e.g., Guzik & Seljak 2002; Hoekstra et al. 2005; Mandelbaum et al. 2006). One can also take a statistical approach and match the expected dark matter halo mass distribution from simulations and the observed galaxy luminosity function (Berlind & Weinberg 2002; van den Bosch et al. 2003, 2007; Vale & Ostriker 2008), which does not require us to know the details of galaxy formation. These halo occupation studies infer a maximum conversion efficiency of baryons to stars at halo mass  $\mathcal{M}_h \sim 2 \times 10^{11} \mathcal{M}_\odot$ , or a stellar mass of  $\mathcal{M}_{G^*} \sim 7.5 \times 10^9 \mathcal{M}_\odot$ . In simulations of galaxy formation, this  $\mathcal{M}_h/L-L$  relation (or alternatively,  $\mathcal{M}_h/\mathcal{M}_{G^*}-\mathcal{M}_{G^*}$ ) requires various heating mechanisms (photoionization, stellar and AGN feedback) to prevent the formation of more stars in galaxies than are observed (Benson et al. 2002; Croton et al. 2006).

What happens if we make the simple assumption that  $N_{GC} \propto \mathcal{M}_h$ , or equivalently  $S_M \propto \mathcal{M}_h/\mathcal{M}_{G^*}$ ? Ideally, we would know the total dynamical mass of each galaxy in our sample. Without this information, we can apply average halo mass-to-light ratios derived from the halo occupation studies, in this case using the parameterization of van den Bosch et al. (2007) and a WMAP3 cosmology (Spergel et al. 2007). We transform their  $L_B$  to  $\mathcal{M}_{G^*}$  using a range of  $\mathcal{M}_{G^*}/L_B$ , monotonically increasing from 1.8 to 4.1 as a function of galaxy luminosity, derived from a polynomial fit to the inferred  $\mathcal{M}_{G^*}/L_B$  and measured  $L_z$  of the ACSVCS galaxies. This results in a simple prediction for the behavior of  $S_M$  as a function of galaxy stellar mass.

In Figure 17 we plot the ACSVCS data for  $S_M$  against  $\mathcal{M}_{G^*}$  with errors for both  $S_M$  and  $\mathcal{M}_{G^*}$ . We extend the data to lower mass galaxies by including the Virgo and Fornax dEs and five Local Group dEs from Lotz et al. (2004). For the latter data, we

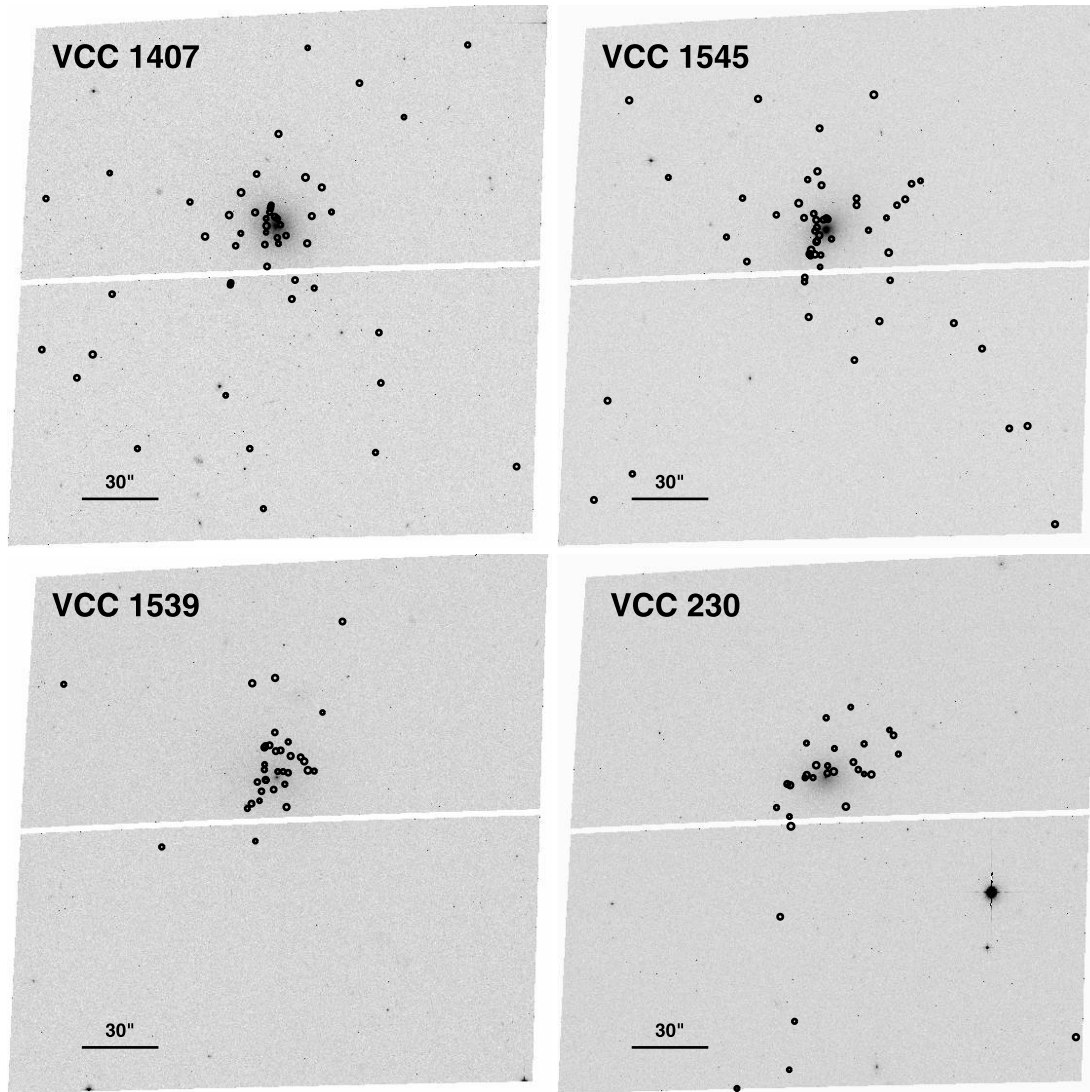


FIG. 12.—Spatial distributions of GC candidates (*circles*) in four dEs with the highest specific frequencies (VCC 1407, VCC 1545, VCC 1539, and VCC 230). In all cases, there is evidence of a rich GC system that is centrally clustered around the galaxy, showing that their elevated  $S_N$  is not due to an enhanced level of interloping GCs from the cD galaxy or an intracluster population. These galaxies have  $\sim 5$  times more GCs than “normal” galaxies with  $S_{N,z} = 1$ . The images show the entire ACS field of view, and the scale bar in the lower left of each image has a length of  $30''$ . The blank diagonal strip is the gap between the two ACS CCDs. Catalogs have been statistically cleaned using expected contamination determined from control fields.

assume  $\mathcal{M}_{G^*}/L_B \sim 2.5$  (an extension of the mass-to-light ratios fitted to the ACSVCS dwarfs) and  $\langle \mathcal{M}_{GC} \rangle = 1.6 \times 10^5 \mathcal{M}_\odot$ , the mean GC mass in our ACSVCS dwarfs. The Virgo and Fornax dE points do not include errors in their stellar masses, and for the Local Group dEs, we have assumed an error in their luminosities of 0.25 mag. Overplotted on the data is the predicted behavior of  $S_M$  assuming  $S_M \propto \mathcal{M}_h/\mathcal{M}_{G^*}$ . The only truly free parameter in the plot is the vertical normalization, which is chosen to match both the massive elliptical M49 and the Local Group dwarf spheroidal galaxies Fornax and Sagittarius. The assumption that the GC mass fraction is proportional to the inverse of the stellar mass fraction,  $S_M \propto \mathcal{M}_h/\mathcal{M}_{G^*}$ , predicts that  $S_M$  should be high at both high and low masses. At high stellar masses, the rise reflects the increasing  $\mathcal{M}/L$  of massive elliptical galaxies. One galaxy, M87, is well off the predicted relation, but this is because its GC population is more likely representative of the entire Virgo A subcluster within which M87 resides (McLaughlin 1999a). The sharp cutoff in the galaxy luminosity function and the nearly universal luminosity of BCGs mean that  $S_M$  can vary widely depending on whether the luminosity is associated with a galaxy

halo or the cluster halo (Blakeslee 1999; Jordán et al. 2004a). At low masses,  $S_M$  is also predicted to rise steeply, but this is at masses *lower* than those in our sample. Many of the ACSVCS dwarfs around  $\mathcal{M}_{G^*} \sim 10^9 \mathcal{M}_\odot$  have  $S_M$  higher than we would infer by a factor of 4–8, and these are the same ones that are in the inner regions of the Virgo Cluster. At masses below  $\mathcal{M}_{G^*} \sim 10^8 \mathcal{M}_\odot$ , the dEs from Lotz et al. (2004) are either consistent with or below the expected rise, although the observational error is quite large (and not shown for clarity). So, while the linear scaling of  $\mathcal{M}_{GC} \propto \mathcal{M}_h$  does seem to be relevant to  $S_M$  across a large galaxy mass range, the  $S_M$  for the ACSVCS dwarf galaxies in this diagram is not explained by a dependence on galaxy mass alone. The scatter in  $S_M$  for dwarfs (and also possibly for intermediate-mass galaxies) also cannot be explained by the intrinsic scatter in galaxy  $\mathcal{M}/L$  (from van den Bosch et al. 2007). The scatter is even harder to explain for  $S_L$ , where the errors are considerably smaller.

Both Forbes (2005) and Bekki et al. (2006) have also attempted to explain the increase in  $S_N$  for dwarf galaxies by assuming a relation between GCs and galaxy mass. These attempts followed similar arguments by Durrell et al. (1996a) and McLaughlin

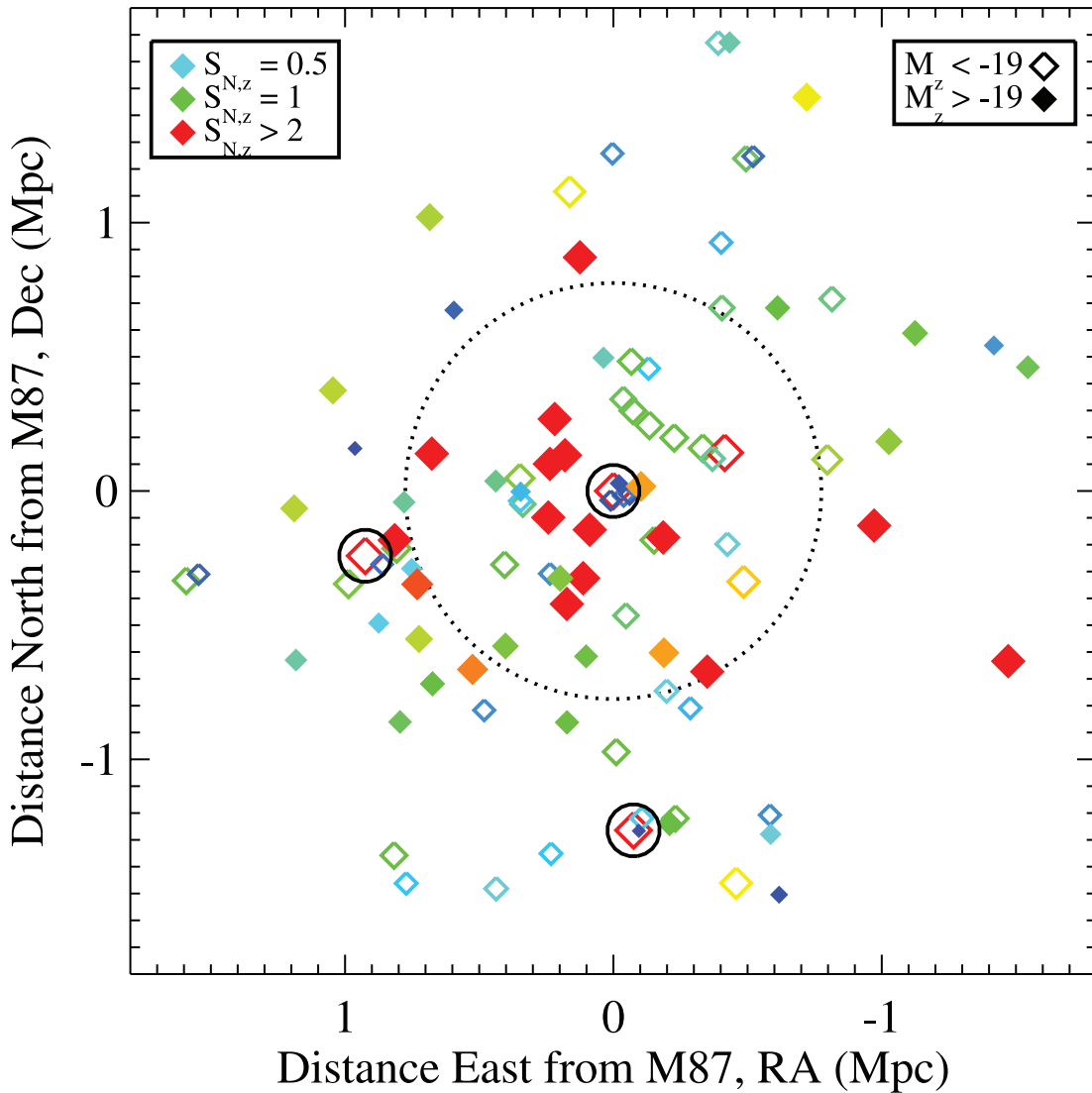


FIG. 13.—Spatial distribution of ACSVCS galaxies color-coded by  $S_{N,z}$ . Specific frequency increases as colors change from blue to red. Filled diamonds are dwarfs ( $M_z > -19$ ), and open diamonds are the more luminous galaxies. The three large solid circles represent the three massive ellipticals in the Virgo Cluster: the cD galaxy M87 (center), M49 (bottom), and M60 (VCC 1978/NGC 4649; left). The large dotted circle represents  $r_{200}/2 = 775$  kpc for the M87/Virgo A subcluster. The high- $S_{N,z}$  galaxies are preferentially around the cD, while no enhancement appears around the cluster's most luminous galaxy M49. Dwarfs immediately in the vicinity of the giants have low  $S_{N,z}$  and may be tidally stripped.

(1999a). Forbes (2005) used the data for ACSVCS galaxies presented in Paper IX to compare the approximate  $S_N$  of blue GCs in dwarf galaxies with the scaling relation of Dekel & Birnboim (2006), where  $\mathcal{M}/L_V \propto \mathcal{M}^{-\alpha}$ . Forbes (2005) used  $\alpha = \frac{2}{3}$  from Dekel & Birnboim (2006) for masses below the critical galaxy stellar mass of  $3 \times 10^{10} \mathcal{M}_\odot$  where galactic winds should be more efficient at blowing gas and metals from the galaxy (Kauffmann et al. 2003; Tremonti et al. 2004). They also assume a constant number of GCs as a function of galaxy mass, giving  $S_N \propto \mathcal{M}^{-5/3}$ . Although this formulation predicts a rapid rise in  $S_N$  and  $S_M$  for fainter galaxies, it is much too steep to accommodate the data in the way that they present it. The normalization of this relation as presented in Forbes (2005) is not constrained, but because the relation is steep, the normalization is critical for comparing to the data. If we fix to the mean  $S_M$  at the critical mass, the relation overpredicts  $S_M$  for all the dwarfs. This is due to the fact that dwarf galaxies do not have a constant number of GCs. Bekki et al. (2006) simulate GC systems and test various values for  $\alpha$ , concluding that the observations are not consistent with  $\alpha = 0$  (a constant  $\mathcal{M}/L$ ), and it would be interesting to determine  $\alpha$  again

with this new data set. Whatever the best value of  $\alpha$ , however, a simple relation where the specific frequency is determined solely by the present-day stellar mass of a galaxy is unable to describe all of the data on the low-mass end.

## 5.2. Environment and Biased GC Formation

### 5.2.1. Cluster Dwarfs in the Millennium Simulation

As we showed in Figures 10–15, GC specific frequency is not purely a function of galaxy mass, but can also depend on environment. We do not see environmental effects on  $S_N$  for massive galaxies with the exception of M87, whose location defines the cluster center. For intermediate-mass galaxies in particular,  $S_N$  does not appear to depend on environment. For dwarf galaxies, however, specific frequency is a strong function of proximity to the center of the Virgo Cluster. Nearly all dwarfs with high  $S_N$  are within  $R_p = 1$  Mpc.

One of the main problems with building up the blue GC subpopulation in the halos of massive galaxies is that the ratio of metal-poor GCs to metal-poor field stars in galaxies is very high

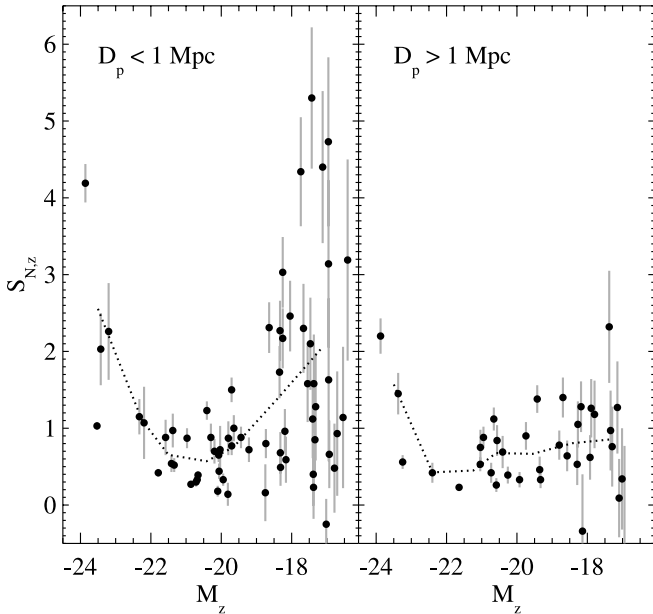


FIG. 14.— $S_{N,z}$  vs. galaxy  $M_z$ , as in Fig. 4, except divided by clustercentric distance. Galaxies within 1 Mpc of the cluster center exhibit the full range of  $S_{N,z}$ , but nearly all the dwarfs outside of 1 Mpc have low specific frequencies ( $S_{N,z} \leq 1.5$ ). The only exception is a dE/dI transition object (VCC 21).

(Harris & Harris 2002). In other words, if halos are built up through the accretion of dwarf galaxies or dwarf-mass fragments, then their specific frequencies need to be high in order to keep the main body of the galaxy metal-rich. These central dwarfs with high  $S_N$  may then hold the key to the formation of GC systems, as they are likely to be most similar to the progenitors of any dwarfs that have since merged into the halo of M87. Figure 9 shows that the  $S_{N,z, \text{blue}}$  values for the innermost dwarfs are as high or higher than that of M87 or any of the other giant ellipticals. These dwarfs could be the survivors of an accreted population of protogalaxies that may have had even higher GC mass fractions and redder GCs. In fact, the two galaxies closest to M87, VCC 1327 and VCC 1297, have GC systems that are either entirely stripped or undetectable against the M87 GC system, and their  $g - z$  colors are much redder than those of other galaxies with the same luminosity (see the red outliers in Fig. 16).

What is the cause of the enhanced GC fractions in these central dwarfs? Are they more efficient at producing GCs for their mass, less efficient at forming field stars, or better at keeping the GCs they form? As discussed earlier, one possible way for GC production to be biased toward the central dense regions of the galaxy cluster is for these galaxies to form a larger fraction of their stars earlier and at a higher star formation rate density. Although we do not currently have the ability to determine the detailed star formation histories of the ACSVCS galaxies at these early times, we can use simulations to study global trends in the star formation histories of low-mass early-type galaxies in the cluster environment. By using simulations, we can test the consistency of the hypothesis that the central dwarfs must both be older and have higher SFR densities.

The simulation we use for comparison is the ‘‘Millennium Simulation’’ carried out by the Virgo Consortium (Springel et al. 2005) coupled with the semianalytic models presented by De Lucia et al. (2006). This simulation consisted of  $2160^3$  dark matter particles followed from  $z = 127$  to the present day in a volume  $500 h^{-1}$  Mpc on a side. The spatial resolution is  $5 h^{-1}$  Mpc, and the simulation is essentially complete for all galaxies with stellar

mass greater than  $3 \times 10^8 M_\odot$ . This mass limit is comparable to the lowest mass galaxies in the ACSVCS. From this simulation, we selected a sample of low-mass early-type cluster galaxies and their progenitors. We identified 126 massive galaxy clusters using the same criteria as in De Lucia & Blaizot (2007), requiring a halo mass greater than  $7 \times 10^{14} M_\odot$  at  $z = 0$ . This mass limit also corresponds roughly to the mass of Virgo (Bohringer et al. 1994). Within each of these clusters, we then selected all galaxies with  $M_{z, \text{SDSS}} > -19$  that have early-type morphologies, matching the low-luminosity sample we have in the ACSVCS. Early-type galaxies were defined by the bulge-to-total luminosity ratio with  $\Delta M < 1.56$  ( $\Delta M = M_{\text{bulge}} - M_{\text{total}}$ ), using the empirical criteria of Simien & de Vaucouleurs (1986). In addition, we selected only galaxies with  $g_{\text{SDSS}} - z_{\text{SDSS}} > 0.5$  at  $z = 0$  in order to best match the colors of the ACSVCS galaxies. In total, we selected 15,506 simulated galaxies at  $z = 0$  and all of their progenitors over 63 snapshots back to  $z \sim 12$ .

In Figure 18 we show the mass-weighted age of the stellar populations in the selected early-type dwarfs as a function of distance from the center of the cluster’s dark matter halo. The error bars show the  $1 \sigma$  width of the age distribution in each bin of radius. This figure clearly shows that the stellar age of the simulated dwarfs has a strong dependence on clustercentric radius, at least within a distance of 1.5 Mpc, with the central dwarfs having mean ages of 11.5 Gyr decreasing to 9 Gyr at the cluster outskirts. The radius within which the age gradient is evident is also the radius in Virgo within which we see elevated GC fractions.

The central dwarfs in the simulation may be older, but they also need to form their stars more intensely (higher peak star formation rates and densities) in order to have more of their stars in massive clusters. The top panel of Figure 19 shows the average normalized star formation rate of early-type cluster dwarfs as a function of look-back time. The star formation rate is normalized to the final stellar mass of the galaxy. The two curves represent SFR histories for dwarfs divided into two bins of clustercentric radius at 1 Mpc from the cluster center. We note that these are average star formation histories and that for any individual galaxy, the bursts of star formation are more intense, brief, and stochastic. Combining galaxies allows us to see that the central dwarfs have a more peaked SFR at earlier times with a rapid falloff, whereas the outer dwarfs have not only a lower peak but also more star formation extending to later times.

It has been suggested that massive star clusters form preferentially in high-pressure environments (Harris & Pudritz 1994; Elmegreen & Efremov 1997; Ashman & Zepf 2001), of which SFR surface density is one possible indication. The bottom panel of Figure 19 shows the ratio of SFR surface density in central dwarfs to that in outer dwarfs. We calculate the SFR surface density using the SFR and the disk radius, which the semianalytic models calculate using the analytic model of Mo et al. (1998). Not only is the peak star formation rate higher in central dwarfs, but the *intensity* of star formation, as measured by the SFR surface density, is also higher during the epoch when these galaxies are forming most of their stars. At later times, the SFR surface density in the outer dwarfs is higher on average than that in central dwarfs, but this is at much lower absolute star formation rates and SFR surface densities. If GCs are formed in the same events that produced the bulk of these low-mass galaxies, then we would expect that the central dwarfs would retain a higher fraction of their stellar mass in massive star clusters.

### 5.2.2. Inferring Cluster Formation Histories

We can quantify these effects more directly by inferring the cluster formation rate (CFR) history, using an empirical relationship

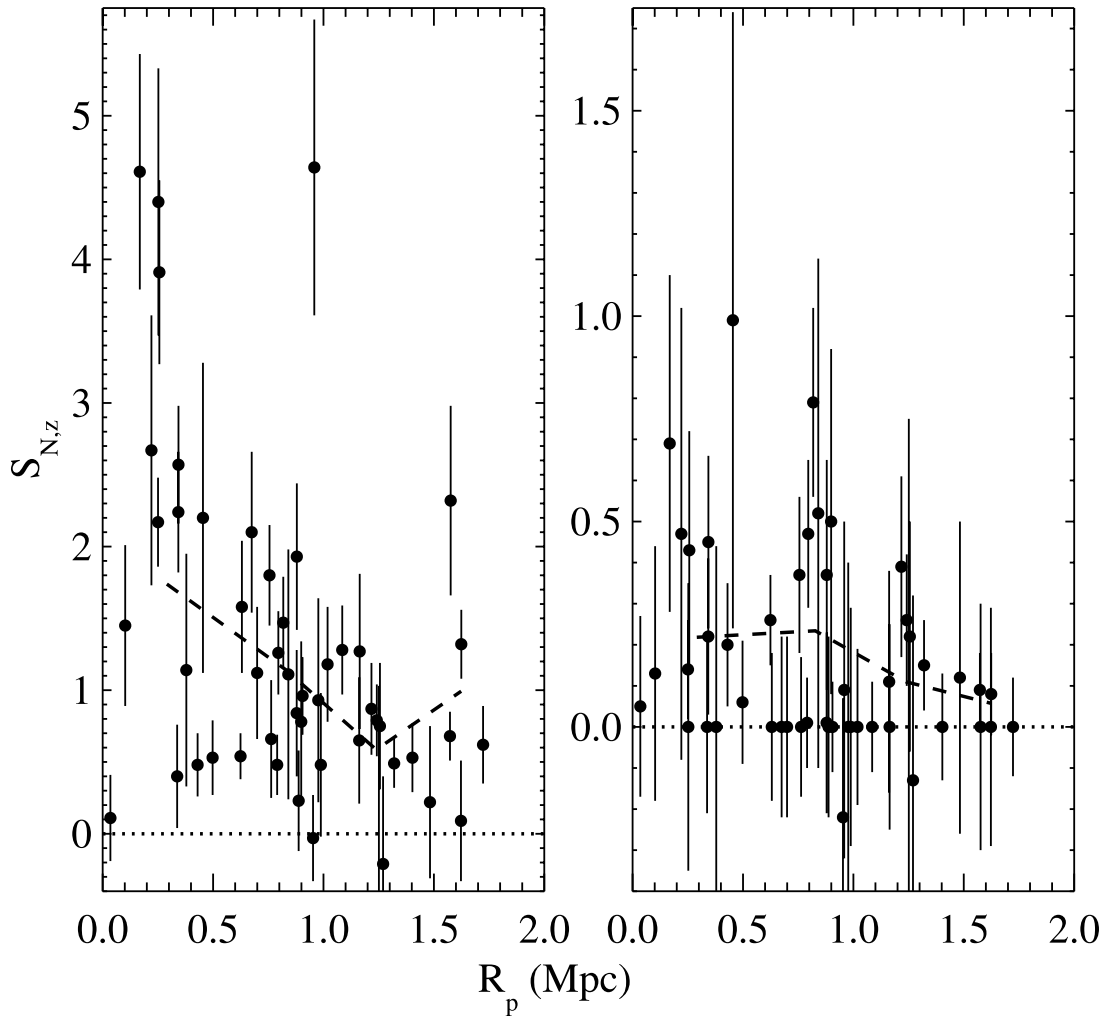


FIG. 15.— $S_{N,z}$  vs. projected clustercentric distance for blue GCs (*left*) and red GCs (*right*) in low-luminosity galaxies ( $M_z > -19$ ). Given that most GCs in dwarfs are metal-poor, it is not surprising that the trend for blue GCs mirrors that in Fig. 10. The  $S_{N,z}$  for the red GCs may also show a tendency to be enhanced at small clustercentric radii, but the numbers involved are small. [See the electronic edition of the *Journal* for a color version of this figure.]

between  $S_M$ , the fraction of stellar mass formed in massive star clusters, and the SFR surface density,  $\Sigma_{\text{SFR}}$ . Fitting the data of Larsen & Richtler (2000), we adopt  $S_M \propto (\Sigma_{\text{SFR}})^{0.8}$ , up to a maximum efficiency of 100%. The cluster formation rate then scales with a combination of the SFR and the SFR surface density as  $\text{CFR} \propto \text{SFR}(\Sigma_{\text{SFR}})^{0.8}$ . The Larsen & Richtler (2000) data measured the fraction of luminosity in young massive clusters (YMCs), which had ages between 10 and 500 Myr and more massive than  $\sim 3 \times 10^4 M_\odot$ . Given that most clusters will be destroyed over a Hubble time, we need to estimate what fraction of these clusters will survive. We adopt the two-stage disruption law of Whitmore et al. (2007) where “infant mortality” causes clusters to disrupt in constant numbers,  $dN/d\tau \propto \tau^{-1}$ , for ages less than a few times  $10^8$  Gyr (Fall et al. 2005) and a constant mass loss over a Hubble time of  $\mu_{\text{ev}} = 1.9 \times 10^4 M_\odot \text{ Gyr}^{-1}$  due to two-body relaxation (Fall & Zhang 2001; Paper XII). Starting with the data of Larsen & Richtler (2000), we estimate that  $\sim 1.8\%$  of their young massive clusters survive to the present day. Using this formalism, we produce the corresponding cluster formation rate for central early-type dwarfs, normalized to the final mass in GCs.

Figure 20 compares the average SFR and SFR surface density as a function of redshift for the sample of central dwarfs (*top*) and outer dwarfs (*bottom*). While the SFR peaks at  $z = 3.5\text{--}4$ ,  $\Sigma_{\text{SFR}}$

is highest at  $z = 10$  and falls a factor of 35 by  $z = 4$ . Because the formation rate of massive star clusters depends on both quantities, their peak epoch of formation will be *earlier* than that of the stars. The result is that the CFR peaks at earlier redshift,  $z \sim 4.5\text{--}5$ , than the total star formation rate, a difference of  $\sim 350\text{--}500$  Myr. In all dwarfs, the dependence of massive star cluster formation on the SFR surface density naturally produces GCs that are older and more metal-poor than the stars.

Depending on how massive star cluster formation scales with SFR and SFR density, we also expect that the mean age difference between the GCs and the field would be different in the central and outer dwarfs. Because GCs in these dwarfs will mostly form at or before the SFR peak, they will always be old, but the field stars can continue to form at later times. This is evident when comparing the difference in look-back time for the SFR peaks for the central and outer simulated dwarfs in the top panel of Figure 19 and the difference in their mean ages in Figure 18. The outer dwarfs have mass-weighted mean ages that are younger by  $\sim 2.5$  Gyr, but the times of their peak star formation rates differ by only 0.1 Gyr. Figure 20 shows how the CFR peak happens at similar redshifts for all dwarfs, but the peak in star formation happens later in outer dwarfs. One prediction is then that, on average, the difference between the mean age of the metal-poor GCs and the mean age of the stars should be larger in dwarf

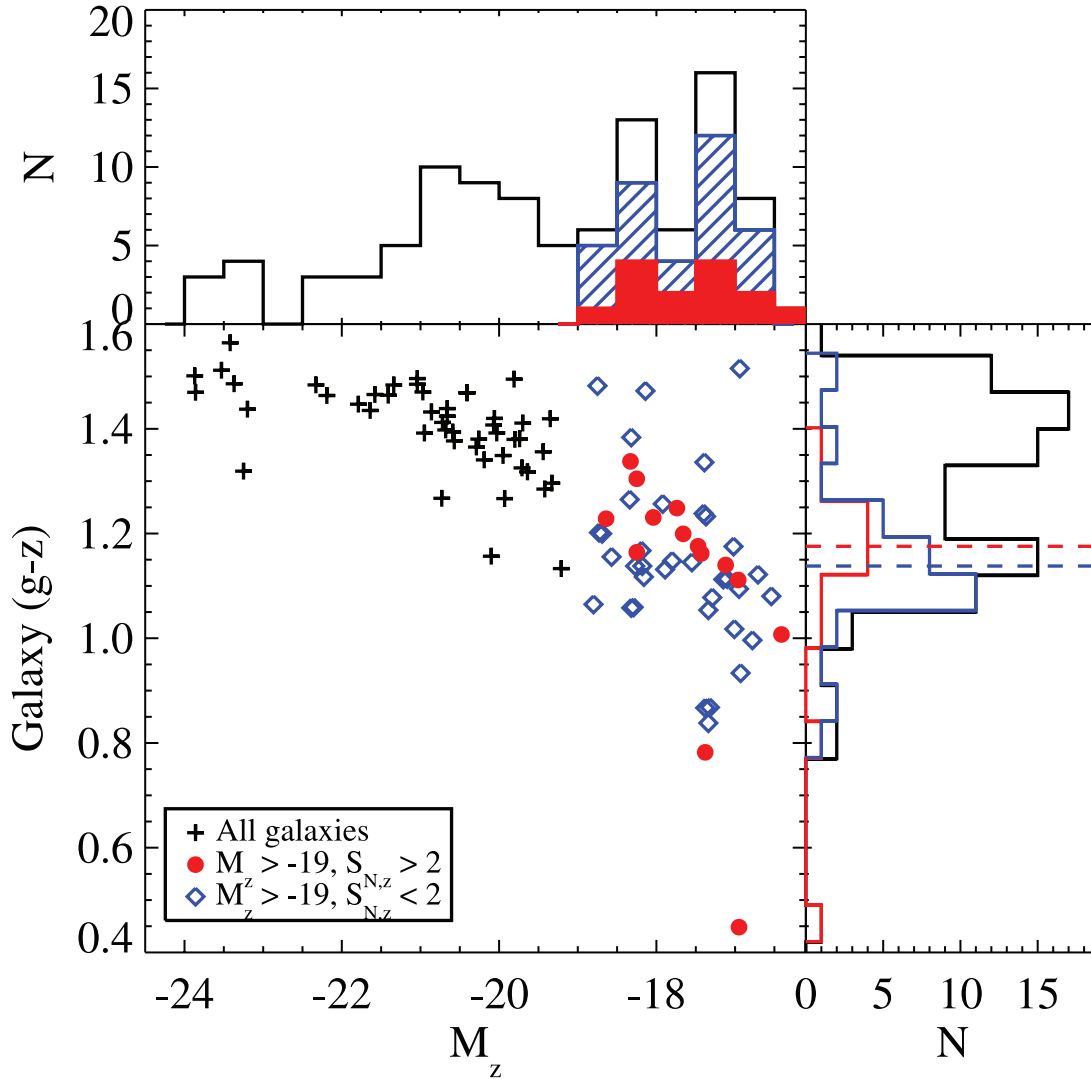


FIG. 16.—Color-magnitude properties of dwarfs with high and low  $S_{N,z}$ . We show  $(g-z)$  vs.  $M_z$  and accompanying histograms for the ACSVCS sample of galaxies. Galaxies with  $M_z > -19$  are separated at  $S_{N,z} = 2$ . The luminosity distributions of the two subsets are nearly identical, while the high- $S_{N,z}$  dwarfs have a median color that is formally redder than that of the low- $S_{N,z}$  dwarfs (shown by dashed lines in the right histogram).

galaxies with lower GC mass fractions (i.e., the mean age of the galaxy is proportional to  $S_M$ ).

Lastly, we test whether this framework can produce higher GC mass fraction in the central dwarfs. Figure 21 shows  $S_M$  against clustercentric radius for simulated dwarfs. The absolute scale for  $S_M$  is highly sensitive to the assumed destruction but is not important for our purposes. It is the relative comparison between central and outer simulated dwarfs that is of interest, and the high- $S_M$  simulated dwarfs do appear preferentially in the central regions, similar to what the data show in Figures 10 and 11.

The simulations do not match the data exactly, as there also appear to be many low- $S_M$  dEs within 1 Mpc, and the metallicities of the GCs as derived from the simulation are generally too high, but it is encouraging that this simple scaling for cluster formation in the Millennium Simulation can reproduce many of the observed trends. A more detailed treatment of ram pressure and tidal stripping in cluster cores in a higher resolution simulation may help resolve some discrepancies.

### 5.3. Possible Mechanisms for Quenching Star Formation

The central dwarfs may produce a higher fraction of GCs at or before their peak SFR, but it is also important that their sub-

sequent star formation is rapidly quenched or kept at a low level. At these and lower masses, photoionization heating of the interstellar medium by the UV background is believed to be an important mechanism for suppressing star formation in dwarf galaxies (Bullock et al. 2000; Benson et al. 2002), and this could universally truncate the epoch of efficient formation of metal-poor GCs, although detailed work still needs to be done to show that reionization is a plausible mechanism for halting GC formation. At least one scenario (Cen 2001) has reionization as the cause for GC formation, another suggests that starburst-driven shocks may trigger GC formation (Scannapieco et al. 2004), and yet another shows that GCs themselves could plausibly be the source of reionizing photons (Ricotti 2004). So, the relationship between GC formation and reionization, if any, is still very much undetermined.

One potential problem with the scenario where reionization halts GC formation is that in the simulations of Bekki et al. (2006), the bias introduced by the truncation redshift is not very strong. Changing the truncation redshift from  $z = 15$  to 6, for example, does not alter the  $S_N$  of galaxies by very much, and at least in these simulations, reionization-induced biased GC formation by itself does not seem to be able to explain the scatter in  $S_N$

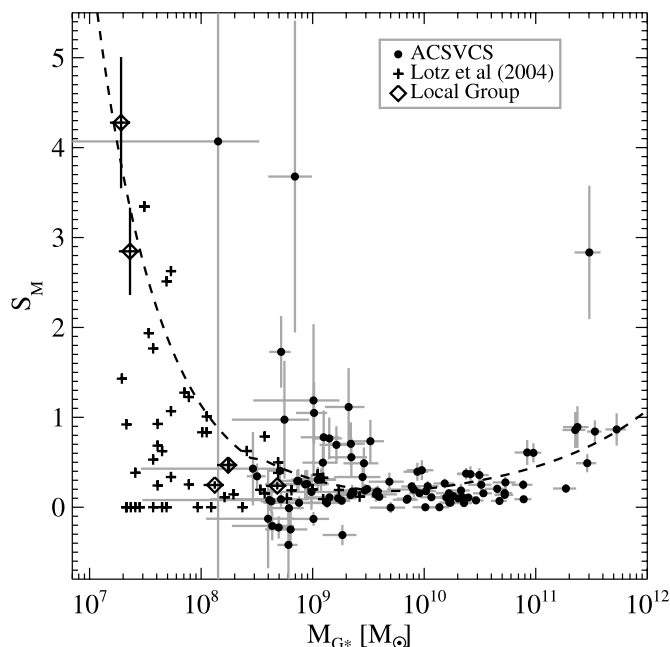


FIG. 17.—Mass fraction of GCs,  $S_M$ , vs. galaxy stellar mass  $M_{G*}$  for ACSVCS galaxies (*filled circles*), Virgo and Fornax dEs from Lotz et al. (2004) (*plus signs*), and Local Group early-type dwarfs (*diamonds*). We overplot the expected behavior of  $S_M$  assuming that the mass in GCs follows the total halo mass as inferred from the  $M_h/L-M_h$  relation of van den Bosch (2007). This assumption can explain the rise in  $S_M$  for luminous galaxies, although the cD galaxy M87 is off the relation because its stellar mass is not representative of its halo mass. At the low-mass end, we also expect a sharp rise in  $S_M$ , which is traced by some cluster and Local Group dEs. At masses of  $\sim 10^9 M_\odot$ , however, many dEs in the ACSVCS sample have higher  $S_M$  than might be expected. The high- $S_M$  and low- $S_M$  dwarfs cannot be explained simultaneously if GC mass fraction is solely a function of galaxy stellar mass.

for dEs. However, it may be sufficient to explain the general trend toward higher GC fractions at lower masses seen in Figure 17.

Another plausible explanation is that the environment plays an important role in quenching the later, lower level star formation that builds up the field. Ram pressure stripping of gas will be more efficient in low-mass galaxies at the center of the cluster. Since these halos are the earliest to fall into the cluster, they will have their star formation quenched earlier and more efficiently than their counterparts in the outer cluster regions. Moore et al. (1998) also showed through numerical simulations that “galaxy harassment” in clusters is efficient at transforming small spirals into low-mass spheroidal galaxies. More recently, Mayer et al. (2007) used hydrodynamic simulations to demonstrate that the environmental effects of tidal shocks and ram pressure stripping can combine to create the most dark matter-dominated dwarf spheroidals (dSphs) in the Local Group. Although the Virgo dEs are much more massive than dSphs such as Draco or Ursa Minor, it is possible that similar processes caused them to fail in their conversion of gas to field stars. Ram pressure stripping in Virgo has been observed out to a cluster radius of 0.8 Mpc (Kenney et al. 2004), similar to the radius within which we see enhanced GC fractions. The comparison of central dEs and those on the outskirts, however, shows no strong differences between their observed properties. If harassment was important, we might expect different morphologies or surface brightness profiles (Sérsic  $n$  or  $r_e$ ) in the two groups, but this is not the case. The cumulative GC color distributions of these two populations of dEs are also indistinguishable. If the central dEs were once much more luminous, we would expect them to be outliers on the color-magnitude re-

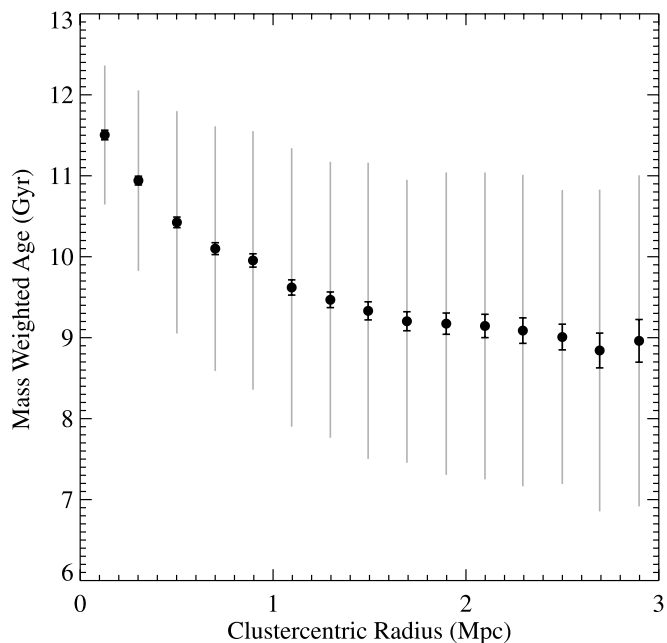


FIG. 18.—Mass-weighted stellar age of low-mass ( $M_z > -19$ ) early-type cluster galaxies in the Millennium Simulation. Galaxies in the inner regions of the galaxy cluster have older mean ages. The black error bars with hats depict the  $2\sigma$  error in the mean, while the gray error bars indicate the  $1\sigma$  width of the age distribution in each bin.

lation, but the ones with high  $S_N$  are not. Only the innermost galaxies that are well within the M87 GC system and have no GCs *are* outliers and are red for their luminosities. These galaxies are good candidates for having been harassed.

#### 5.4. Globular Cluster Destruction

Up until now, we have only discussed the possible variation in formation histories to explain the observed trends in GC fraction. However, star clusters can also be destroyed, and we expect that a substantial fraction of clusters initially formed will not survive a Hubble time. Observations of young star clusters in nearby galaxies show that most star clusters are disrupted very early (Fall et al. 2005; Bastian et al. 2005). In addition, subsequent evolution through two-body relaxation, stellar mass loss, and tidal shocks will destroy even more low-mass clusters (Fall & Rees 1977; Vesperini 1998; Fall & Zhang 2001). Globular clusters are thus the survivors of what was once a much larger population of star clusters. Any differences in the survival rates between different galaxies could produce different GC fractions in the present day.

Unfortunately, the cosmological simulations used to study galaxy formation cannot at the present time model GC destruction. One benefit of studying trends in terms of  $S_M$ , or GC mass fraction, is that while most of the destroyed star clusters after the initial 1 Gyr are preferentially low mass, most of the mass remains in the higher mass objects. Thus, GC mass fraction is more robust against disruption processes that preferentially destroy low-mass clusters. If, however, there was a mechanism that could affect the survival efficiency of massive clusters, perhaps variable infant mortality or dynamical friction as influenced by the dark matter halo profile, then destruction could play a role in driving the trends in GC mass fraction. This mechanism would require central dwarfs to be able to retain a larger fraction of the massive star clusters than their counterparts at larger cluster radii. This is perhaps counter to what one might expect if the cluster tidal field



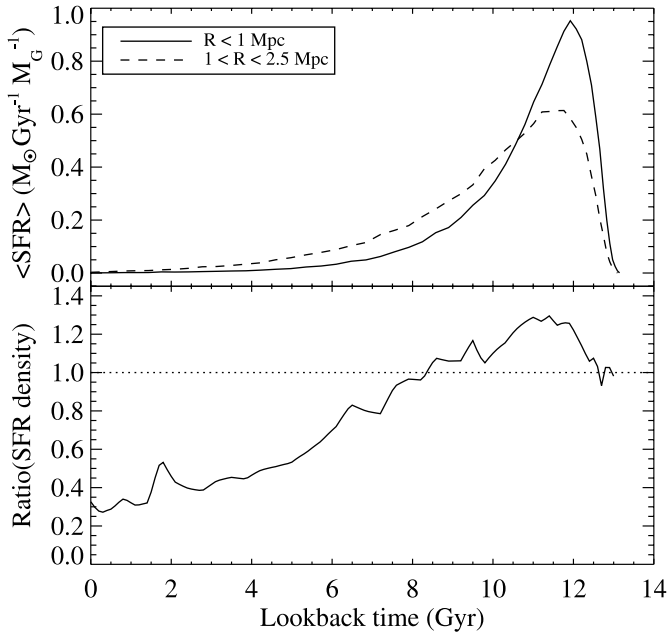


FIG. 19.—*Top*: Average normalized star formation rate vs. look-back time for low-mass early-type cluster galaxies in the Millennium Simulation. The solid line shows the average SFR history for “central” galaxies within 1 Mpc of the cluster center, and the dashed line shows “outer” galaxies. Low-mass early-type galaxies in the inner cluster regions formed their stars earlier and with higher peak SFR than those in the outskirts. Higher peak SFRs could result in higher globular cluster mass fractions. *Bottom*: Average ratio in SFR surface density between central and outer low-mass early-type cluster galaxies as a function of look-back time. Galaxies in the inner regions have a higher SFR surface density during the period of highest total star formation ( $t_{\text{look-back}} > 8$  Gyr). This results in central galaxies having a higher mass fraction in massive star clusters.

had a role in stripping or destroying GCs. We also see no signs in our data that central dwarfs have higher mean GC masses or different GCLFs, things that might point to internal destruction being a dominant mechanism driving GC fractions.

For a real test of the mechanisms that drive GC formation efficiency, it would be extremely beneficial to have more and better simulations of galaxies and their star cluster populations. Both hydrodynamic simulations and dark matter simulations coupled with semianalytic models can provide useful quantitative predictions. These would ideally be able to produce a  $z = 0$  volume comparable to the Virgo Cluster with a stellar mass resolution that would allow the resolution of GC mass objects ( $\sim 10^5 M_{\odot}$ ).

### 5.5. $S_N$ Variation in Metal-poor and Metal-rich GCs

Another result of this study is the behavior of the specific frequencies of blue and red GCs. The variation in  $S_N$  is dominated by the blue, metal-poor GCs, although the red GCs do exhibit similar trends with fewer GCs and lower significance. Figure 8 shows that the fraction of red GCs rarely exceeds 50% and that the trend of increasing red fraction with galaxy mass either flattens or turns over for the most massive galaxies. This suggests that massive ellipticals have accreted the most dwarflike galaxies, the presumed original hosts of metal-poor GCs, or that the dwarfs that they accreted had the highest  $S_N$ , or both. Also, the rise in red GC fraction from  $M_z = -19$  to  $-22$  requires that  $S_{N,z,\text{blue}}$  and  $S_{N,z,\text{red}}$  respectively fall and rise in equal proportions to produce the nearly constant  $S_{N,z}$  over this range. The fraction of blue GCs is lowest where the total  $S_{N,z}$  is at a minimum, implying again that they are driving the overall trends.

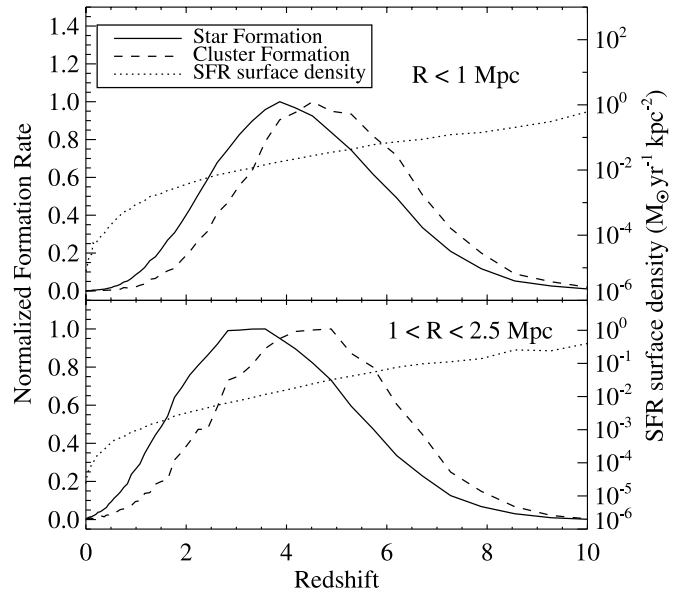


FIG. 20.—For central (*top*) and outer (*bottom*) early-type dwarfs in the Millennium Simulation, a comparison of the average star formation rate (*solid line*), the average star formation rate surface density (*dotted line*), and the inferred average cluster formation rate (*dashed line*). The highest SFR surface densities occur at earlier times than the highest SFRs because the disks within which stars form are smaller at high redshift. If massive star clusters preferentially form at high SFR surface densities, then GCs in early-type dwarfs will, on average, have older ages and lower metallicities than the field stars. Formation rates are normalized to their maximum values.

The higher values of  $S_{N,z,\text{red}}$  for the massive galaxies, in particular M87, mean that the formation of metal-rich GCs also occurs at different efficiencies. For massive galaxies, these results are consistent with the data of Rhode et al. (2005), who found that the mass-normalized number of GCs,  $T_{\text{red}}$ , is positively correlated with galaxy mass (shown in Fig. 4 of Brodie & Strader 2006). Could these metal-rich GCs have been accreted from dwarfs in a scenario similar to that described for the metal-poor GCs? Figure 15 shows that the number fraction of red GCs may also be enhanced in central dwarfs, like the metal-poor GCs, but it is unclear whether a biasing scenario similar to the one described above provides an explanation since we would expect that the formation of metal-rich GCs at a later time might actually be suppressed. However, the numbers we observe are very small and uncertain and may be sensitive to the dividing color between blue and red GCs, which is chosen to be the same in all galaxies at these luminosities. Another possibility is that the trend in the red GCs is determined more by the mass of the host, with red GCs forming in the most massive progenitor, and the more massive protogalaxies forming their stars earlier and with higher star formation rates. This is consistent with observations of the stellar populations of elliptical galaxies, which find that the stars in massive ellipticals ( $L > L^*$ ) have old ages and high  $[\alpha/\text{Fe}]$ , i.e., they formed early and rapidly, while lower mass early-type galaxies have younger mean ages and lower  $[\alpha/\text{Fe}]$  (Thomas et al. 2005). Similar trends in  $[\alpha/\text{Fe}]$  are seen in the GCs of these galaxies (Puzia et al. 2006).

The fact that we see the same trends for both the blue and red GC subpopulations suggests the possibility that the two populations are not very distinct, and that their separation is merely the result of placing a dividing line on a continuum of GC properties. This possibility has been raised by Yoon et al. (2006), who suggest that the GC metallicity bimodality in early-type galaxies

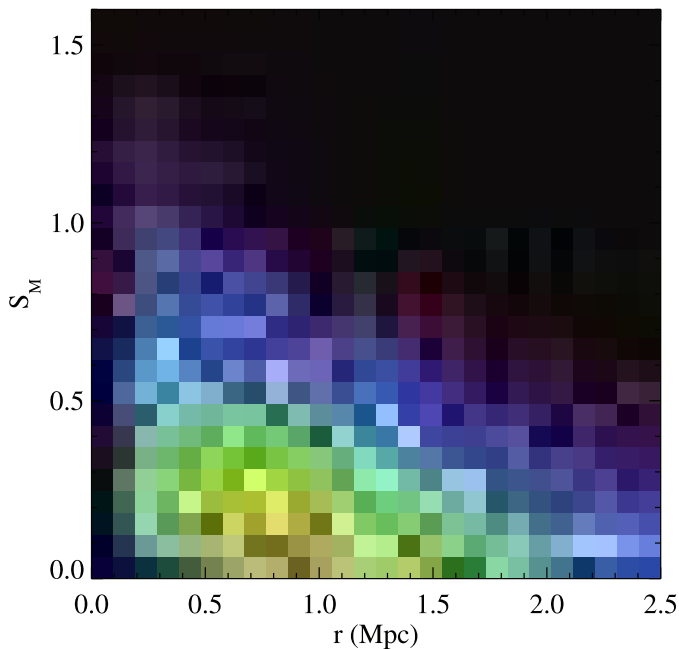


FIG. 21.— Simulated  $S_M$  vs. clustercentric radius for early-type dwarfs in the Millennium Simulation. Colors from purple to red symbolize lower to higher numbers of simulated dwarfs in a particular cell in this diagram. Like in the observations (Figs. 10 and 11), there is a tendency for dwarf galaxies with higher  $S_M$  to be at the center of the cluster.

is a result of a nonlinear metallicity-color relation. Although our data do not constrain this hypothesis, we do not observe a clear distinction between blue and red GC fractions across galaxy mass or environment.

## 6. CONCLUSIONS

We have measured globular cluster specific frequencies ( $S_N$ ), luminosity fractions ( $S_L$ ), and stellar mass fractions ( $S_M$ ) for 100 early-type galaxies in the ACS Virgo Cluster Survey. These galaxies span the mass range from giants to dwarfs ( $-22 < M_B < -15$ ), and these represent the largest homogeneous catalog of GC number and mass fractions to date. We have studied these quantities as a function of galaxy mass and environment and find the following:

1. Globular cluster fractions can be high ( $S_{N,z} > 2$ ) for both high- and low-luminosity early-type galaxies but are universally low for intermediate-luminosity galaxies ( $-22 < M_z < -19$ ).
2. There is a large spread in GC fraction in early-type dwarfs ( $0 < S_{N,z} < 5.5$ ), which can be understood as an underlying dependence on environment. Almost all dwarfs ( $M_z > -19$ ) with  $S_{N,z} > 2$  are within a projected radius of 1 Mpc from M87 and the cluster center (or 1.5 Mpc in 3D cluster radius). The spatial distributions of GCs in the high- $S_N$  dwarfs are centrally concentrated, showing that they are intrinsic to the host dwarf and not interloping GCs from M87 or an intracluster population. We do not detect higher GC fractions in dwarfs around the other massive ellipticals in the cluster. We present this as evidence that GC formation in low-mass galaxies is biased toward the densest environments.
3. Galaxies within  $\sim 40$  kpc of M87 and M49 have few or no GCs and are likely to have had their GC systems tidally stripped by their giant neighbors. Galaxies out to  $R_p \sim 100$  kpc from M87 (VCC 1185) also appear to be affected.

4. Analyzing the blue and red GC populations separately, we find that the fraction of red GCs increases with galaxy luminosity until  $M_z \sim -22$ , at which point it flattens or declines. Trends in GC fraction are dominated by the blue GCs, although the  $S_{N,z}$  of red GCs does exhibit similar relative enhancement for massive galaxies and possibly shows a weak enhancement for central dwarfs.

5. We use a globally averaged  $\mathcal{M}_h/\mathcal{M}_{G^*}-\mathcal{M}_{G^*}$  relation from halo occupation studies to test whether trends in GC fraction can be explained by the assumption that the mass in GCs is directly proportional to the halo mass,  $\mathcal{M}_{GC} \propto \mathcal{M}_h$ . While this may be able to explain the mean trend in  $S_M$ , it is unable to account simultaneously for all dwarfs, particularly the dwarfs with high GC fractions in the cluster's central regions.

6. Comparisons with semianalytic models of galaxy formation in the Millennium Simulation show that early-type dwarfs in the central 1 Mpc of massive galaxy clusters are expected to be older and have higher peak star formation rates and SFR surface densities than their counterparts on the outskirts of the cluster. The higher stellar mass fractions in globular clusters for central dwarfs can be explained if higher SFR surface densities are responsible for more efficient formation of massive star clusters.

7. The peak SFR surface density in simulated dwarfs occurs before the peak SFR, which we propose as an explanation for why GCs are, on average, both older and more metal-poor than the field stars in their host galaxies.

We present a picture of globular cluster system formation in the Virgo Cluster, where the highest GC mass fractions are formed in the oldest systems. These progenitors formed a larger fraction of their stars at higher peak SFR surface densities and also have star formation suppressed at later times. In regions of high density, halos collapse and star formation starts earlier, and the mechanisms for truncating and suppressing subsequent star formation are stronger. For all dwarfs, the highest SFR surface densities occur earlier than the peak SFR, naturally producing the metal-poor GC populations we see today: old ( $>10$  Gyr) and more metal-poor than the bulk of the field stars. Merging of many low-mass progenitors with high GC fractions produces the extremely high  $S_N$  for blue GCs seen in the cD galaxy M87. They may also, at a lower level, produce the elevated blue GC fractions seen in other massive ellipticals. Future detailed simulations of GC system formation will be crucial to test, in a quantitative way, the scenarios of galaxy and GC system formation that we are beginning to assemble.

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*Facilities:* HST(ACS,WFC2)

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